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Development and evaluation of a suite of isotope reference gases for methane in air

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Abstract. Measurements from multiple laboratories have to be related to unifying and traceable reference material in order to be comparable. However, such fundamental reference materials are not available for isotope ratios in atmospheric methane, which led to misinterpretations of combined data sets in the past. We developed a method to produce a suite of synthetic CH₄-in-air standard gases that can be used to unify methane isotope ratio measurements of laboratories in the atmospheric monitoring community. Therefore, we calibrated a suite of pure methane gases of different methanogenic origin against international referencing materials that define the VSMOW (Vienna Standard Mean Ocean Water) and VPDB (Vienna Pee Dee Belemnite) isotope scales. The isotope ratios of our pure methane gases range between -320 and +40% for $\delta^{2}H$ –CH₄ and between -70 and -40% for δ^{13} C–CH₄, enveloping the isotope ratios of tropospheric methane (about -85 and -47% for δ^2 H–CH₄ and δ^{13} C–CH₄ respectively). Estimated uncertainties, including the full traceability chain, are $< 1.5\,\%$ and < 0.2% for δ^2 H and δ^{13} C calibrations respectively. Aliquots of the calibrated pure methane gases have been diluted with methane-free air to atmospheric methane levels and filled into 5L glass flasks. The synthetic CH₄-in-air standards comprise atmospheric oxygen/nitrogen ratios as well as argon, krypton and nitrous oxide mole fractions to prevent gasspecific measurement artefacts. The resulting synthetic CH₄in-air standards are referred to as JRAS-M16 (Jena Reference Air Set – Methane 2016) and will be available to the atmospheric monitoring community. JRAS-M16 may be used as unifying isotope scale anchor for isotope ratio measurements in atmospheric methane, so that data sets can be merged into a consistent global data frame.

1 Introduction

Isotope ratios of CH₄ in the present and the past atmosphere (e.g. from ice cores) are a powerful tool to study the biogeochemical processes that cause the variation of CH₄ in the atmosphere (Stevens and Rust, 1982; Quay et al., 1991, 1999; Lowe et al., 1994; Sapart et al., 2012; Möller et al., 2013; Sperlich et al., 2015; Schaefer et al., 2016). Recently, two conflicting publications highlighted (i) the interpretative power when data sets from multiple laboratories are combined for spatiotemporal analysis of CH₄ isotope ratios (Kai et al., 2011) and (ii) the pitfalls when differences due to laboratory offsets are misinterpreted as spatial variability of CH₄ sources (Levin et al., 2012). Levin et al. (2012) identified calibration offsets between three laboratories by comparing their long-term observations in Antarctic background air, where the δ^{13} C of CH₄ is assumed to be free of spatial gradients. However, this technique is a temporary work-around that excludes the use of data sets from laboratories without a history of observations in Antarctica or a traceable link to Antarctic observations. This dilemma could be solved if suitable reference materials (RMs) were available to all laboratories that measure isotope ratios of atmospheric CH₄.

Certified reference materials (CRMs) are provided by the IAEA, NIST and others for many analytes. The lack of CRMs for CH₄ isotope ratios has long been recognised in the literature, ranging from pioneering papers (e.g. Craig, 1953; Schiegl and Vogel, 1970) to recent publications on analytical systems to measure isotope ratios in atmospheric CH₄ (e.g. Sapart et al., 2011; Sperlich et al., 2013; Bock et al., 2014; Tokida et al., 2014; Eyer et al., 2015) as well as papers that present and interpret such data (e.g. Levin et al., 2012; Sapart et al., 2013; Schaefer et al., 2016). In the absence of CRMs for isotope ratios of CH₄, many laboratories have developed methods to calibrate purified CH₄ against CRMs that were available as a "second-best solution", thereby accepting the shortcoming that those CRMs comprised of different physicochemical properties and are therefore not ideal (IAEA, 2003). For example, δ^{13} C–CH₄ calibrations were made against NBS 20 (limestone) and NBS 21 (graphite) by Stevens and Rust (1982), against NBS 16 (CO₂) and NBS 20 (limestone) by Quay et al. (1991), against IAEA-CO-9 (Barium carbonate) by Lowe et al. (1994), against NBS 19 (limestone) by Quay et al. (1999) and against RM 8563 (CO₂) by Sperlich et al. (2012). Dumke et al. (1989) calibrated against the natural gas mixtures NGS 1, NGS 2 and NGS 3, which were not of the highest purity level with 81, 53 and 99 % CH₄ respectively (e.g. IAEA, 2003; Brand et al., 2014). It is furthermore important to understand the variation of uncertainties of the applied CRMs, ranging from assigned values of 0.00% (NBS 19, the only primary measurement standard for VPDB) up to 0.56 % (NGS 2) (Brand et al., 2014). The situation becomes even more complicated because the δ^{13} C values of some of the applied CRMs were revised and changed by as much as 0.4% over time (e.g. NBS 21; Brand et al., 2014). As a consequence, this would require the adjustment of dependent δ^{13} C–CH₄ data. The use of different calibration methods, CRMs and the change of their assigned δ^{13} C values have undoubtedly contributed to calibration offsets between laboratories. This fact highlights the importance that applied CRMs and their δ^{13} C values are reported in the metadata of the measurement results and that their uncertainty is included in the uncertainty budget of the measurements. Fortunately the situation is more homogenous for δ^2 H–CH₄ calibrations, which were only made against CRM waters, such as VSMOW2, SLAP2 or their precursors (e.g. Schiegl and Vogel, 1970; Dumke et al., 1989; Quay et al., 1999; Sperlich et al., 2012). Brand et al. (2014) provide a comprehensive overview on the variation of δ^2 H–H₂O values and associated uncertainties. Another common method for laboratories to anchor CH₄ measurements to the VPDB or VSMOW isotope scales is to get their working standard (WS) calibrated by an external laboratory (e.g. Behrens et al., 2008; Brass and Röckmann, 2010; Bock et al., 2014; Schmitt et al., 2014; Rella et al., 2015; Brand et al., 2016). It is important to keep in mind that propagating isotope scales between laboratories also requires inclusion and propagation of the uncertainty of the respective isotope scale anchor.

In summary, the absence of unique CRMs for δ^2 H–CH₄ and δ^{13} C–CH₄ led to a diversity of calibration trajectories. Significant calibration offsets between laboratories on the order of 0.05–0.09‰ for δ^{13} C–CH₄ were identified through co-located measurements by Levin et al. (2012) and Schaefer et al. (2016), while Bock et al. (2014) reported laboratory offsets of up to 15‰ for δ^2 H–CH₄. Even though interlaboratory differences can be established experimentally, e.g. by co-located measurements or regular round robins, such comparisons are not intended to re-define local scale anchors to the VPDB and VSMOW isotope scales (WMO, 2014) and can therefore not replace a unifying scale anchor.

Until recently, a comparable problem existed for observations of isotope ratios in atmospheric CO₂. Ghosh et al. (2005) established a method to produce synthetic CO₂in-air standards, comprising of isotopically calibrated CO₂ and CO₂-free air. The concept of these CO₂-in-air standards is to provide a matrix reference material (m-RM), which is defined as RM that is mixed with matrix material to match the composition of the samples (IAEA, 2003). Since 2005, the ISOLAB of the Max Planck Institute for Biogeochemistry (MPI-BGC) in Jena, Germany, distributes a suite of m-RMs, known as JRAS (Jena Reference Air Set), which is accepted as an isotope scale anchor by the community (WMO, 2012). Calibrating against the JRAS reduces laboratory offsets and has proven a successful method to reach and maintain the compatibility goal for isotope ratios in atmospheric CO_2 (Wendeberg et al., 2013).

This paper describes an analogue method to produce synthetic CH₄-in-air standards for δ^2 H–CH₄ and δ^{13} C–CH₄, which we refer to as JRAS-M16 (short for JRAS-Methane 2016). We present new methods to calibrate a suite of isotopically different CH₄ gases, which span over a large isotopic range. We calibrate two CH₄ gases for δ^2 H and δ^{13} C and compare our results to independent calibrations made at a partnering laboratory to demonstrate the comparability of our new methods, thereby fulfilling the requirement to use two independent analytical methods during the development of quality control materials (QCMs) when CRMs are not available (IAEA, 2003). We produce synthetic CH₄-in-air standards by diluting aliquots of calibrated CH₄ with CH₄free synthetic air and include the full traceability chain in the uncertainty budget. Calibrated δ^2 H–CH₄ and δ^{13} C–CH₄ values in our synthetic CH₄-in-air standards bracket tropospheric values and enable two-point calibrations to account for scale compression effects (Coplen et al., 2006a). Our synthetic CH₄-in-air standards can be tested by other laboratories in the community; alternatively, compressed air cylinders from other laboratories can be calibrated at MPI-BGC. Our long-term strategy is to establish JRAS-M16 as m-RM for δ^2 H–CH₄ and δ^{13} C–CH₄ in the future. We hope that our efforts help the community to reach the scale anchor compat-



Figure 1. Calibration hierarchy to produce synthetic CH₄-in-air standards including links of the traceability chain. The long, central arrow shows that the primary CH₄ gases were directly calibrated against CRMs for δ^{13} C but not for δ^{2} H. The uncertainty (*U*) associated with each calibration hierarchy level is indicated by indices that are described in Sect. 2.4.

ibility goals of 1 and 0.02% for δ^2 H–CH₄ and δ^{13} C–CH₄ respectively (WMO, 2014).

2 Materials and methods

Throughout this paper, we use the terminology of "calibration" and "measurement" with different intentions. We use calibration when samples are repeatedly compared against measurement standards of the highest possible hierarchy level (possible hierarchy levels include CRMs and WSs) in order to determine the isotopic composition of the analyte under consideration of the full traceability chain. In contrast, we use the measurement term when the analysis is not necessarily based on measurement standards of highest possible hierarchy level, when the achievable uncertainty of the analysis is not of primary importance or when the uncertainty does not necessarily include the full traceability chain. For example, we use the measurement term for the experiments to establish the dependence of isotope ratios in the analyte on reactor temperatures of the analytical system.

The aim of our method is to calibrate and prepare synthetic CH₄-in-air standards, as outlined in the flow diagram of Fig. 1. Therefore, we calibrate two pure CH₄ gases for their δ^2 H–CH₄ and δ^{13} C–CH₄ isotope ratios against CRMs and WSs, where the latter are of comparable chemical composition to the former. We refer to these two CH₄ gases as primary CH₄ gases. The primary CH₄ gases are then used to calibrate a suite of pure CH₄ gases, which we refer to as secondary CH₄ gases. The analytical methods we developed for δ^2 H–CH₄ and δ^{13} C–CH₄ calibrations are based on well-established IRMS methods, thereby complying with the requirements to use established analytical systems for the production of QCMs when CRMs are not available (IAEA, 2003). Once calibrated, aliquots of both primary and secondary CH₄ gases are diluted with CH₄-free air to atmospheric CH₄ mole fractions. We analyse the resulting synthetic CH₄-in-air standards on a new analytical system that is designed to analyse atmospheric samples, thereby complying with the principle of identical treatment (PIT; Werner and Brand, 2001) during the analysis of the synthetic CH₄-in-air standards. This enables us to determine the calibration difference between JRAS-M16 and the hitherto adopted method to reference δ^2 H–CH₄ and δ^{13} C–CH₄ in atmospheric samples to the VSMOW and VPDB scales respectively. This difference represents the laboratory specific correction that has to be applied to anchor all measurements from MPI-BGC to the new JRAS-M16 scale.

2.1 Gases, reference materials and hierarchy levels of calibrations

Our study is based on a suite of CH_4 gases that differ in their methanogenic origin and therefore in their isotopic composition. We identify our CH_4 gases by names as shown in Table 1. "Biogenic" and "Fossil" have been calibrated at the Centre for Ice and Climate (CIC), which is a department of the Niels Bohr Institute at the University in Copenhagen, Denmark (Sperlich et al., 2012). These gases allow testing and evaluating the performance of the analytical systems at MPI-BGC with independent methods, which is a required control mechanism for the development of QCMs when CRMs are not available (IAEA, 2003). Six other CH_4 gases were purchased from suppliers of commercial gases or laboratory equipment (Air-Liquide, Westfalen AG, Linde, Messer, Campro Scientific) and were used as purchased or as mixtures thereof. The purity level of all our CH_4 gases is



Figure 2. Configuration of manual the two-position 10-port valve with two 1 mL sample loops shown in grey dashed box and TC/EA–IRMS system for δ^2 H–CH₄ calibration. The TC/EA–IRMS reactor (displayed as in Gehre et al., 2004) is fed either by the sample line from the 10-port valve or by the syringe via autosampler (not shown). The size of components is chosen to increase clarity.

 \geq 99.995 %. Our goals were to produce (i) a suite of CH₄ gases that encompasses the isotopic composition of tropospheric CH₄, and (ii) CH₄ gases that closely match the isotopic composition of tropospheric CH₄. For δ^2 H–CH₄ this was achieved by spiking fossil CH₄ gases with pure CH₃D to yield "Martha-1", "Martha-2" and "Mike-1". Mike-1 was then mixed with a fossil CH₄ gas to produce "Mike-2" while Martha-1 was spiked with pure CH₃D to produce "Martha-2". Martha-1 and Mike-1 were thereby transitional CH₄ mixtures.

We calibrated "Megan" and "Merlin" for δ^2 H–CH₄ and δ^{13} C–CH₄ as primary CH₄ gases (Fig. 1) against CRMs and WSs to the VSMOW and VPDB isotope scales respectively. Applied WSs are identical or similar in chemical composition to available CRMs in most cases (Table 2). All secondary CH₄ gases were calibrated against the primary CH₄ gases and are therefore of lower hierarchy level in the calibration scheme (Fig. 1). Megan was used as primary CH₄ gas for all initial experiments and our first calibrations of secondary CH₄ gases, until it was accidentally vented to ambient in March 2015. In order to compensate for the loss, we calibrated Merlin against CRMs and WSs as primary CH₄ in replacement of Megan.

2.2 Referencing pure CH₄ for δ^2 H against VSMOW/SLAP and against other pure CH₄ gases

We use a high-temperature conversion elemental analyser (TC/EA) coupled to an isotope ratio mass spectrometer (IRMS; Delta Plus XL, Thermo Finnigan, Bremen, Germany) via an open split (ConFlo III, Thermo Finnigan, Bremen, Germany). The system at MPI-BGC is operated for δ^{2} H–H₂O and δ^{18} O–H₂O analysis with high precision and negligible systematic errors since more than a decade (Gehre et al., 2004; Brand et al., 2009a) and is depicted in Fig. 2. Because TC/EA–IRMS systems are also used for δ^2 H analysis in hydrocarbons (e.g. Hilkert et al., 1999; Schimmelmann et al., 2016), this method is particularly suitable to calibrate δ^2 H–CH₄ against reference H₂O (CRM and WS, Table 2). Therefore, we inject CH₄ and H₂O samples through an externally heated septum (kept at 130 °C) into the glassy carbon reactor of the TC/EA-IRMS (kept at 1450 °C), where both species are converted to H_2 (+ carbon or CO). A helium carrier gas transports the sample gases from the hightemperature reactor through a gas chromatographic (GC) column (1/4 in. OD, 60 cm length, 5-Å zeolite, 75 °C) and into the open split. CH₄ injections are made with a two position 10-port valve (VICI, USA), which is configured as depicted in Fig. 2. A helium stream of 15 mL min⁻¹ carries CH₄ sam-

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	Cylinder				
Gas name	volume (L)	Pressure (bar)	Function in study	CH ₄ source	Gas supplier
Megan	10	_	first primary CH ₄ (lost)	fossil CH ₄	Air Liquide, Germany
Merlin	10	190	second primary CH ₄ (replacement of Megan)	fossil CH ₄	Air Liquide, Germany
Mike-1	_	_	secondary CH ₄	MPI mixture	MPI-BGC
Mike-2	5	45	secondary CH ₄	MPI mixture	MPI-BGC
Merkur	2	100	secondary CH ₄	fossil CH ₄	Messer Griesheim, Germany
Merida	10	175	secondary CH ₄	unknown	Westfalen AG, Germany
Martha-1	-	_	secondary CH ₄	MPI mixture	MPI-BGC
Martha-2	10	165	secondary CH ₄	MPI mixture	MPI-BGC
Minion	3	150	secondary CH ₄	unknown	Messer Griesheim, Germany
Melly	50	193	secondary CH ₄	unknown	Westfalen AG, Germany
δ^2 H-spike gas	0.4	2.5	CH ₃ D spiking gas		Campro Scientific, Germany
Fossil	30	2	secondary CH ₄ and comparison	fossil CH ₄	Air Liquide, Denmark
Biogenic	30	2	secondary CH ₄ and comparison	biogas plant	Biogas Plant, Germany
synthetic air	50	200	Synthetic air matrix		Linde, Germany
Krypton	2	200	synthetic air matrix		Westfalen AG, Germany
Carina-1	50	200	working standard and scale comparison	Jena air	MPI-BGC
Carina-2	50	200	working standard and scale comparison	Jena air	MPI-BGC

Table 1. Gases used for this study. Note that Mike-1 and Martha-1 were transitional CH₄ mixtures and do not exist anymore.

Table 2. Measurement standards used in this study. "CRM" and "WS" identify certified reference material and in-house working standards respectively. The uncertainties of the δ^2 H and δ^{13} C data from MPI-BGC correspond to the 95 % confidence limit of the error of the mean. We include the uncertainty estimate that the IAEA recently suggested for LSVEC. Publications and additional comments related to the standards are listed in the last column.

Name	Material	CRM/WS	$\delta^2 \mathrm{H} [\%]$	δ ¹³ C [‰]	Source	Reference/comment
VSMOW2	H ₂ O	CRM	0 ± 0.3	-	IAEA	Gröning et al. (2007), Brand et al. (2014)
SLAP2	H ₂ O	CRM	-427.5 ± 0.3	_	IAEA	Gröning et al. (2007), Brand et al. (2014)
GISP	H ₂ O	CRM	-189.7 ± 0.9		IAEA	Brand et al. (2014)
NBS 19	CaCO ₃	CRM		$+1.95\pm0.00$	IAEA	Brand et al. (2014), exhausted
LSVEC	Li ₂ CO ₃	CRM	-	-46.6 ± 0.15	IAEA	Coplen et al. (2006b), Qi et al. (2016),
						Schimmelmann et al. (2016)
CO-9	BaCO ₃	CRM	-	-47.32 ± 0.06	IAEA	Coplen et al. (2006b)
RM 8563	CO_2	CRM		-41.59 ± 0.06		Coplen et al. (2006b), exhausted
WWW-J1	H_2O	WS	-67.0 ± 0.4		MPI-BGC	-
BGP-J1	H_2O	WS	-187.1 ± 0.6		MPI-BGC	-
MAR-J1	CaCO ₃	WS		$+1.96\pm0.01$	MPI-BGC	Brand et al. (2009b)
ALI-J3	Acetanilide	WS		-30.06 ± 0.05	MPI-BGC	-
Cecily	CO_2	WS		-3.84 ± 0.015	MPI-BGC	-
Carina-1	Jena air	WS	-82.7 ± 4.0	-47.61 ± 0.09	MPI-BGC	calibration (T. Röckmann, personal communication, 2013)
Carina-2	Jena air	WS	-85.5 ± 4.0	-47.62 ± 0.12	MPI-BGC	calibration (T. Röckmann, personal communication, 2013)

ples from the 1 mL sample loops into the TC/EA reactor. Typical CH₄ flow rates range between 2 and 3 mL min⁻¹. For the calibrations of primary CH₄ gases, the two sample loops are fed by the same CH₄ gas (connecting vent 1 and CH₄ port 2). The sample loops are fed by two different gases for the calibration of secondary against primary CH₄ gases (Table 1). While CH₄ gases are injected manually, H₂O is introduced via autosampler.

It is recommended to measure samples against standards with identical material-specific properties (PIT; Werner and Brand, 2001). Under such conditions, measurement artefacts are likely to cancel when, for example, H_2O samples are calibrated against H_2O standards. However, great care has to be taken when chemically identical or similar CRMs are not available so that sample and standard comprise materials with different chemical properties, which is the case when calibrating CH_4 against H_2O . Calibration errors may arise when only one or both materials are fractionated during analysis, where the latter is likely to occur with different fractionation factors.

We performed a range of experiments to test for systematic errors during H₂O and CH₄ analysis. (i) System memory occurs during the isotopic analysis of H₂O due to adhesion of H₂O onto internal surfaces. System memory is sufficiently minimised by repeated H₂O injections and rejection of the first sample in every sequence. Remaining memory effects are corrected for in the evaluation routine as shown by Gehre et al. (2004). System memory is not created by CH₄ injections but some δ^2 H–CH₄ analyses may be affected by desorption of H₂O, stemming from previous injections. (ii) We observe a systematic effect of the septum temperature on the resulting δ^2 H–H₂O and operate the system with a septum temperature of 130 °C, where δ^2 H–H₂O was found stable. δ^2 H–CH₄ analysis is not affected by septum temperature. (iii) We experimentally optimised the TC/EA reactor temperature and found highest H₂ yields, quantitative conversion and hence smallest isotopic fractionation at 1450 °C during both δ^2 H–H₂O and δ^2 H–CH₄ analysis. Appendix A describes these experiments in greater detail.

The introduction of H₂ samples into the ion source of an IRMS leads to the formation of H_3^+ ions that are registered on the HD⁺ detector, which is accounted for by the so called "H3-factor correction" (Friedman, 1953; Sessions et al., 2001). The H₃-factor correction is experimentally determined and assumed to be constant until re-determined. Determining the H₃-factor correction is part of the daily preparation routine at MPI-BGC and shows only minor variation with time. Theoretically, the H_3^+ formation could be dynamic during the experimental period with unknown variability. We matched the H₂ peak heights resulting from both CH₄ and H_2O injections around $5.5 \pm 0.5 V$ in order to minimise the impact of imperfect H3-factor correction. Peak widths ranged around 45 and 60 s for H₂O- and CH₄-derived H₂ peaks respectively. A typical chromatogram of the δ^2 H–CH₄ calibration including details on peak shape and background is shown in Fig. 3. The similarity between CH₄-derived and the H₂O-derived H₂ peaks allows the use of the standard integration software (ISODAT, Thermo Finnigan, Bremen, Germany).

Megan and Merlin (Table 1) were calibrated in three independent sequences during 3 days against the in-house working standards "WWW-J1" and "BGP-J1" with a wide δ^2 H range from -67.0 to -187.1% (Table 2). WWW-J1 and BGP-J1 are independently calibrated against international reference waters VSMOW2 and SLAP2 (Table 2). Other CH₄ gases were initially also measured against working standards (WWW-J1 and BGP-J1) but were finally calibrated against Megan or Merlin, which were co-analysed in the same measurement sequence in a one-point calibration.



Figure 3. Chromatograms of δ^2 H–CH₄ calibration sequences using TC/EA–IRMS with traces of m/z 2 and m/z 3 shown in black and blue respectively. The bottom panel shows an example of an entire calibration sequence which begins with three square-shaped peaks of pure H₂, followed by alternations of three to four H₂O- and three to four CH₄-derived H₂ peaks before the sequence ends with another three square-shaped peaks of pure H₂. The top left panel enlarges H₂ peaks from H₂O (peak no. 6–7) and CH₄ (peak no. 8–9) injections respectively. A zoom into baseline details of H₂O-derived peak no. 7 and CH₄-derived peak no. 8 is shown in the top right panel. Red lines indicate the sections used for peak integration (weak widths are 43 and 59 s for H₂O- and CH₄-derived H₂ peaks respectively) by the IRMS software.

2.3 Referencing pure CH₄ for δ^{13} C against LSVEC/MAR-J1 and against other pure CH₄ gases

We calibrated δ^{13} C–CH₄ in pure CH₄ gases after conversion to CO₂ using an elemental analyser (EA 1100, CE, Rodano, Italy) coupled to an IRMS (Delta Plus, Thermo Finnigan, Bremen, Germany) via open split (ConFlo III, Thermo Finnigan, Bremen, Germany). This system is routinely used for the analysis of ¹³C and ¹⁵N in samples with solid or liquid matrices (Werner et al., 1999; Brooks et al., 2003). We fitted a 1/16 in. tube of 70/30 % Cu/Ni alloy to the EA and used the previously described 10-port valve to inject the CH₄ samples into the EA with a 10 mL min⁻¹ helium flow (Fig. 4).

The plumbing of the system is designed so that gaseous CH₄ and solid CRMs/WSs are applied to the same location inside the combustion reactor of the EA. All samples are combusted at a reactor temperature of 1020 °C (Werner et al., 1999) and experience identical analytical treatment thereafter. Following the combustion, each sample passes through a reduction reactor filled with elemental copper, which is kept at 650 °C to remove excess O₂ and to reduce NO_x if present. The sample is dried by passing through a NafionTM membrane (Perma Pure LLC, Toms River, NJ, USA; not shown in



Figure 4. The 10-port valve for manual CH₄ injections is coupled to the EA–IRMS system through custom-made gas inlet into combustion (oxidation) unit for δ^{13} C–CH₄ calibration. The proportions of illustrated components are chosen to increase clarity.

Fig. 4) and a Mg(ClO₄)₂ trap before it enters the GC column (3 m, 1/4 in.; Porapak PQS, CE instruments) held at 80 °C. Thereafter, the sample enters the IRMS through the open split.

Measurement sequences to calibrate primary CH₄ gases to the VPDB isotope scale are created by alternating blocks of manual CH₄ injections and CRM/WS (Table 2) applications via autosampler. We applied one WS and one CRM (LSVEC) to calibrate the primary CH₄ gases in a two-point calibration. While MAR-J1 was used as WS in most experiments, ALI-J1 was used once, during a calibration of Merlin. Megan and Merlin were each calibrated on 3 different days to determine the external reproducibility of the δ^{13} C results. Chromatograms resulting from CH₄ and from carbonate analyses using EA-IRMS are displayed in Fig. 5 and show very similar peak shapes for CH₄ and carbonates. Typical m/z 44 amplitudes and peak widths were $\sim 7.4 \pm 0.2$ V and 101 ± 1 s for both materials respectively. We connected a primary CH₄ and a secondary CH₄ gas to the 10-port valve to calibrate the secondary CH₄ gases (Table 1) for δ^{13} C in a one-point calibration. All measurement results were corrected for scale compression based on the method suggested in Verkouteren and Klinedinst (2004), using an empirical, mass spectrometer specific correction factor of 1.0056.

2.4 Measurement uncertainty and error propagation

The fully propagated uncertainty for the primary CH₄ gases $(U_{pCH4-tot})$ is calculated as

$$U_{\rm pCH_4-tot} = \sqrt{u_{\rm CRM}^2 + u_{\rm WS}^2 + u_{\rm pCH_4}^2},$$
 (1)

where u_{CRM} , u_{WS} and u_{pCH4} indicate the uncertainty of the CRM, the applied working standards and the respective primary CH₄ gas respectively. Both u_{WS} and u_{pCH4} are calculated as the standard error of the mean of all measurements,



Figure 5. Chromatograms of δ^{13} C–CH₄ calibrations using EA–IRMS with traces for m/z 44, 45 and 46 in green, brown and black respectively. Bottom panels show complete chromatograms of CH₄ and Li₂CO₃ analyses while the two top panels zoom into the baseline of the traces. The first three square-shaped peaks stem from injections of a pure CO₂ WS while the more Gaussian-shaped peaks result from CH₄- and Li₂CO₃-derived CO₂ analysis. The two red lines indicate the sections that the IRMS software uses for peak integration (CO₂ peak widths are 101 and 100 s for CH₄ and Li₂CO₃ analysis respectively).

multiplied by t, Student's factor for a 95 % confidence limit to account for the limited number of measurements.

The uncertainty for the secondary CH_4 gases ($U_{sCH4-tot}$) is then calculated as

$$U_{\rm sCH_4-tot} = \sqrt{U_{\rm pCH_4-tot}^2 + u_{\rm sCH_4}^2},$$
 (2)

where u_{sCH4} is the standard error of the mean of all measurements of the respective secondary CH₄ gas, multiplied by *t*, Student's factor for a 95 % confidence limit. Therefore, $U_{pCH4-tot}$ and $U_{sCH4-tot}$ indicate the fully propagated uncertainty onto the VPDB or VSMOW isotope scales, representing the traceability chain.

Note that the isotopic composition of LSVEC (Table 2) was recently found to show significant variability, most likely due to adhesion of H₂O and reaction with air-CO₂ (e.g. Oi et al., 2016; Schimmelmann et al., 2016). Until this problem is solved, the IAEA, one of the providers of LSVEC, advised to increase the uncertainty of LSVEC, which was hitherto assigned to 0.00 %. We follow the recommendation by S. Assonov (Sergey Assonov, IAEA, personal communication, 2016) and Schimmelmann et al. (2016) and adopt an uncertainty of 0.15 % for the δ^{13} C of LSVEC. Note that the new 0.15 % uncertainty of LSVEC represents the largest single contributor to the total uncertainty budget in our δ^{13} C calibrations. As a consequence we present the combined uncertainty of the full traceability chain in two versions, the first being the hitherto adopted method using an uncertainty of 0.00% for LSVEC and the second being the method with uncertainty for LSVEC of 0.15 %.

2.5 Producing synthetic CH₄-in-air standards from pure CH₄ and CH₄-free air (JRAS-M16)

The MPI-BGC operates an analytical system (named ARAMIS) to dilute pure CO₂ with CO₂-free air to atmospheric CO₂ mole fraction without isotopic fractionation (Ghosh et al., 2005). We use ARAMIS to dilute an aliquot of primary or secondary CH₄ with CH₄-free air to atmospheric CH₄ mole fractions ($\sim 2 \text{ ppm}$) in 5 L glass flasks with a final filling pressure of 1.8 bar absolute. The produced synthetic CH₄-in-air standards represent the JRAS-M16 reference gases. The CH₄-free matrix air has been target-mixed from ultra-pure constituents and contains N2, O2, N2O and Kr at atmospheric levels, so that the composition of the produced CH₄-in-air standards is as close to ambient air as possible. Krypton was added to this matrix air to account for the measurement artefact during GC-IRMS analysis of CH4 for δ^{13} C (Schmitt et al., 2013). A blank analysis of the CH₄free air yielded a maximum CH₄ blank of 0.5 ppb. Because such a CH₄ blank is too small for accurate isotopic analysis on our atmospheric system (Sect. 2.6 and Brand et al., 2016), we choose a mass-balance calculation to determine the maximum blank effect in the synthetic CH₄-in-air standards. Let us assume δ^2 H–CH₄ and δ^{13} C–CH₄ values for the CH_4 blank of -150 and -40%, respectively, that are typical for fossil CH₄. Let us now mix this blank with target CH₄ comprising the most depleted δ^2 H–CH₄ and δ^{13} C–CH₄ values we find in our CH₄ gas suite with δ^2 H of -320%and δ^{13} C of -70%. The mixing ratio of blank:target CH₄ is 1 : 4000, which reflects the ratio within a synthetic CH₄-inair mixture with 2 ppm CH₄. The maximum blank contribution in this extreme scenario would be 0.04 and 0.007 % for δ^2 H–CH₄ and δ^{13} C–CH₄, respectively, which is negligible in both cases.

2.6 Analytical systems to measure the isotopic composition of CH₄ in air at IMAU and MPI-BGC

IMAU: Brass and Röckmann (2010) and Sapart et al. (2011) described the system for the analysis of both δ^2 H and δ^{13} C in atmospheric methane at IMAU. CH₄ is separated from the other air components by cryogenic traps and gas chromatography before it is converted by either oxidation or pyrolysis for IRMS analysis on CO₂ or H₂ respectively. For δ^{13} C an additional GC column (PoraPlotQ, 12.5 m, 0.32 mm ID, Agilent, the Netherlands) was added between the Nafion drying unit and the open split interface in order to remove the interferences from Kr (Schmitt et al., 2013).

MPI-BGC: a new system to measure δ^2 H–CH₄ and δ^{13} C–CH₄ in-air samples was recently developed at the MPI-BGC and is described in greater detail in Brand et al. (2016). The system at MPI-BGC is referred to as *i*SAAC, in abbreviation for integrated System for Analysis of Atmospheric Constituents. iSAAC consists of a 16-port sample carousel to take two consecutive 100 mL aliquots of air from a glass flask or high-pressure cylinders for parallel analysis of δ^2 H–CH₄ and δ^{13} C–CH₄, respectively, by continuous-flow GC–IRMS. The two air samples are routed through two identical but independent pre-concentration lines, one for the analysis of δ^2 H–CH₄ and one for δ^{13} C–CH₄. In each line, CH₄ is cryogenically separated from the main air constituents in a Hayesep D-filled trap at -130 °C and cryo-focussed in a further Hayesep D-filled trap at -110 °C. Each of the two analytical lines is equipped with its own cooling compressor to avoid the use of cryogenic liquids. The separated and cryofocussed CH₄ sample is released into a GC column from where it is routed either through a pyrolysis furnace (kept at 1400 °C) to convert the CH₄ sample to H₂ for δ^2 H–CH₄ analysis or through a combustion furnace (kept at 1000 °C) to convert the CH₄ sample to CO₂ for δ^{13} C–CH₄ analysis. A post-combustion GC column separates the CH₄-derived CO₂ from Kr (Schmitt et al., 2013). CH₄-derived H₂ and CO₂ samples are introduced via open splits into dedicated IRMS instruments, one each for δ^2 H–CH₄ and δ^{13} C–CH₄ analysis. iSAAC has been operational since 2012 to measure air samples with a precision of 1.0 and 0.12 % for δ^2 H–CH₄ and δ^{13} C–CH₄ respectively. The precision is determined by the performance chart method (Werner and Brand, 2001), determined by the standard deviation (1σ) of all quality control standard measurements, which has been analysed once in every measurement sequence (Brand et al., 2016). The reproducibility of δ^{13} C–CH₄ analyses ranges around 0.06 % over the course of 1 day. All measurements on *i*SAAC so far have been allocated to the VPDB and VSMOW scales using an inhouse WS that was calibrated against "Carina-1" (Table 1).

2.7 Histories to anchor δ^2 H–CH₄ and δ^{13} C–CH₄ to the VSMOW and VPDB scales at IMAU and MPI-BGC

It is the intention of all laboratories analysing δ^2 H–CH₄ and δ^{13} C–CH₄ to reference their samples relative to the VSMOW and VPDB scales respectively. However, possible accuracy errors in the laboratory specific scale anchors often result in inter-laboratory offsets. In order to retrace the potential for calibration offsets between IMAU, MPI-BGC and JRAS-M16, we describe the history of the scale anchors for each laboratory.

IMAU: the calibration strategy at IMAU, including traceability chain and long-term control, is different for δ^2 H–CH₄ and δ^{13} C–CH₄ (Brass and Röckmann, 2010). (i) Three synthetic gas mixtures with CH₄ mole fractions of ~ 9000 ppm were calibrated for δ^2 H–CH₄ at the Max Planck Institute for Chemistry (MPI-C) in Mainz, Germany, using a tunable diode laser absorption spectrometer (TDLAS) technique. The TDLAS is described by Bergamaschi et al. (1994) with a measurement precision for δ^2 H–CH₄ of 5.1 ‰ and an accuracy estimate of similar magnitude. The accuracy estimate is based on a comparison with the calibrations to the VSMOW scale by Dumke et al. (1989), which marks the origin of the isotope scale anchor for δ^2 H–CH₄ at IMAU. Aliquots of the gases from Bergamaschi et al. (1994) were diluted with synthetic CH₄-free air at IMAU to yield reference gases ("Cal1", "Cal2", "Cal3") with the δ^2 H–CH₄ values initially assigned at MPI-C and atmospheric CH₄ levels. Improved measurement precision and inter-laboratory comparisons lead to a δ^2 H–CH₄ refinement in Cal1, Cal2 and Cal3 with recent values of +21.1, -19.0 and -164.9% respectively. Cal1, Cal2 and Cal3 represent the primary reference gases for δ^2 H–CH₄ at IMAU and were used to calibrate the δ^2 H–CH₄ in the working standard ("SiL") to the VSMOW scale. While Cal2 and Cal3 have become exhausted, Cal1 is still used in regular checks of the calibration scale, together with a set of firn air samples (see ii) that are used for $\delta^{13}C$ calibration. (ii) IMAU's working gas SiL has also been calibrated for δ^{13} C–CH₄. This was achieved by co-analysing SiL with a suite of Antarctic firn gas samples, where the δ^{13} C–CH₄ of the latter had been determined by two laboratories (MPI-C and the Laboratoire de Géologie et Géophysique de l'Environnement (LGGE), Grenoble, France), using two different techniques (Bräunlich et al., 2001). The δ^{13} C–CH₄ scale anchors at LGGE and MPI-C are calibrated at MPI-C against a pure CO₂ WS, which itself has been calibrated against NBS 19 (Bergamaschi et al., 2000), which represents the ultimate link to the VPDB scale for the scale anchor at IMAU. Using that method, the suite of firn gas samples was treated as a set of working standards to calibrate SiL to the VPDB scale by propagation from MPI-C and LGGE to IMAU. It is important to note that Brass and Röckmann (2010) highlighted that the firn gas itself is a set of samples and not to be taken for a set of calibration standards. The calibration strategy was revised during 2013 to account for the Kr interference (Schmitt et al., 2013).

MPI-BGC: all measurements on *i*SAAC use a natural air WS that was calibrated against Carina-1 at MPI-BGC. Carina-1 and Carina-2 are natural air samples that were calibrated for δ^2 H–CH₄ and δ^{13} C–CH₄ at IMAU (Table 2), using the analytical setup described by Brass and Röckmann (2010) and Sapart et al. (2011). While the calibration results of Carina-1 and Carina-2 from IMAU show excellent agreement in CH₄ mole fractions (both 1910 ppb), in δ^{13} C–CH₄ (within 0.01%), their δ^{2} H–CH₄ values differed by 2.8 % (Table 2). Because both Carina cylinders were filled at the MPI-BGC with Jena air on the same day within a short period of time during stable meteorological conditions, and because their δ^{13} C–CH₄ and CH₄ mole fractions are in excellent agreement, a true difference in δ^2 H–CH₄ between the two Carina cylinders seems unlikely. The magnitude of the δ^2 H–CH₄ offset was smaller than the former δ^2 H–CH₄ measurement precision at IMAU of ±4 % (Brass and Röckmann, 2010) and was accepted as "agreement within measurement uncertainty" at the time. It is important to note that Carina-1 and Carina-2 were each calibrated on different days and in separate measurement sequences, which does not enable a direct comparison of the two gases. Therefore, a systematic calibration error in one of the two Carina gases is possible. In contrast, the superior measurement precision of *i*SAAC for δ^2 H–CH₄ of 1.0% can resolve a true δ^2 H–CH₄ difference of 2.8%. However, both Carina-1 and Carina-2 appear indistinguishable in δ^2 H–CH₄ on *i*SAAC, as determined during several direct comparisons in independent measurement sequences. Therefore, the δ^2 H–CH₄ offset between Carina-1 and Carina-2 must be due to an artefact of the calibration at IMAU. Our experiments at MPI-BGC indicate that the calibration of Carina-1 is indeed flawed. The choice to use Carina-1 as scale anchor for all iSAAC measurements at MPI-BGC was made arbitrarily, before it was known that it's calibration was impacted by an artefact. In hindsight, Carina-2 would have been a better choice as VSMOW scale anchor for δ^2 H–CH₄ at MPI-BGC. This calibration offset will be furthermore addressed a future comparison with IMAU, where a new system has been developed with an improved precision in for δ^2 H–CH₄ (Röckmann et al., 2016). All iSAAC measurements are anchored to the VSMOW and VPDB isotope scales based on the described scale propagation from IMAU to MPI-BGC, until JRAS-M16 is established as new m-RM for δ^2 H–CH₄ and δ^{13} C–CH₄ in air.

2.8 Comparison of the existing isotope scales at MPI-BGC with new, synthetic CH₄-in-air standards

The synthetic CH₄-in-air standards produced in this study (Sect. 2.5) were analysed at MPI-BGC using *i*SAAC (Sect. 2.6). In that, the synthetic CH₄-in-air standards

are treated as unknown samples and their δ^2 H–CH₄ and δ^{13} C–CH₄ values are determined using Carina-1 as scale anchor (Sect. 2.7). We calculate the isotopic difference ($\delta_{iSAAC} - \delta_{pure}$) between the measurements on *i*SAAC and the calibrations of the pure CH₄ gases (Sects. 2.2 and 2.3), which indicates the correction to anchor the measurements at MPI-BGC to JRAS-M16.

2.9 Comparison between CIC and MPI-BGC

Two CH₄ gases, Biogenic and Fossil, were previously calibrated at CIC by Sperlich et al. (2012), who analysed the CH₄-derived CO₂ for δ^{13} C–CH₄ by dual-inlet IRMS and the CH₄-derived H₂O for δ^2 H–CH₄ by either cavity ringdown spectroscopy (CRDS) or TC/EA-IRMS. Sperlich et al. (2012) presented the data with the measurement reproducibility, calculated as the pooled standard deviation of the measurements. Therefore, their uncertainty does not include the uncertainties of the full traceability chain. Furthermore, a statistical provision that accounts for the small number of measurements has not been made by Sperlich et al., (2012). This imposes a hurdle in the comparison with data from MPI-BGC. Therefore, we revise the uncertainty of the CIC data and calculate the full traceability chain as described in Sect. 2.4. Furthermore, all δ^{13} C measurements from CIC are affected by a small offset of RM 8563 that has been reported by Coplen et al. (2006b) and are therefore shifted by 0.03 % towards more depleted δ^{13} C values. Moreover, the δ^{13} C data presented in Sperlich et al. (2012) have not been corrected for scale compression. We are able to correct all CIC data for this effect, because the scale compression factor of the instrument at CIC has been determined (1.0025) at the time the study of Sperlich et al. (2012) was published. Applying the scale compression correction shifts the δ^{13} C–CH₄ of Fossil and Biogenic by 0.01 and 0.05 % towards more depleted δ^{13} C values respectively. The revised data and uncertainties from CIC and the results from MPI-BGC for Biogenic and Fossil are shown in Table 4 for δ^{13} C–CH₄ and in Table 5 for δ^2 H–CH₄.

We perform two comparisons between CIC and MPI-BGC. (i) The calibration results for Fossil and Biogenic from CIC as published in Sperlich et al. (2012) are compared to the calibrations at MPI-BGC using the methods to calibrate pure CH₄ gases for δ^2 H–CH₄ and δ^{13} C–CH₄ as described in Sects. 2.2 and 2.3. (ii) We performed new combustion experiments at CIC using Fossil and Biogenic and analysed the resulting CO₂ for δ^{13} C at both CIC and MPI-BGC. These combustion experiments were made in 2012 but after the publication of Sperlich et al. (2012). Therefore, these experiments provide new data to evaluate the method at CIC. Following the δ^{13} C analyses at CIC, the remaining CO₂ gases were cryogenically transferred and flame sealed in glass ampules for δ^{13} C analysis at MPI-BGC. The δ^{13} C analyses at MPI-BGC were made on "Cora", a MAT 252 dual-inlet IRMS (Thermo Finnigan, Bremen, Germany) that is used for δ^{13} C and δ^{18} O analysis of CO₂ in air or pure CO₂ gases (Brand et al., 2009b). Unfortunately, the comparison based on the new combustion experiments made at CIC could not include δ^2 H–CH₄ because the system was not capable to process CH₄ samples large enough to provide sufficient amounts of H₂O.

We use the indices $_{CIC-old}$ for experiments made at CIC and published by Sperlich et al. (2012) and $_{CIC-new}$ for the new combustion experiments at CIC. We use the index $_{MPI-BGC}$ for the analysis at MPI-BGC of CO₂ samples that were combusted at CIC and $_{MPI-BGC}$ for the calibrations of the two CH₄ gases from CIC using the analytical methods at MPI-BGC presented above (Sect. 2.2 and 2.3).

3 Results

3.1 Results for primary CH₄ gas calibrations on the international VSMOW and VPDB isotope scales

We performed 214 repetitive calibration measurements for Megan and Merlin for δ^2 H–CH₄ and δ^{13} C–CH₄; the results are given in Table 3. Megan and Merlin have δ^2 H–CH₄ values of –168.1±0.7% and –165.7±0.7%, respectively, and δ^{13} C–CH₄ values of –40.76±0.04 and –39.06±0.02%, respectively. Both the δ^2 H–CH₄ and δ^{13} C–CH₄ values are typical for fossil CH₄ (e.g. Quay et al., 1999; Mikaloff Fletcher et al., 2004). The δ^{13} C–CH₄ uncertainty in Megan and Merlin increases to 0.16 and 0.15%, respectively, when the suggested uncertainty of 0.15% for LSVEC is taken into account in the traceability chain (Qi et al., 2016; Schimmelmann et al., 2016). However, we will use the uncertainty budget without the new uncertainty for LSVEC for the evaluation of internal results.

3.2 Results for secondary CH₄ gas calibrations against primary CH₄ gases

We made a total of 260 calibration measurements for the secondary CH₄ gases for δ^2 H–CH₄ and δ^{13} C–CH₄. Altogether, the secondary CH₄ gases cover a large range in δ^2 H (-320 to +36%) and δ^{13} C (-70 to -39%), where the former was achieved by spiking some of the gases with pure CH₃D. The results for secondary CH₄ gas calibrations are shown in Table 3, including the uncertainties of the full traceability chain. We found typical uncertainties on the order of 0.8% for δ^2 H–CH₄ calibrations and on the order of 0.07 and 0.17% for δ^{13} C–CH₄ calibrations, where the latter includes the uncertainty of 0.15% in LSVEC.

3.3 Results from the comparison between CIC and MPI-BGC

Our comparison results for δ^{13} C–CH₄ show overall agreement within the uncertainties of the traceability chains (Table 4). The δ^{13} C results from the previous and the new com-

Table 3. Results of CH₄ isotope calibrations. Gas names as used in main text and their function as primary or secondary CH₄ are shown in column 1 and 2 respectively. All uncertainty estimates include the full traceability chain (Sect. 2.4). Note that we provide uncertainty estimates for δ^{13} C–CH₄ without and with the uncertainty of 0.15% in LSVEC in column 6 and 7 respectively. Martha-1 and Mike-1 were intermittent gases and used to produce Martha-2 and Mike-2.

Gas name	Function	$n~(\delta^2 \mathrm{H})$	δ ² H–CH ₄ [% ₀]	$n (\delta^{13}\mathrm{C})$	δ^{13} C–CH ₄ [‰] $u_{LSVEC} = \pm 0.00\%$	δ^{13} C-CH ₄ [‰] $u_{LSVEC} = \pm 0.15 \%$
Megan	primary	116	-168.1 ± 0.7	15	-40.76 ± 0.04	-40.76 ± 0.16
Merlin	primary	51	-165.7 ± 0.7	32	-39.06 ± 0.02	-39.06 ± 0.15
Martha-1	secondary	15	-176.6 ± 0.8	10	-48.84 ± 0.07	-48.84 ± 0.17
Martha-2	secondary	9	$+36.2 \pm 1.0$	19	-48.92 ± 0.06	-48.92 ± 0.16
Mike-1	secondary	12	$+44.5\pm0.9$	8	-40.79 ± 0.09	-40.79 ± 0.17
Mike-2	secondary	15	-80.3 ± 0.5	13	-42.76 ± 0.05	-42.76 ± 0.16
Merida	secondary	12	-171.7 ± 0.9	13	-60.39 ± 0.09	-60.39 ± 0.18
Melly	secondary	19	-177.5 ± 0.7	15	-70.04 ± 0.07	-70.04 ± 0.17
Minion	secondary	12	-182.7 ± 0.8	15	-58.19 ± 0.05	-58.19 ± 0.16
Merkur	secondary	15	-195.8 ± 0.9	19	-43.05 ± 0.04	-43.03 ± 0.16
Fossil	secondary	15	-171.9 ± 0.9	16	-39.71 ± 0.08	-39.71 ± 0.17
Biogenic	secondary	25	-319.8 ± 0.8	10	-56.60 ± 0.07	-56.60 ± 0.17

bustion experiments measured at CIC differ by -0.03 and -0.06% ($\delta_{CIC-new} - \delta_{CIC-old}$) for Fossil and Biogenic, respectively, which is within the uncertainty of the full traceability chain and furthermore within the system reproducibility as stated in Sperlich et al. (2012). The δ^{13} C differences between the results from the new combustion experiments measured at MPI-BGC* and at CIC ($\delta_{MPI-BGC*} - \delta_{CIC-new}$) are 0.10% for Fossil and 0.11% for Biogenic, respectively, and agree well within the combined uncertainty of both methods. Table 4 shows even better agreement for $\delta_{MPI-BGC*} - \delta_{CIC-old}$. Altogether, the comparisons highlight the reproducibility of δ^{13} C measurements in the combustion-derived CO₂ at both laboratories.

When comparing the δ^{13} C results from the new calibrations at MPI-BGC to the results based on combustion experiments at CIC, the results from MPI-BGC appear slightly more depleted in δ^{13} C for both Fossil and Biogenic in all comparisons (Table 4). We find the smallest δ^{13} C differences between $\delta_{MPI-BGC}$ and $\delta_{CIC-new}$, accounting for -0.08 and -0.09% for Fossil and Biogenic respectively. The respective differences increase to -0.18 and -0.20% between $\delta_{MPI-BGC}$ and $\delta_{MPI-BGC^*}$. It is important to note that only the difference found in Biogenic between $\delta_{MPI-BGC}$ and $\delta_{MPI-BGC^*}$ is outside of the sum of the uncertainties (Table 4). In contrast, we find excellent agreement in all comparisons when the uncertainty of 0.15% in LSVEC is taken into account. Table 4 also shows excellent agreement in the determination of the differences between Fossil and Biogenic in all δ^{13} C measurements, which is an important quantity for the evaluation of scale compression.

Comparing the results for δ^2 H–CH₄ between CIC and MPI-BGC shows overall agreement (Table 5). The differences we find in the δ^2 H–CH₄ calibrations between Sperlich

et al. (2012) and MPI-BGC ($\delta_{MPI-BGC} - \delta_{CIC-old}$) are -1.8and -2.4% for Fossil and Biogenic respectively. Albeit it is slightly larger than the sum of the uncertainties of the measurements at CIC and MPI-BGC for Biogenic, the difference in Fossil is just within the uncertainties of the two methods. Note that the isotopic difference Fossil–Biogenic is homogenously resolved with 147.9% at MPI-BGC and 147.3% at CIC respectively.

3.4 Results of δ^2 H–CH₄ and δ^{13} C–CH₄ measurements in synthetic CH₄-in-air standards to determine compatibility between the propagated isotope scale from IMAU and JRAS-M16 at MPI-BGC

The isotopic difference $(\delta_{iSAAC} - \delta_{pure})$ is shown in Table 6 and indicates the offset between the scale that was propagated from IMAU to MPI-BGC (Sect. 2.8) and the new synthetic CH₄-in-air standards (JRAS-M16), assuming no isotope fractionation during the dilution process. Our experiments show excellent agreement for δ^{13} C–CH₄ with an average difference of $+0.03 \pm 0.10\%$, thus confirming that the propagated scale from IMAU was already very close to the newly determined scale anchor for δ^{13} C–CH₄. For unknown reasons, the δ^{13} C–CH₄ measurements of Melly, Fossil and Biogenic show a larger discrepancy between the two methods. Because the discrepancy for Biogenic exceeds the measurement uncertainty by a factor of 3, we have excluded this result from the determination of the laboratory offset. The values for Melly and Fossil are within two uncertainties and are therefore included. For δ^2 H–CH₄ a systematic offset of $+4.2 \pm 1.2\%$ is found, confirming that the calibration of Carina-1 and hence the scale propagation from IMAU to MPI-BGC is flawed by an artefact. Obviously,

Table 4. Results for comparison in δ^{13} C between CIC and MPI-BGC. Indices of the header are explained in Sect. 2.9 of the main text. The CIC data are corrected for the offset in RM 8563 (Coplen et al., 2006b) and scale compression effects. They are furthermore presented with revised uncertainties to include the full traceability chain (Sect. 2.4). The uncertainties of the full traceability chains with the recently suggested uncertainty in LSVEC of 0.15% are shown in brackets. The $\delta^{13}C_{MPI-BGC^*}$ measurements used a system that is virtually unaffected by scale compression (Ghosh et al., 2005) and a WS calibration that is based on NBS 19 as the only CRM; therefore, the $\delta^{13}C_{MPI-BGC^*}$ data do not suffer from the uncertainty in LSVEC. The difference Fossil–Biogenic can be used to compare scale compression effects between the respective methods.

Gas name	$\delta^{13}C_{CIC-old}$ [%]	$\delta^{13}C_{CIC-new}$ [%]	$\delta^{13}C_{MPI-BGC^*}$ [%]	$\delta^{13}C_{MPI-BGC}$ [%]
Fossil	$-39.60 \pm 0.07 \; (0.17)$	$-39.63 \pm 0.14 \; (0.20)$	-39.53 ± 0.11	$-39.71 \pm 0.08 \ (0.17)$
Biogenic	$-56.45 \pm 0.10 \; (0.18)$	$-56.51 \pm 0.08 \ (0.17)$	-56.40 ± 0.04	$-56.60 \pm 0.07 \ (0.17)$
Fossil-Biogenic	16.85	16.88	16.87	16.89

Carina-2 would have been a closer choice as scale anchor for δ^2 H–CH₄ (Table 2).

4 Discussion

4.1 Discussion on the experimental artefact elimination during δ^2 H–CH₄ and δ^{13} C–CH₄ calibrations in primary and secondary CH₄ gases

We present $\delta^2 H$ and $\delta^{13}C$ calibrations in pure CH₄ gases against CRMs, WSs and other CH₄ gases. Samples and reference materials were always analysed in the same analytical systems, thereby complying with the PIT as much as possible. The only limitation of the PIT is due to the chemical difference between unknown samples (CH₄) and the known reference materials (carbonates, H₂O) used for anchoring the CH₄ gases to the respective isotope scales. In order to calibrate the primary CH₄ gases accurately, we need to exclude or eliminate material- and method-specific errors (IAEA, 2003), which we discuss in the following.

Quantitative oxidation of CH₄ during δ^{13} C–CH₄ analysis requires high reaction temperatures (e.g. Dumke et al., 1989). A major complication during δ^{13} C–CH₄ analysis arises when oxidation yields are significantly lower than 100 % (Merritt et al., 1995; Fig. 4 in Sperlich et al., 2012). CH₄ is a potent source of protonation in the IRMS ion source (Anicich, 1993). Introducing unconverted CH₄ together with the CH₄-derived CO₂ sample into the IRMS results in the formation of CO₂H⁺ in the ion source, which produces an isobaric interference on the m/z 45 trace, where the δ^{13} C signal is measured. This artefact can be prevented when CO₂ and CH₄ are separated after the oxidation, which we achieve with the post-combustion chromatographic column in both the EA-IRMS system (Sect. 2.3) and iSAAC (Sect. 2.6). Note how this effect would cause an accuracy shift towards more enriched δ^{13} C–CH₄ values predominantly during primary CH₄ gas calibrations, because CH₄ samples would be affected by CO_2H^+ formation in the ion source while the analysis of the used CRMs would not.

We carefully checked the completeness of CH₄ conversion (EA-IRMS and TC/EA-IRMS) by monitoring for residual CH₄ with the IRMS instruments. In the ion source, CH₄ molecules are subject to fragmentation and re-combination processes, resulting in CH₄-typical mass spectra during mass abundance scans in the IRMS (Brunnée and Voshage, 1964). The strongest CH₄-specific signal occurs on the m/z 15 trace (CH_{3}^{+}) , which makes the m/z 15 signal a good indicator for incomplete CH₄ conversion (Sperlich et al., 2012). The CH_4^+ signal at m/z 16 is not suitable for CH_4 quantification due to the interference with the O^+ signal from CO_2^+ fragmentation. We tune the m/z 44 collector of the IRMS to monitor the m/z 15 trace during the analysis of a CH₄ sample and find an amplitude of 0.12 mV. From Sperlich et al. (2012) we estimate that about 40% of the total CH₄ signal in a mass abundance scan is recorded on m/z 15. The total CH₄ signal in the mass abundance scan would therefore amount to $\sim 0.3 \,\mathrm{mV}$, which we can compare to the \sim 7000 mV on m/z 44 from a typical CH₄ injection into the EA–IRMS (e.g. Fig. 5). This approximation suggests a CH_4 oxidation efficiency of >99.9 %. An analogue experiment on the TC/EA-IRMS system (Sect. 2.2) shows a conversion efficiency of CH_4 of >99.9 % as well. Because the ionisation energy of CH₄ is comparable to that of both CO₂ and H₂, we can ignore this effect in the above determinations. Therefore, we conclude that the CH₄ conversion at MPI-BGC is complete and that we can rule out incomplete conversion as source for measurement errors.

It has been demonstrated that the introduction of carbonates into the high-temperature oxidation furnace of the EA–IRMS yields a high CO₂ conversion rate and δ^{13} C results of high precision and accuracy (Coplen et al., 2006b). In order to test for the completeness of carbonate digestion, we added tungsten trioxide (WO₃) to some of the carbonate samples during weighing (about 1 : 1 by weight). The goal of this experiment is to increase the instantaneous reaction temperature and to provide additional oxygen during the liberation of CO₂ from different carbonates. While the addition of WO₃ had no effect on the analysis of CaCO₃ and Li₂CO₃, it improved the peak shape during BaCO₃ analysis (Table 2).

Table 5. Comparison of δ^2 H results between CIC and MPI-BGC. Indices of the header are explained in Sect. 2.9 of the main text. The uncertainty of all data includes the full traceability chain (Sect. 2.4), which includes revised uncertainties of the CIC data (Sect. 2.9). The difference Fossil – Biogenic allows us to compare scale compression effects between both methods.

Gas name	$\delta^2 H_{CIC-old}$ [%]	$\delta^2 H_{MPI-BGC}$ [%]
Fossil	-170.1 ± 0.9	-171.9 ± 0.9
Biogenic	-317.4 ± 0.9	-319.8 ± 0.8
Fossil-Biogenic	147.3	147.9

However, it did not impact on its δ^{13} C. We conclude that the carbonate digestion is not limited by either temperature or oxygen availability and omitted the addition of WO₃ in further reactions. Note that the accurate analysis of carbonates is critical for accurate CH₄ calibrations, even if CH₄ injections themselves are not compromised.

A considerable advantage of the conversion of carbonates in the high-temperature oxidation furnace of the EA–IRMS over other methods (e.g. acid reaction) is that the oxygen isotopic composition is homogenised for all samples. This balances the ¹⁷O correction, which accounts for the isobaric interference between δ^{13} C–CO₂ and δ^{17} O–CO₂ on m/z 45. The ¹⁷O correction is statistically dependent on the δ^{18} O–CO₂ of each individual sample. Hence, any uncertainty arising from the ¹⁷O correction during the calculation of δ^{13} C values from m/z 45 ion currents tends to cancel out. The applied ¹⁷O correction is a function built into the evaluation software of the IRMS. The algorithm and ratio assumptions are based on Assonov and Brenninkmeijer (2001). The same technique had been used to revise the VPDB scale by adding LSVEC as a second scaling point (Coplen et al., 2006b).

The EA–IRMS analysis of carbonates includes a wellcharacterised blank contribution that is due to the carbon impurities within the tin capsules that are used for carbonate analyses (Werner et al., 1999). In contrast, no such blank is expected when samples are analysed without tin capsules, as would be the case for gaseous CH₄ samples. While we did not observe a significant δ^{13} C difference when tin capsules were added to CH₄ injections and the δ^{13} C bias was subsequently corrected for or when the δ^{13} C–CH₄ analysis was performed without tin capsules. We continuously added the tin capsules to each δ^{13} C–CH₄ analysis and applied the routine blank correction to all measurements in compliance with the PIT between analyses of carbonate reference materials and CH₄ samples.

For $\delta^2 H$ analyses, we chose an analogue approach and process both H₂O and CH₄ using the high-temperature reactor of the TC/EA–IRMS system. Possible artefacts can arise mainly from the stronger surface activities of H₂O vs. CH₄ prior to the conversion to H₂ (and CO or carbon). H₂O injections can lead to memory effects, which need to be taken into

Table 6. Differences in δ^2 H–CH₄ and δ^{13} C–CH₄ between primary/secondary CH₄ gas calibrations and *i*SAAC measurements of the synthetic CH₄-in-air standards using the scale anchor based on Carina-1. Differences are calculated as $\delta_{iSAAC} - \delta_{pure}$. The bottom line shows the average and the standard deviation (1 σ) of considered differences, excluding the value of Biogenic (°) as described in main text.

Gas name	$\Delta \delta^2$ H–CH ₄ [‰]	$\Delta \delta^{13}$ C–CH ₄ [‰]
Megan	3.9	0.05
Merlin	5.6	-0.04
Minion	2.7	-0.05
Melly	4.3	0.13
Mike-1	5.7	-0.03
Martha-1	3.2	-0.06
Fossil	5.1	0.19
Biogenic	3.0	0.31 (°)
Average	$+4.2\pm1.2$	$+0.03\pm0.10$

account in δ^2 H–H₂O and subsequent δ^2 H–CH₄ analyses, either by discarding initial injections or making appropriate corrections (Werner and Brand, 2001). H₂O injections produced highest H₂ yields and stable δ^2 H–H₂O values at reactor temperatures of 1450 °C. Therefore we kept the reactor at 1450 °C during all calibration measurements. In addition, we found a minor dependence of $\delta^2 H - H_2 O$ on the septum temperature. We experimentally determined a septum temperature of 130 °C at which the effect on δ^2 H–H₂O was insignificant and kept the septum at 130 °C during all calibrations. We describe the experiments on reactor temperature and septum temperature in Appendix A in more detail. Note that it is essential to exclude systematic, material-specific errors to make H₂O and CH₄ reactions directly comparable for δ^2 H calibration. Based on these experiments we conclude that the δ^2 H–CH₄ calibrations do not contain measurement errors introduced by bracketing δ^2 H–H₂O analyses.

4.2 Discussion of the comparison between CIC and MPI-BGC

We compare the results of δ^2 H–CH₄ and δ^{13} C–CH₄ calibrations achieved by the two independent methods from CIC and MPI-BGC in Tables 4 and 5. Note that the verification of the principle calibration method (MPI-BGC) by an independent method (CIC) is required for the preparation of QCMs when CRMs are not available (IAEA, 2003). The comparison between CIC and MPI-BGC is to some degree representative of the situation of the community analysing atmospheric δ^2 H–CH₄ and δ^{13} C–CH₄ without access to international reference air but locally produced or propagated standard gases.

Even though there is no significant difference between the intercomparison results for δ^{13} C–CH₄, and the difference in δ^{2} H–CH₄ is rather small, there seems to be a systematic pattern that the samples combusted at CIC are generally more

enriched in both δ^2 H and δ^{13} C (Tables 4 and 5). The cause for this offset is not yet fully understood but will be discussed in more detail. The δ^{13} C–CH₄ calibrations presented in Table 4 were made on three different IRMS systems with three different working standards. All δ^{13} C measurements were corrected for potential scale compression effects, except from the MPI-BGC* analyses, which were made on an IRMS system specifically tuned to render scale compression effects for δ^{13} C, as demonstrated by Ghosh et al. (2005). Because the difference in δ^{13} C between Fossil and Biogenic is remarkably well resolved in all comparison measurements (Table 4), we conclude that our δ^{13} C comparison does not suffer from a significant scale compression error. Rather, the difference in δ^{13} C between the methods seems related to the method of CH₄ conversion. In principle, incomplete CH₄ combustion in the experiments at CIC would create a δ^{13} C pattern where the affected experiments appeared more enriched in δ^{13} C. This is because the remaining CH₄ fraction in the combustion-derived CO2 gas would be introduced into the dual-inlet IRMS together with the CO₂, and form CO₂H⁺ ions, which creates an artefact on m/z 45 (Sect. 4.1). However, we carefully tested every sample for residual CH₄ and are confident that the CH₄ combustions at CIC have been complete. Therefore, we cannot resolve this difference further.

We also observe a small δ^2 H–CH₄ offset between CIC and MPI-BGC. The δ^2 H measurements at CIC were made using combustion-derived H₂O with two different methods (TC/EA-IRMS and CRDS). Moreover, the measurement procedures at CIC included WSs covering the full VSMOW/SLAP scale. In contast, the direct δ^2 H–CH₄ analysis of the secondary CH₄ gases at MPI-BGC was performed as a one-point calibration against Megan or Merlin with a δ^2 H–CH₄ similar to that of Fossil (Table 3). Please note that $\delta^2 H$ scale compression often arises during the analysis of H₂O because it interacts with all sorts of surfaces in the analytical system. However, CH₄ gas behaves very much like pure H₂ in the high-temperature conversion system and a careful H_3^+ -factor determination often results in accurate isotopic distances. If the control of scale compression at MPI-BGC was limited due to the one-point calibration, we would expect the isotopic difference between Biogenic and Fossil to be smaller in the results from MPI-BGC than CIC. However, this is clearly not the case. The isotopic difference between Biogenic and Fossil ($\delta_{Fossil} - \delta_{Biogenic}$) appears to be very similar in the calibrations of both laboratories with 147.3 % at CIC and 147.9 % at MPI-BGC, even showing a slightly larger difference at MPI-BGC (Table 5). Therefore, we are confident that the observed, small $\delta^2 H$ offset is not caused by scale compression effects in one of the laboratories. Moreover, the excellent agreement between the experimentally controlled scale compression at CIC and the method at MPI-BGC proves that the analysis at MPI-BGC is free of significant scale compression artefacts over the tested isotopic range of $\sim 150 \%$.

The comparisons show small differences in the calibration results, but we found no evidence that either one of the two analytical methods is more accurate. Note that the difference in both δ^2 H–CH₄ and δ^{13} C–CH₄ exceeds the compatibility goal of 1 and 0.02‰ by a factor of 2 to 10 respectively (WMO, 2014). We interpret the results of this comparison to reflect calibration differences between laboratories that are to be expected, when CRMs are not available. Finally, we conclude that our new method is as capable to calibrate CH₄ gases to the international isotope scales and that it is as accurate as the method presented by Sperlich et al. (2012). However, we think that our new methods are more suitable for the task to produce and maintain a suite of calibration gases for the following reasons.

- The methods at MPI-BGC are more time efficient than the method of Sperlich et al. (2012). While the new methods at MPI-BGC can be used to calibrate an entire suite of CH_4 gases within a relatively short time, the method of Sperlich et al. (2012) is capable of processing only one sample per day.
- The new MPI-BGC methods are based on continuousflow IRMS and follow the PIT to the highest possible degree. In comparison, the method of Sperlich et al. (2012) is based on the combustion of CH₄ in an offline reactor, which requires re-oxidation after every sample and partial dismantling of the system to retrieve the sample for isotopic analysis. Because the analytical system at CIC could theoretically be at a different state for every sample (oxidation state, air leak rate) and because the system at CIC does not allow us to compare two CH₄ gases directly against each other, the methods at MPI-BGC are superior in the ability to fulfil the PIT. Even though the method at CIC proved to be very reproducible, we cannot rule out that a variation in the oxidation state of the reactor or an undetected air leakage into the system would affect the analysis of some CH₄ samples more than others. Because fulfilling the PIT is of paramount importance for isotope ratio analysis (e.g. Werner and Brand, 2001; Schimmelmann et al., 2016), we believe the method at MPI-BGC is less vulnerable to measurement errors in future calibrations.

4.3 Discussion on the compatibility between the scale anchors for δ^2 H–CH₄ and δ^{13} C–CH₄ as propagated from IMAU to MPI-BGC and JRAS-M16

We interpret the excellent agreement between the δ^{13} C and CH₄ calibrations in Carina-1 and Carina-2 from IMAU (Table 2) that both gases are precisely referenced and suitable for scale propagation from IMAU to MPI-BGC. The synthetic CH₄-in-air standards were analysed on *i*SAAC for δ^2 H–CH₄ and δ^{13} C–CH₄ and their isotope values were assigned using a WS that was calibrated against Carina-1. We can then

interpret the δ^{13} C difference between the *i*SAAC measurement and the calibrated synthetic CH₄-in-air standards of $+0.03 \pm 0.10\%$ as an accurate estimate for the calibration offset between the propagated scale anchor at MPI-BGC and the newly developed JRAS-M16.

Unfortunately, the situation is currently less straightforward for δ^2 H–CH₄. The two WSs Carina-1 and Carina-2 were calibrated at IMAU with a difference in δ^2 H–CH₄ of 2.8% that was insignificant at the time (Table 2). Because Carina-1 and Carina-2 appear indistinguishable in δ^2 H–CH₄ when compared to iSAAC with a measurement precision for δ^2 H–CH₄ of 1.0% (Sect. 2.6), we cannot determine the laboratory offset with the same certainty as for δ^{13} C–CH₄. If either Carina-1 or Carina-2 were representative for the calibrations at IMAU, the δ^2 H–CH₄ offset between the laboratories would amount to $+4.2 \pm 1.2$ or $+1.4 \pm 1.2\%$ respectively. A further comparison that includes new measurements on the current system at IMAU is required to determine the offset δ^2 H–CH₄ accurately. This offset can be resolved, for example, when a set of synthetic CH₄-in-air standards (JRAS-M16) is analysed at IMAU in future.

4.4 Discussion on possible use of synthetic CH₄-in-air standards in future

We demonstrated the ability to test the compatibility between IMAU and MPI-BGC by comparing scale anchors that were previously propagated from IMAU to MPI-BGC to JRAS-M16 gases. Future developments include an interlaboratory comparison to test whether a dedicated set of our synthetic CH₄-in-air standards (JRAS-M16) could provide a community anchor to the VPDB and VSMOW scales with documented accuracy. A further important test would be to determine to what extent the use of centrally calibrated standard gases could increase compatibility. A recent incidence provides a good example for the vulnerability of δ^{13} C–CH₄ observations in the atmosphere without suitable m-RM.

LSVEC, the second CRM anchor to the VPDB scale, has recently been discovered to be less reliable than anticipated. Until further notice, LSVEC is suggested to be treated with an enhanced δ^{13} C uncertainty of 0.15% (S. Assonov, personal communication, 2016). It is important to appreciate that this uncertainty is fully added to the uncertainty of δ^{13} C–CH₄ measurements, due to the similarity of LSVEC (-46.6%) and tropospheric CH₄ (-47.5%) in δ^{13} C. That is, the new uncertainty of LSVEC contributes the largest component in the full error budget of $\delta^{13}C$ –CH₄ analysis. Note that the suggested uncertainty of LSVEC is (i) on the order of the seasonal δ^{13} C–CH₄ cycle in the Southern Hemisphere and (ii) a multiple of the analytical precision of laboratories monitoring δ^{13} C–CH₄. If measurements of δ^{13} C–CH₄ considered the new uncertainty for LSVEC, the significance of signals such as the seasonal variability in the Southern Hemisphere would be lost on the cost of a better representation of accuracy. Including the uncertainty of LSVEC may further impact on the compatibility between several laboratories and, for example, suggest an artificially imposed spatial δ^{13} C–CH₄ gradient, based on calibration artefacts. We advocate the scientific gain when accuracy and compatibility are differentiated (WMO, 2014). The community benefits from a referencing method that enables a compatibility level that is smaller than the atmospheric δ^{13} C–CH₄ signal to resolve spatiotemporal δ^{13} C–CH₄ differences as primary goal. We think that establishing JRAS-M16 as community scale anchor could be a valuable step towards reaching this goal. As appropriate for any scale anchor that is intended to be usable for the whole community over long periods of time, the scale anchors will have to be re-calibrated frequently in order to detect possible drifts or to improve and correct previous assignments. The results of these efforts will be made available to the public at regular intervals.

We propose the distribution of JRAS-M16, a set of synthetic CH₄-in-air standards in 5L glass flasks. While two JRAS-M16 gases shall be used as calibration standard, an optional third JRAS-M16 gas can be used as unknown that is calibrated against the known JRAS-M16 gases as measurement control standard. This experiment would simulate the case when all participating laboratories measure the same sample directly against the same m-RM using the method that is otherwise applied to every sample in the respective laboratory and has the potential to determine the achievable compatibility. A further possibility to share the JRAS-M16 scale anchor would be to send cylinders with air-WSs to MPI-BGC for calibration. Because a dedicated target of this work is to achieve best possible accuracy with JRAS-M16, we provide the uncertainty of the full traceability chain. Once a new CRM has been found in replacement of LSVEC, the δ^{13} C–CH₄ and the traceability chain of JRAS-M16 will be revised accordingly. This will also be made upon future CRM revisions or replacements.

5 Conclusions

The number of laboratories that measure isotope ratios of atmospheric CH₄ is growing and combining data from multiple laboratories could enable new science and increasingly powerful analysis. However, merging data from multiple laboratories for analysis is currently hampered by the lack of reference materials that enable the community to produce a unified data set. To overcome this problem and to improve compatibility between laboratories, we produced synthetic CH₄-in-air standards (JRAS-M16). We modified standard online IRMS techniques to calibrate pure CH₄ gases for δ^2 H and δ^{13} C on international VSMOW and VPDB isotope scales respectively. Because such instrumentation is available to many isotope laboratories, our technical modifications and experiments can be reproduced elsewhere. Eight of the calibrated CH₄ gases were diluted with CH₄-free air in 5L glass flasks to produce synthetic CH₄-in-air standards with known δ^2 H–CH₄ and δ^{13} C–CH₄ values. These synthetic gas mixtures were then analysed on a newly developed system (*i*SAAC) to measure δ^2 H–CH₄ and δ^{13} C–CH₄ in air samples. Hitherto, iSAAC used working standards as scale anchors for δ^2 H–CH₄ and δ^{13} C–CH₄, which were calibrated at a partnering institute (IMAU). The history of the propagated isotope scales goes more than 2 decades back in time and includes the propagation between several laboratories. We determine δ^2 H–CH₄ and δ^{13} C–CH₄ in our synthetic CH₄-in-air standards using the scale anchor propagation from IMAU and compare the results with our calibration results for δ^2 H–CH₄ and δ^{13} C–CH₄. We use this method to determine the δ^{13} C–CH₄ offsets between the scale anchor that was propagated from IMAU and JRAS-M16, thereby providing a method to improve compatibility. Further comparisons are required to determine the offset for $\delta^2 H - CH_4$.

We welcome other laboratories to further test our calibrations by analysing JRAS-M16 air sets, which will be available upon request. Another possibility could be to have cylinders with air-WSs sent to MPI-BGC for calibration using JRAS-M16 as scale anchor. JRAS-M16 may help laboratories to anchor δ^2 H–CH₄ and δ^{13} C–CH₄ observations to unified community scale anchors. This might be a useful step towards reaching the compatibility goals between laboratories, leading to an improved understanding of atmospheric CH₄. Future work includes a revision of the δ^{13} C–CH₄ calibrations once the replacement for LSVEC is established. This will reduce the uncertainty of the δ^{13} C–CH₄ scale anchors significantly. The LSVEC replacement should extend to the δ^{13} C-depleted range of biogenic CH₄ gases.

6 Data availability

The results of our final calibrations with the associated uncertainties of the full traceability chains are published as a Supplement to this paper. The supplementary data file also contains the revised calibrations of the data by Sperlich et al. (2012). These include corrections for the offset in RM8563 and for scale compression effect in the IRMS at CIC.

Appendix A: Experiments to enhance the performance of the analytical system for the calibration of δ^2 H–CH₄ with H₂O



Figure A1. The δ^2 H variation of H₂O injections with septum temperatures. Blue circles show average δ^2 H–H₂O values for septum temperatures above 90 °C, the black line is the quadratic polynomial fit to the data above 90 °C while red diamonds display δ^2 H–H₂O values at septum temperatures below 90 °C. The error bars show 1 σ standard deviations and the grey-dashed lines indicate the typical precision limit of 1‰ for δ^2 H–H₂O analysis (Gehre et al., 2004) around the δ^2 H–H₂O value of the polynomial fit for the septum temperature of 130 °C (set point during calibration experiments). The grey dashed lines show that our δ^2 H–H₂O analyses remain within a typical precision level as long as the septum temperature is controlled to ~ 130 ± 10 °C.

The injection of H₂O samples into the reactor is critical because it is prone to isotopic fractionation (Werner and Brand, 2001). This fractionation is mainly caused by system memory due to adhesion of injected H₂O to the reactor walls. The isotopic fractionation can be overcome by repetitive injections of H₂O samples with identical isotopic composition, thereby overwriting the memory effect until it reaches a marginal level. For H₂O analyses under constant analytical conditions (e.g. constant reactor temperature), the adhesion effect is a function mainly of the amount of injected H₂O sample. Moreover, the effect on the isotopic composition scales with the isotopic difference between two consecutive samples (Gehre et al., 2004). Because there is no adhesion of the sample during CH₄ analysis, this memory effect is most pronounced only during the analysis of H₂O in our study. Subsequent CH₄ analysis does not contribute to system memory but can still be affected by H₂O desorption from internal surfaces of the analytical system. Therefore, memory effects of H₂O can propagate into the CH₄ calibrations. Memory effects are identified in a series of replicate H₂O measurements and are corrected for by modelling the memory function as described in Gehre et al. (2004) and Brand et al. (2009a) on a routine basis, as our system has been used for isotope analysis of H₂O samples for more than a decade. We conclude that our results are free of artefacts arising from sample memory.



Figure A2. The dependence of $\delta^2 H$ and H_2 peak areas of H_2O and CH_4 injections from reactor temperatures between 1300 and 1450 °C. Top and bottom panels show H_2O and CH_4 experiments respectively. $\delta^2 H$ isotope ratios are shown in blue for H_2O and green for CH_4 and refer to the left-hand axes. Average H_2 peak areas are indicated by grey crosses and refer to the right-hand axes. All error bars indicate the standard deviation. The red diamond shows the average peak area and the respective standard deviation including the outliers (see Appendix text). *Y* axes ranges are matched between top and bottom panels to enable direct comparison of the temperature effect for H_2O and CH_4 . Equations describe the fits in both panels, displayed by dashed lines. Continuous lines in the bottom panel indicate the 95 % confidence interval of the linear fit.

Isotopic fractionation during the analysis of the reference waters can also be caused by insufficiently heated septa (Gehre et al., 2004). We injected 106 identical H₂O samples while we increased the septum temperatures in nine steps from 76 to 137 °C. In general, we observed a δ^2 H enrichment with increasing septum temperature. A systematic increase of δ^2 H–H₂O with septum temperature is apparent above 90 °C until δ^2 H–H₂O values plateau at septum temperatures around 130 °C (Fig. A1, blue circles). The stabilising δ^2 H–H₂O at high temperatures suggests quantitative H₂O processing without significant isotope fractionation, in line with previous observations (Gehre et al., 2004). In contrast, the three δ^2 H–H₂O values below 90 °C (red diamonds) show an insignificant but slight increase in δ^2 H–H₂O with septum temperature, which deviated from the pattern above 90 °C. We cannot explain the mismatch between the two patterns above and below 90 °C. We speculate that the initial heating of the septum to temperatures between 70 and 90 °C caused the desorption of accumulated of H₂O, which was desorbed once the septum was heated to temperatures above 90 °C.

Quantitative conversion of both CH_4 and H_2O in the hightemperature reactor is of utmost importance for our study, because incomplete conversion causes isotopic fractionation in the reaction products (e.g. Burgoyne and Hayes, 1998; Hilkert et al., 1999; Gehre et al., 2004). The reactor temperature is critical for the efficiency of the conversion process. We performed an experiment with CH₄ and H₂O injections at different reactor temperatures (Fig. A2). For water injections we observe a pronounced, nonlinear $\delta^2 H - H_2 O$ change of $\sim 15\%$ with reactor temperature increase from 1300 to 1450 °C, reaching a plateau above 1400 °C. The pattern is consistent with previous observations in both trend and magnitude (Gehre et al., 2004). In contrast, the linear fit for δ^2 H–CH₄ increases by only about 1 % over the 150 °C temperature range. However, the slope is statistically insignificant as shown by the 95 % confidence interval of the linear fit (Fig. A2). This analyte-specific isotope variation is also reflected in the areas of the H₂O and CH₄-derived H₂ peaks (Fig. A2) (with some significant scatter in the data). While the H₂O-derived H₂ peak areas increase with increasing reactor temperature, the CH₄-derived H₂ peak areas remain constant within the error bars throughout the experiments. For an unknown reason, three out of six H₂ peaks that resulted from H₂O injections at 1400 °C were by 10-15 standard deviations smaller than the remaining three peaks. We present the averages and 1σ standard of the H₂ peaks with and without removal of these outliers in Fig. A2, which shows the exceptional pattern at 1400 °C. Despite this peak size variability, the isotopic composition of all H₂O injections at 1400 °C is in good agreement. Our experiments indicate that reactor temperatures in excess of 1400 °C are required especially for quantitative conversion of H₂O, while the effects of reactor temperature on both yield and the isotopic composition of CH₄-derived H₂ are comparably small. Therefore, we operate the reactor at a temperature of 1450 °C to guarantee quantitative conversion without isotope fractionation of both H₂O (Gehre et al., 2004) and CH₄ (Burgoyne and Hayes, 1998; Hilkert et al., 1999).

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