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Proton-detected MAS NMR experiments based on dipolar transfers for backbone assignment of highly deuterated proteins



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ABSTRACT

Proton-detected solid-state NMR was applied to a highly deuterated insoluble, non-crystalline biological assembly, the *Salmonella typhimurium* type iii secretion system (T3SS) needle. Spectra of very high resolution and sensitivity were obtained at a low protonation level of 10–20% at exchangeable amide positions. We developed efficient experimental protocols for resonance assignment tailored for this system and the employed experimental conditions. Using exclusively dipolar-based interspin magnetization transfers, we recorded two sets of 3D spectra allowing for an almost complete backbone resonance assignment of the needle subunit PrgI. The additional information provided by the well-resolved proton dimension revealed the presence of two sets of resonances in the N-terminal helix of PrgI, while in previous studies employing ¹³C detection only a single set of resonances was observed.

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1. Introduction

Various sophisticated isotope labeling schemes, such as carbon selective labeling [1–3] and deuteration [4,5], facilitate the study of complex biomolecules by solid-state NMR (ssNMR). These approaches require, in turn, development and revision of existing protocols for an optimal performance. For complex protein investigations, protein perdeuteration with subsequent proton back-substitution on hydrogen exchangeable sites using a buffer containing a defined H₂O–D₂O mixture is a very efficient approach to reduce the proton density and attenuate proton–proton dipolar couplings. For such protein systems, traditional ssNMR approaches, limited to the observation of ¹³C and ¹⁵N nuclei, can be extended to proton detection [6,7], which provides better sensitivity compared to ¹³C detection and improved resolution via the additional dimension. Even without high power proton decoupling, proton and nitrogen line widths in strongly deuterated proteins can be as sharp as 20 Hz and 12 Hz, respectively [4,8]. Since no high power proton decoupling is required, long chemical shift evolution times can be used to obtain highly resolved multidimensional spectra. Relaxation studies on perdeuterated proteins provide unique information about dynamics in the solid state [9–11], which is not

accessible in protonated samples at the same level of accuracy. High quality data has been obtained on longitudinal [12] and differential transverse relaxation [10,13] and order parameters [10,14]. These data can be used for an extended model-free relaxation analysis [9,10]. Furthermore, proton–proton distance measurements were performed in structural elucidations of several deuterated microcrystalline proteins [11,15–18].

The first essential step in structural NMR studies of proteins is the sequential assignment of the resonances. Even though deuterated proteins have only been studied for a relatively short time by solid-state NMR [19,20], already a plethora of different methods for ¹H, ¹³C and ¹⁵N assignment exists [15,18,21–26]. These approaches can be classified into two categories according to the protonation level and the accompanied typical experimental parameters: (1) fully protonated proteins [21,26,27] or perdeuterated proteins with complete reprotonation on exchangeable sites [15,18], which require ultrafast MAS frequencies (40–60 kHz) and very high external magnetic fields; (2) perdeuterated proteins with partial reprotonation on exchangeable sites (ca. 20–30%) [22–25], which can be studied even at low external magnetic fields of 9.4 T and moderately high MAS rates between 20 and 32 kHz, achievable with 3.2 mm and 2.5 mm probe heads. The rotor size is an essential parameter in protein studies, since they are often limited by the sensitivity and require therefore the highest possible sample amount [28]. In general, the sensitivity and resolution of proton-detected spectra are additionally depending on the MAS rate, protonation level, external magnetic field and the homogeneity of the investigated sample.

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All assignment approaches are based on the joined analysis of proton-detected triple resonance three dimensional (3D) experiments correlating H, N, CA, CO and rarely CB atoms from residues i and $i - 1$. This allows for a “sequential walk” along the protein backbone resonances. In perdeuterated proteins with 100% re-protonation, (H)CANH and (H)CONH 3D spectra can be recorded using straightforward heteronuclear cross polarization transfer for each step, including the initial polarization of CA and CO [15,18,21]. (H)CA(CO)NH and (H)CO(CA)NH spectra require a specific CA–CO magnetization transfer step, which can be achieved by pure through-bond INEPT transfer [18], utilizing the relatively high CA–CO scalar coupling or by through-space transfer based on DREAM [29] or HORROR [30] recoupling, optimized for one bond transfer [21]. The sensitivity of these experiments compared with (H)CANH and (H)CONH is determined by the efficiency of the CA–CO transfer which is around 20–60%. In close analogy, similar approaches for CA–CB magnetization transfer are used in (H)CB(CA)NH or (H)CB(CACO)NH experiments [21]. Solution state assignment approaches employing only scalar couplings are also applicable at ultrafast spinning and low reprotonation degree [10,31].

In partially reprotonated samples, (H)CANH and (H)CONH spectra can be readily recorded as in 100% reprotonated samples. Obtaining correlations between CA_{i-1} and $(HN)_i$, however, is a more challenging task, due to the low proton density, which limits the initial polarization of CA_{i-1} . The magnetization transfer from H_{i-1} is reduced because of the low probability of proton occurrence, while the transfer from the weakly coupled H_i is inefficient due to the longer distance and dipolar truncation [32] caused by the strong dipolar coupling between H_i and CA_i . The same issue affects correlations between CO_i and $(HN)_i$. At the same time, using a low reprotonation degree is a pre-requirement imposed by employing widely spread 3.2 mm and 2.5 mm probes with limited MAS rates. Thus, sequential resonance assignment of fractionally protonated samples at moderately high MAS rates (<32 kHz) remains a considerable challenge. In the following we will briefly review the existing solutions.

Reif and coworkers implemented solution-like NMR approaches [22–24], utilizing only hetero- and homonuclear scalar couplings to obtain $H_iN_iCO_{i-1}$, $H_iN_i-CA_i/CA_{i-1}$, $H_iN_iCO_{i-1}CA_{i-1}$, $H_iN_i-CA_iCO_i/CA_{i-1}CO_{i-1}$ and $H_iN_i-CA_iCB_i/CA_{i-1}CB_{i-1}$ correlations, providing extensive inter-residual connectivities for the small perdeuterated microcrystalline protein α -SH3 with fractional proton back-substitution. Outstanding sharp lines, extremely low transverse relaxation rates, increased by elevated temperatures of 25–30 °C during the experiments and exploiting paramagnetic relaxation enhancement (PRE) shortening longitudinal relaxation times by a factor of 5–8 [33], facilitated the applicability of this approach. One drawback of this methodology is that the nitrogens with relatively short transverse relaxation times are simultaneously coupled to two CAs, from the same and the previous residue, via small scalar couplings. This results in split transfers and reduced spectral resolution due to the presence of both $^{15}N-^{13}CA_i/^{13}CA_{i-1}$ correlations. While this approach yielded excellent results for α -SH3, it may be difficult to apply to other samples with broader lines and faster transverse magnetization decays [34]. Linsler [25] has recently suggested to use long range proton carbon cross polarization in order to correlate $(HN)_i$ with CA_{i-1} , CO_i or CB_{i-1} . The experimental schemes, namely hCxhNH, hCAhNH and hCohNH were tested on a Cu(Edta) doped sample of α -SH3 and provided valuable information for resonance assignment. However, dipolar truncation effects [32] arising from strong H_i and CA_i/CO_{i-1} couplings may significantly suppress through-space magnetization transfer between H_i and CA_{i-1}/CO_i in both directions.

Another assignment approach introduced by Ladizhansky and co-workers [35] is based on correlations of H_iN_i with CA_i , CO_{i-1}

and double quantum (DQ) coherences of $CO_{i-1}-CA_{i-1}$, CA_i-CO_i and CA_i-CB_i spins. To create DQ coherences, the SPC-5 sequence is used [36], which employs RF field strength 5 times higher than the MAS frequency, which may pose a restriction to the employed experimental conditions.

The achieved success of the deuteration approach in combination with proton detection motivated us to apply it to an insoluble, non-crystalline biological assembly, the *Salmonella typhimurium* type III secretion system (T3SS) needle [37]. The 80 residue subunit PrgI was perdeuterated and uniformly ^{13}C , ^{15}N labeled with subsequent 10% or 20% proton back-substitution at exchangeable sites using a buffer containing a corresponding H_2O/D_2O mixture. In the present work, we have developed two efficient experimental approaches for backbone assignment of such a fractionally reprotonated system at very fast MAS (25–28 kHz) and moderate external magnetic field (600 MHz 1H Larmor frequency), making use of the high sensitivity and resolution due to proton detection as well as the high efficiency of dipolar-based inter-nuclear magnetization transfer. The first introduced method exploits “out-and-back” $^{13}CA-^{13}CO$ dipolar based magnetization transfer optimized specifically for directly bonded nuclei, yielding high sensitivity and resolution spectra with unambiguous information about inter-spin connectivity. The second approach employs long range H–CO and H–CA magnetization transfers to obtain high sensitivity spectra for backbone sequential assignment.

2. Materials and methods

2.1. Sample preparation

Expression, purification and polymerization of perdeuterated ^{15}N - and ^{13}C -labeled wild-type PrgI protein was performed as described before [3,38,39]. A low protonation degree in the samples was achieved according to the approach described earlier by Reif and coworkers [4]. Two samples were prepared using buffers with a fraction of either 10% or 20% H_2O . Approximately 12 mg of the 10% H_2O sample were packed into a 3.2 mm rotor. Approximately 25 mg and 10 mg of the 20% H_2O sample were packed into 3.2 mm and 2.5 mm rotors, respectively.

2.2. Solid-state NMR spectroscopy

Solid-state NMR experiments were conducted on 850 MHz, 800 MHz and 600 MHz (1H Larmor frequency) spectrometers (Bruker Biospin, Germany) equipped with (1H , ^{13}C , ^{15}N) triple-resonance 2.5 mm (600 MHz) or 3.2 mm (800, 850 MHz) probes. Samples were spun at rates of 25–28 kHz (2.5 mm probe) or 20 kHz (3.2 mm probes). The effective sample temperature was 11 ± 2 °C as measured by the temperature-dependent water proton resonance relative to an internal DSS reference [40]. Chemical shift referencing was achieved using also the internal DSS reference.

The presented experiments are based on proton detection and the constant time approach, as introduced by Zilm and co-workers [20] for optimization of residual water proton signal suppression. Magnetization transfer steps between different nuclei were achieved by recoupling of hetero- or homo-nuclear dipolar interactions. For this purpose common recoupling techniques such as cross polarization (CP) [41], SPECIFIC-CP [42], HORROR or DREAM [29,30] and band-selective homonuclear CP (BSH-CP) [38] were used. Details about RF field strengths and shaped pulses are listed in Table S1 (Supporting Information) for each experiment. For homonuclear double-quantum transfer, continuous ^{13}C RF irradiation was applied in the middle of the CA band, while the CO magnetization was aligned with a trim pulse along the effective RF field defined by chemical shift offset and RF field strength. No

Table 1

Average proton and nitrogen line widths (full width at half-maximum, FWHM) in spectra obtained under different experimental conditions.

^1H , %	B_0 , MHz	ω_r , kHz	^1H FWHM, Hz	^{15}N FWHM, Hz	H-N/J decoupling on ^1H channel
10	800	20	58.4 ± 21.2	19.2 ± 7.5	Single hard 180°
20	800	20	63.9 ± 21.0	22.9 ± 8.4	Single hard 180°
20	600	28	47.0 ± 14.0	18.8 ± 8.7	Single hard 180°
20	600	28		16.3 ± 8.7	WALTZ-64

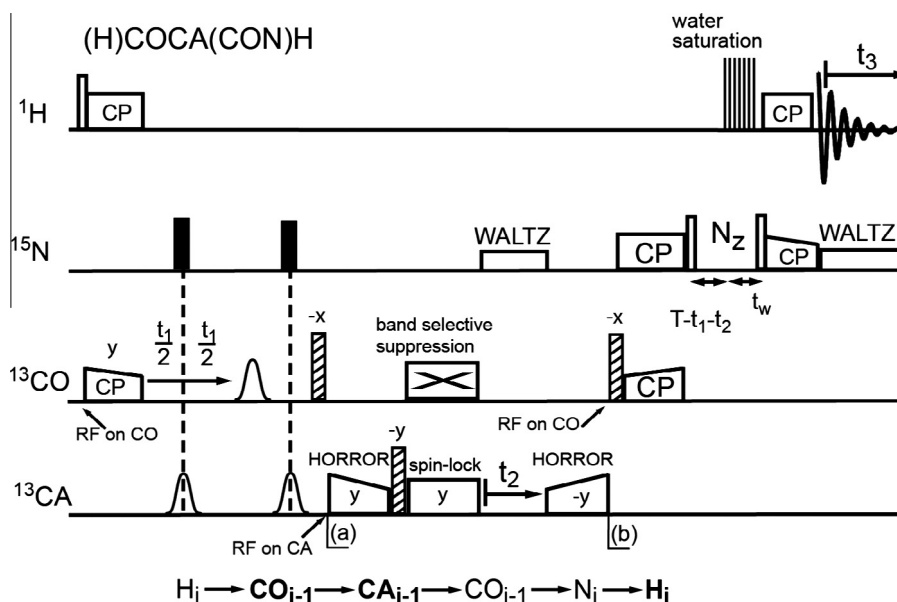


Fig. 2. Pulse sequence for obtaining (H)COCA(CON)H 3D proton-detected inter-residual H^{N} , CA_{i-1} , CO_{i-1} correlation spectra employed for sequential backbone assignment. Open, filled and hatched bars represent 90° , 180° and hard trim pulses, respectively. Bell shapes correspond to band-selective 180° pulses. All magnetization transfer steps are based on recoupling of hetero- or homo-nuclear dipolar interactions. Band-selective suppression of remaining CO magnetization was used after the CO-CA magnetization transfer. Straightforward modifications allow to obtain pulse schemes for (H)CACO(CAN)H and (HCO)CA(CO)NH correlation experiments, as described in detail in the [Supporting Information](#). The coherence transfer steps are schematically indicated at the bottom of the figure and observed nuclei are represented in bold.

the latter step. This concept is more efficient compared to direct CP from protons to CA_{i-1} , because magnetization transfer from the $(i-1)$ hydrogen is very small due to the low probability of simultaneous proton occurrence at H_i and H_{i-1} in a 20% reprotated sample. The direct transfer from the weakly coupled H_i is dipolar truncated [32] because of the stronger dipolar coupling between H_i and CA_i . For back polarization transfer from CA_{i-1} to H_i we used 3 consecutive one-bond transfer steps, which unambiguously correlate CA_{i-1} , CO_{i-1} , N_i and H_i . In the presented experiment (Fig. 2), NC scalar couplings were removed by 180° ^{15}N hard pulses during the CO chemical shift evolution, while carbon band-selective 180° pulses [43] were used to remove the CO-CA J couplings. After the chemical shift evolution, the CO magnetization was flipped by a 65° hard trim pulse along the effective RF field determined by the ^{13}C RF field applied on CA for HORROR/DREAM recoupling. After CO-CA transfer, the remaining CO magnetization was suppressed by band-selective recoupling achieved by a 21.5 kHz spin lock field applied on CA, which results in an effective RF field of 28 kHz (i.e. equal to the MAS rate) acting on CO. A more detailed description of “out” CO-CA and “back” CA-CO transfers with band-selective magnetization suppression (fragment between points a and b in Fig. 2) is shown in the [Supporting Information](#) (Fig. S5). The overall efficiency of the CO-CA-CO module was 33%. During CA chemical shift evolution, NC J couplings were removed by WALTZ-16 [48], but no CO-CA J decoupling was applied. “Back” CA-CO magnetization transfer was achieved by DREAM/HORROR recoupling under continuous RF irradiation applied in the middle of the CA band. By means of a 65° hard trim pulse, the CO magnetization was then flipped to the transverse plane

for subsequent CO-N SPECIFIC CP [42]. The ^{15}N magnetization was preserved along the Z axis during suppression of the remaining water proton magnetization by a proton pulse train [49,50], as shown in detail in Fig. S2 B ([Supporting Information](#)). The intensity of the first FID was 9.5% compared with the signal in an (H)NH experiment. The $\text{H}_i\text{CO}_{i-1}\text{CA}_{i-1}$ correlation spectrum was recorded with excellent resolution and sensitivity within a short time of 42 h, using 16 scans, a 2.5 s inter-scan delay, and maximum isotropic chemical shift evolution times of 45 ms, 5.72 ms and 17.4 ms for ^1H , CA and CO, respectively.

The same approach can be readily applied to correlate H_i and CO_i nuclei. The straightforward modifications of the scheme described above for (HCO)CA(CO)NH and (H)CACO(CAN)H experiments, as used in the present study, are explained in detail in Fig. S4 and Table S3 ([Supporting Information](#)). In general, the transfer efficiency between low γ nuclei was around 45–60% due to the long polarization life times in deuterated proteins. The high transfer efficiency resulted in spectra with high sensitivity even in (HCO)CA(CO)NH, (H)COCA(CON)H and (H)CACO(CAN)H experiments, accommodating 5 transfer steps. The sensitivity of each experiment relative to (H)NH was between 7.5% and 9.5% (see Table S3, [Supporting Information](#)), which is quite high for such complex ssNMR correlation experiments. As a result, the required measurement time was only 42–61 h for a single 3D experiment. In the recorded spectra, each amide proton or amide group correlates only to one CA and/or CO nucleus, yielding twice higher resolution compared to solution-like [24] or dipolar-based methods [25] reported previously for proton-diluted proteins. We found that $\text{H}_i\text{CO}_{i-1}\text{CA}_{i-1}$ and $\text{H}_i\text{CA}_i\text{CO}_i$ correlations are preferable to

$H_iN_iCA_{i-1}$ and $H_iN_iCO_i$ correlations in order to establish H_i connectivity with CA_{i-1} and CO_i , because of the higher resolution of H_iCO_{i-1} and H_iCA_i spectra compared to H_iN_i spectra. A set of four 3D (H)CANH, (H)CONH, (H)COCA(CON)H and (H)CACO(CAN)H experiments yields the necessary information for the backbone resonance assignment using the frequencies of CO and CA nuclei to establish the connectivity between subsequent residues. An additional fifth 3D (HCO)CA(CO)NH experiment was recorded to compare the results with the (H)COCA(CON)H spectrum.

We note that the proposed methodology can be performed as well if band-selective homonuclear (BSH) [38] CP is used instead of HORROR/DREAM for CO–CA magnetization transfer, as required in case of higher external magnetic field used with a 3.2 mm probe providing slower MAS rates. As an example a (HCO)CA $_{i-1}$ (CON)H $_i$ 2D correlation spectrum obtained on the 800 MHz spectrometer at 20 kHz MAS is presented in Fig. S7.

Fig. 3 illustrates a sequential walk along the backbone between residues A73 and V65, based on the set of five 3D spectra. The resonance assignment strategy is based on keeping two frequencies per assignment step constant while adding one new frequency. Interspin connectivity between H_i , N_i , CA_i and CO_{i-1} is established by $H_iN_iCA_i$ and $H_iN_iCO_{i-1}$ experiments. The $H_iN_iCA_{i-1}$

and $H_iCO_{i-1}CA_{i-1}$ spectra provide both the CA_{i-1} frequency. They are therefore not both mandatory to obtain a complete set of sequential assignment experiments, but they mutually corroborate the assignment information of each other. The $H_{i-1}CA_{i-1}CO_{i-1}$ spectrum provides the H_{i-1} frequency and the $H_iN_iCA_i$ spectrum yields the N_{i-1} frequency. The high quality of the set of five 3D spectra allowed for the unambiguous and reliable assignment of all 1H , ^{13}CA , ^{13}CO and ^{15}N backbone resonances (Table S4, Supporting Information; BMRB Entry: 18276; contains also stereospecifically assigned methyl resonances of Val and Leu [51]), except of ^{15}N of the three prolines and all nuclei of the first 4 amino acids, because of their flexibility [52] and a proline in the 4th position. We found, that residues 6, 19, 22–31 and 36 exhibit two independent sets of signals. The doubling is mostly restricted to the 1H dimension, which explains why the second polymorph was not detected before in protonated samples where the ssNMR spectra were based exclusively on the observation of ^{15}N and ^{13}C [39]. We found only slight chemical shift differences between earlier studied protonated [39] samples and the currently studied deuterated sample, which are completely explainable by the deuterium isotope effect [53]. The ^{13}CA and ^{15}N resonances were up-field shifted on average by 0.31 ppm and 0.29 ppm, respectively.

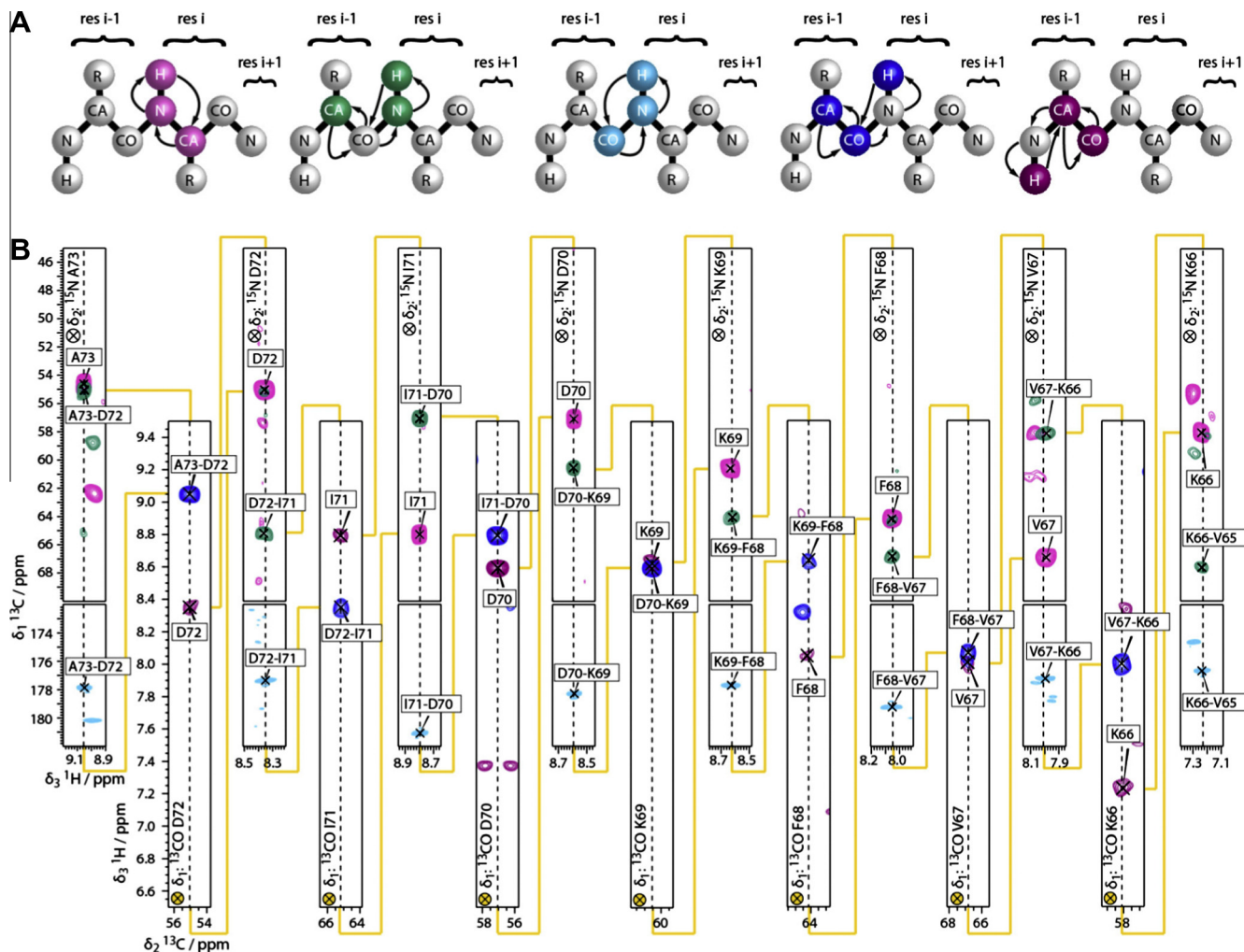


Fig. 3. Panel (A) shows from the left to the right $H_iN_iCA_i$, $H_iN_iCA_{i-1}$, $H_iN_iCO_{i-1}$, $H_iCO_{i-1}CA_{i-1}$ and $H_{i-1}CA_{i-1}CO_{i-1}$ connectivities obtained by (H)CANH, (HCO)CA(CO)NH, (H)CONH, (H)COCA(CON)H and (H)CACO(CAN)H proton-detected 3D experiments. Arrows represent spin polarization transfer pathways. Panel (B) shows strip plots for a sequential backbone walk between the residues A73 and V65 based on the correlations presented in panel (A). The peaks are color coded according to the experimental type, as presented in panel (A). The upper row contains signals from $H_iN_iCA_i$, $H_iN_iCA_{i-1}$ and $H_iN_iCO_{i-1}$ experiments, whilst the lower row presents signals from $H_iCO_{i-1}CA_{i-1}$ and $H_{i-1}CA_{i-1}CO_{i-1}$ experiments. The isotropic chemical shift value in the indirect dimension is represented in the upper left or lower left corner. Dashed black lines are an eye guide along constant 1H or ^{13}C frequency. Yellow lines and yellow indirect dimension symbols refer to conserved isotropic chemical shifts between upper and lower rows. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

3.3. Design of 3D pulse sequences for backbone resonance assignment using long range H–C dipolar transfers

In this section we demonstrate an efficient approach for backbone sequential assignment based on long range H_i – CA_{i-1} and H_i – CO_i cross polarization transfer. The (H)COCAH pulse sequence and main magnetization transfer pathways achieved during this experiment are presented in Fig. 4. Initial proton magnetization was directly transferred to CO and after chemical shift evolution further distributed to CA by HORROR/DREAM recoupling (or at higher field and lower MAS rate: with BSH-CP) with 60% efficiency of the latter step. After subsequent isotropic chemical shift evolution, ^{13}C magnetization was preserved along the Z axis for

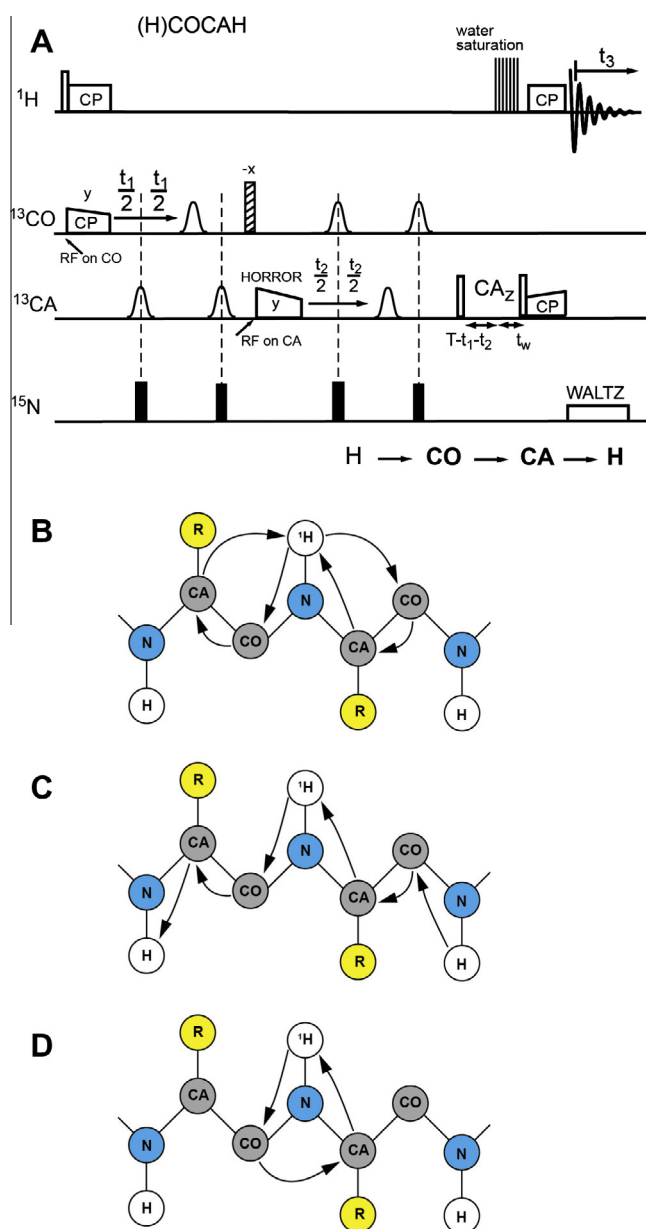


Fig. 4. Panel (A) shows the pulse sequence for the 3D (H)COCAH experiment for obtaining $H^N CA_{i-1} CO_{i-1}$ and $H^N CA_i CO_i$ correlations employed for sequential backbone assignment. Open, filled and hatched bars represent 90° , 180° and hard trim pulses, respectively. Bell shapes correspond to band-selective 180° pulses. All magnetization transfer steps are based on recoupling of hetero- or homo-nuclear dipolar interactions. The general coherence transfer scheme is schematically indicated at the bottom of the panel (A) while panels (B)–(D) show the most probable site specific magnetization transfer pathways.

suppression of the remaining water proton magnetization by a proton pulse train [49,50], as shown in detail in Fig. S2 B (Supporting Information). For back polarization transfer from CA to H we used direct cross polarization. This approach provides $H^N CO_{i-1} CA_{i-1}$ and $H^N CO_i CA_i$ correlations within a single spectrum, as it is visualized in Fig. 4. Panel B shows the most probable magnetization transfer pathways, while panel C represents possible magnetization transfers if adjacent amide hydrogen positions are occupied by protons. If we assume that H_i – CA_{i-1} and H_i – CO_i as well as H_i – CA_i and H_i – CO_{i-1} transfer efficiency are pairwise equal, then the intraresidual $H_i CA_i CO_i$ correlations are expected to be more intense than the $H_i CA_{i-1} CO_{i-1}$ correlations. We found that several residues exhibit $H^N CO_{i-1} CA_i$ correlations as well (see Fig. 4D), despite the expected efficient truncation of magnetization transfer between CO_{i-1} and CA_i due to the strong CO_{i-1} and CA_{i-1} dipolar coupling. One explanation is the relatively low MAS rate of 25 kHz, which results in a relatively weak ^{13}C RF field of ca. 6 kHz. At this condition, isotropic chemical shift variations can significantly affect effective RF field and the efficiency of homonuclear dipolar recoupling. In case of unfavorable isotropic chemical shift values, recoupling between CO_{i-1} and CA_{i-1} might be not optimal, and thus the truncation of magnetization transfer between CO_{i-1} and CA_i nuclei can be partially canceled. Higher MAS frequencies require stronger ^{13}C RF fields to fulfill the HORROR condition, thus the above described mechanism for magnetization transfer between CO_{i-1} and CA_i nuclei could be suppressed. To ensure long range magnetization transfer between H_i – CA_{i-1} and H_i – CO_i we employ relatively long contact times of 4.2–4.5 ms, as optimized in a 1D version of the (H)COCAH experiment. During CO and CA chemical shift evolution periods, NC scalar couplings were removed by 180° ^{15}N hard pulses, while carbon band-selective 180° pulses [43] were used to remove the CO–CA J couplings. The intensity of the first FID was 20% compared with the signal in an (H)NH experiment. The spectrum was recorded with excellent resolution and sensitivity within a short time of 64 h, using 16 scans, a 2.5 s inter-scan delay, and maximum isotropic chemical shift evolution times of 45 ms, 22.7 ms and 7.0 ms for 1H , CO and CA, respectively.

In the recorded spectrum, each amide proton correlates to two CA–CO couples yielding twice less resolution compared to the method introduced above. On the other hand this approach yields at least 1.5 times higher sensitivity for $H^N CO_{i-1} CA_{i-1}$ and $H^N CO_i CA_i$ correlations (see Table S3), despite the fact that it is based on truncated long-range 1H – $^{13}CA/^{13}CO$ transfer. The higher sensitivity of this method compared to the above discussed approach is due to fewer transfer steps and the fact that it contains both H_i – CA_{i-1} and H_i – CO_i correlations. Also the experiment is very straightforward to setup and implement. A set of three 3D (H)CANH, (H)CONH and (H)CACOH experiments provides the necessary information for the backbone resonance assignment using the frequencies of CO and CA nuclei to establish the connectivity between subsequent residues. During CA chemical shift evolution in the (H)CANH experiment, CA–CO J couplings were removed by band selective 180° pulses as plotted in Fig. S3. (H)CONH and (H)CACOH experiments were recorded within a short time of 26 and 18 h, respectively.

A sequential walk along the backbone between residues Q48 and P41, based on the set of three 3D spectra is demonstrated by aligned strips in Fig. 5. If CO_{i-1} and H_i frequencies are known, the $H^N CO_{i-1} CA_{i-1}$ and $H^N CO_i CA_i$ correlations, represented in blue color, provide the CA_{i-1} and H_{i-1} frequency, respectively. $H_i N_i CA_i$ (green) and $H_i N_i CO_{i-1}$ (magenta) correlations establish the N_{i-1} and CO_{i-2} frequencies. The high quality of the spectra allowed for the unambiguous and reliable assignment of most 1H , ^{13}CA , ^{13}CO and ^{15}N backbone resonances, except of ^{15}N of the three prolines, all nuclei of the first 4 amino acids, because of their flexibility [52] and a proline in the 4th position, residues 22 and 23, as well

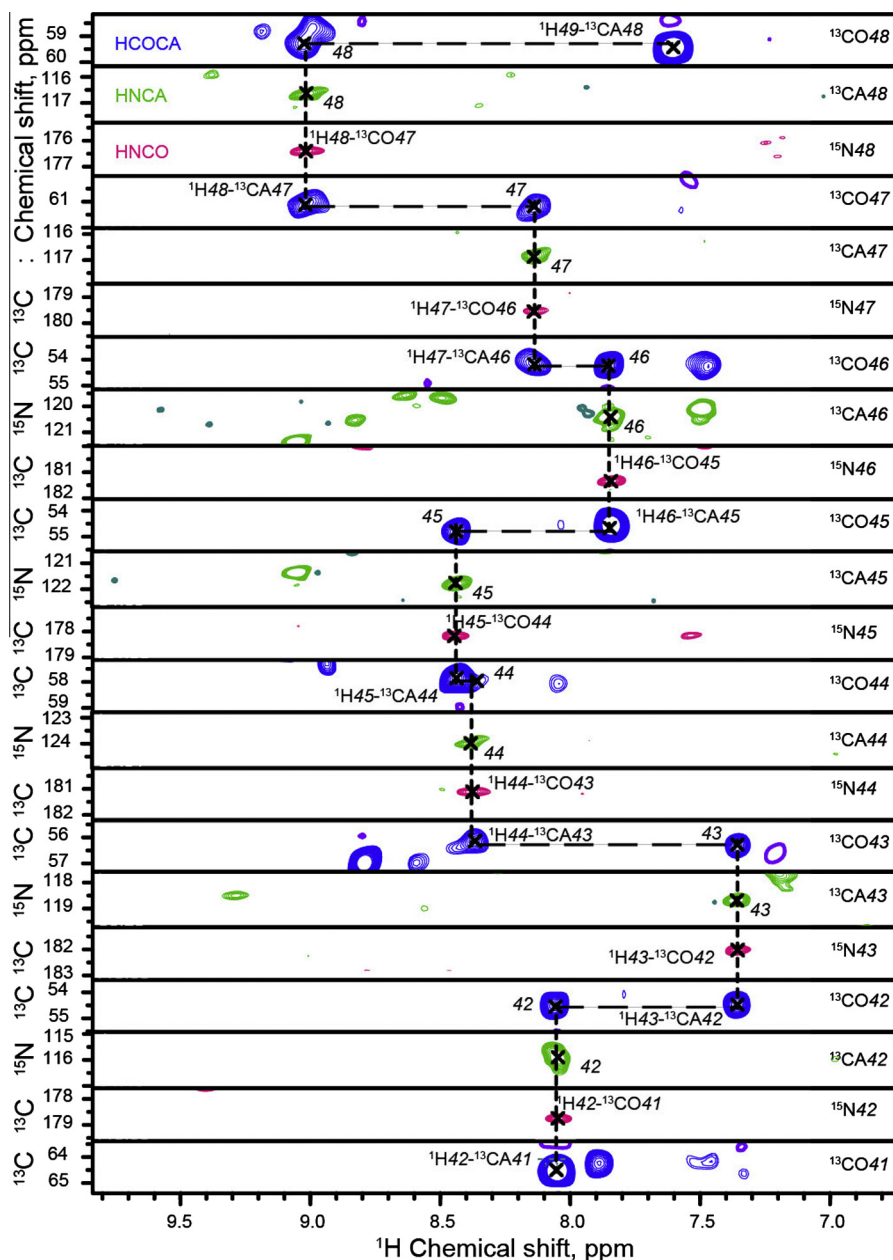


Fig. 5. Strip plots for a sequential backbone walk between the residues Q48 and P41 based on the correlations obtained by long range H–C CP. The peaks obtained from (H)COCAH, HCANH and HCONH experiments are shown in blue, green and magenta, respectively. The isotropic chemical shift value in the indirect dimension corresponds to the nucleus given on the right side of the strips. Dashed black lines are an eye guide along conserved ^1H or ^{13}CA frequencies during assignment steps. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

residues 50 and 62, whose signals were only observable in the direct (H)NH 2D correlation spectrum during this second measurement session. We could confirm that residues 6, 19, 24–31 and 36 exhibit two independent sets of signals.

3.4. Efficiency of different experimental schemes

Fig. 6 compares the efficiency of the above introduced HCACONH, HCOCACONH and HCOCAH experiments and presented earlier pulse sequences: (i) HCOCACONH based on out-and-back CO–CA INEPT transfer [54] and (ii) HCACONH and HCOCANH experiments employing DREAM for magnetization transfer between CO and CA spins [21]. All heteronuclear transfers in these pulse schemes were achieved by CP. Detailed descriptions of the latter experiments are given elsewhere [21,54], their analysis is

presented in the introduction and pulse sequences are depicted in the Supporting Information, Fig. S8. All data were obtained at an MAS rate of 25 kHz and were recorded within the same experimental session. Number of scans, experiment type and efficiency relative to the HNH experiment are given in the panels. Essentially, the efficiency of all experiments is similar and a choice of the pulse sequence is then mostly determined by the protonation level and the transverse coherence decay time T_2' of ^{13}C .

Strong ^{13}CO – ^{13}CA J couplings and a long ^{13}C T_2' suggest the use of INEPT for “out-and-back” (CO)CA(CO) steps [54]. Considering a CO transverse relaxation time of 25–40 ms, the expected efficiency of the (CO)CA(CO) block is 48–63% in case of ideal performance of the applied pulse train. However, non-negligible magnetization losses can occur due to RF field inhomogeneity, because the INEPT block comprises 6 pulses. In general, the efficiency of this method

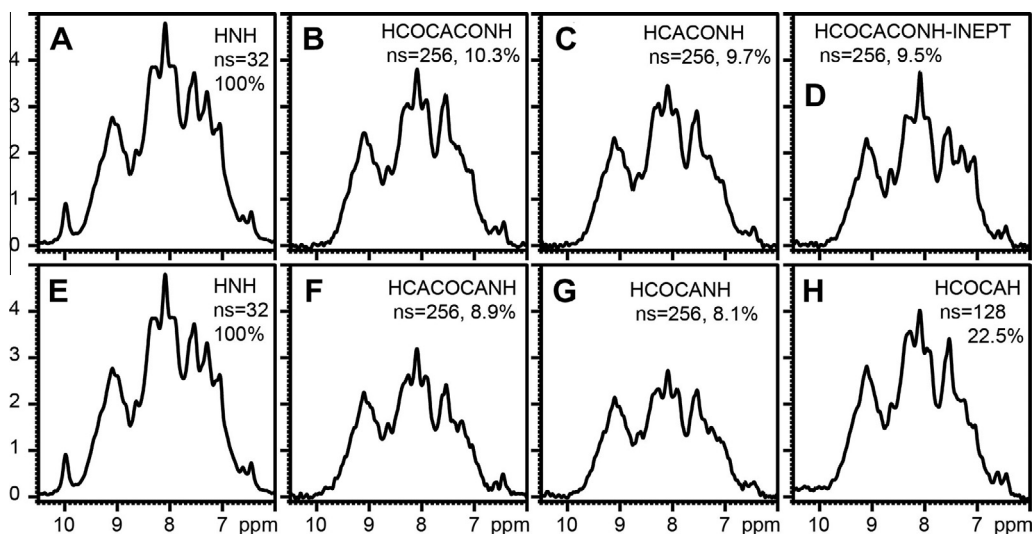


Fig. 6. Proton detected 1D spectra obtained by different pulse sequences on deuterated PrgI T3SS needles. Panels A and E show the reference HNHN spectrum (Fig. S1). Spectra in the panels B and F were obtained by using dipolar-based CA–CO out-and-back transfers (Figs. 2 and S4). Spectra in the panels C and G were obtained by using direct CP transfer from remote protons to CA and CO to create inter-residual H_iCA_{i-1} and intra-residual H_iCO_i connectivity [21]. D) The spectrum was recorded using out-and-back CO–CA INEPT transfer in a HCOCACONH experiment [54]. H) The spectrum was obtained by employing long range C–H transfers to establish H–C inter- and intra- residual correlations, according to the pulse sequence in Fig. 4.

is strongly influenced by the transverse magnetization life time, which can be quite short for some biomolecules due to specific internal motions, low temperature and presence of paramagnetic centers for relaxation enhancement. Recently, Reif and co-workers found a T_2' of ^{15}N for a number of systems within 10–15 ms [34], while T_2' of CO is typically shorter than of ^{15}N . For a CO transverse relaxation time of 10–15 ms, the expected efficiency of the (CO)CA(CO) block drops to 18–31% in case of ideal performance of the applied pulse train, thus making scalar coupling based transfer between CO and CA less practical.

As it has been discussed above, initial polarization of CA_{i-1} in H_i – CA_{i-1} correlation experiments using direct CP transfer from protons is not efficient in highly deuterated proteins, because of the low probability of H_{i-1} proton occurrence and because the transfer from the weakly coupled H_i is inefficient due to the longer distance and dipolar truncation [32] caused by the strong dipolar coupling between H_i and CA_i . The presented results demonstrate that 2-step transfer comprising H_i – CO_{i-1} and CO_{i-1} – CA_{i-1} magnetization transfers is more preferable than direct magnetization transfer to CA_{i-1} from H_i and H_{i-1} at the here considered low protonation degree. At the same time the found efficiency of HCA–CONH and HCOCANH experiments was higher than expected with the used 20% reprotonation buffer. One possible explanation is that the real protonation degree might be higher, because water can be absorbed from the atmosphere during sample preparation.

4. Conclusion

We have demonstrated that the deuteration approach applied to PrgI needles yields proton-detected solid-state NMR spectra of high resolution and sensitivity. We then designed for proteins with extremely low proton content of 10–20% on exchangeable sites two efficient assignment approaches which employ proton detection and exclusively dipolar-based magnetization transfer steps. The method based on “out-and-back” ^{13}CA – ^{13}CO magnetization transfers provides unambiguous information about inter-residual spin connectivity, high sensitivity and better resolution compared to other methods presented before, making this approach highly attractive for even more complex proteins. The approach utilizing long range ^1H – ^{13}CA and ^1H – ^{13}CO magnetization transfers yields

higher sensitivity and is very simple to setup, making this approach highly useful as well. The excellent data quality allowed for an almost complete backbone resonance assignment of the 80-residue needle subunit PrgI employing a moderate external magnetic field of 14.1 T (600 MHz ^1H Larmor frequency) and an MAS rate of 25–28 kHz. All 3D spectra were recorded in a relatively short time frame (see Table S3), ranging from 18 h to 61 h for a single experiment. A previously unidentified second polymorph of the PrgI needles was assigned, due to the additional information granted by the proton dimension. All the doubled residues are close to W5 in the 3D structure [3] which may indicate that the line splitting is due to ring current effects caused by the side chain of W5. This would explain why protons are more affected than carbons. We will investigate this effect more carefully in future studies.

5. Data deposition

Chemical shifts are deposited in the BMRB (BMRB Entry: 18276).

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Appendix A. Supplementary material

Supplementary data associated with this article can be found, in the online version, at <http://dx.doi.org/10.1016/j.jmr.2014.02.020>.

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