

# Ricci-flat metrics on complex cones

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# Motivation

- Our wish is to understand the  $AdS/CFT$  correspondence away from the maximally supersymmetric case, but in situations where dual theories can still be under control ( $\mathcal{N} = 1$ )
- However, in this talk we will mainly concentrate on the geometric aspects of the problem

# $AdS/CFT$ with $\mathcal{N} = 1$ SUSY

- Classical example of  $AdS/CFT$ :  $AdS_5 \times S^5$ , dual to  $\mathcal{N} = 4$  SYM in  $d = 4$  [Maldacena \[1997\]](#)
- There exist extremal black hole solutions to IIB supergravity preserving  $\mathcal{N} = 1$  SUSY, whose ‘near-horizon’ geometry is  $AdS_5 \times X^5$  [Morrison, Plesser \[1998\]](#)
- $X^5$  is a Sasaki-Einstein manifold (this implies the existence of one Killing spinor)

# Sasaki-Einstein manifolds

- $X$  is Sasaki-Einstein iff the cone over it is Kähler and Ricci-flat:

$$ds^2 = dr^2 + r^2 (\widetilde{ds^2})_X$$

- $\widetilde{ds^2}$  Kähler & Ricci-flat  $\Leftrightarrow$   $(ds^2)_X$  Sasaki-Einstein, of positive curvature
- The metric can be written as  $(ds^2)_{X^5} = (d\phi - J)^2 + (ds^2)_{\mathcal{M}}$  where  $(ds^2)_{\mathcal{M}}$  is Kähler-Einstein (but not necessarily smooth),  $J$  is the Kähler current
- $r = 0 \rightarrow$  singularity

# Resolving the singularity of the cone

- It is possible to resolve the singularity of the conical metric by ‘blowing-up’ the vertex, i.e. by replacing it with a cycle of non-zero size
- The metric at infinity, i.e. at  $r \rightarrow \infty$ , will still be asymptotic to the cone:

$$ds^2 = dr^2 + r^2 (\widetilde{ds^2})_X \quad \text{for } r \rightarrow \infty$$

- Apart from simplest cases, resolved metrics on the cones are not known  $\Rightarrow$  Our study

# Some examples

- Eguchi, Hanson, 1978

Complex dimension 2, singularity of the form

$$\mathbb{C}^2/\mathbb{Z}_2 : (z_1, z_2) \sim (-z_1, -z_2)$$

- Introducing invariant coordinates

$X = z_1^2, Y = z_2^2, Z = z_1 z_2$ , we get an

equation  $XY = Z^2$  in  $\mathbb{C}^3$

- This corresponds to the cone in the embedding of  $\mathbb{CP}^1$  by the linear system  $|\mathcal{O}(2)|$ , i.e. the *anticanonical* embedding

# The Eguchi-Hanson metric

- One can look for the Kähler potential of the form  $K = K(|z_1|^2 + |z_2|^2)$ .  
The metric is, as usual,  $ds^2 = \partial_i \bar{\partial}_j K dz^i d\bar{z}^j$
- For a Kähler metric the Ricci tensor can be expressed as  $R_{i\bar{j}} = -\partial_i \bar{\partial}_j \log \det g$
- Set  $R_{i\bar{j}} = 0$ , solve for the Kähler potential:

## The Eguchi-Hanson metric

$$K = \sqrt{r^2 + 4x^2} + r \log \left( \frac{\sqrt{r^2 + 4x^2} - r}{2x} \right), \quad r > 0$$

# More examples

- 2d case: Eguchi-Hanson = anticanonical cone over  $\mathbf{CP}^1 \Rightarrow \text{SE } X_3 = S^3/\mathbb{Z}_2$
- ‘3d Eguchi-Hanson’ = anticanonical cone over  $\mathbf{CP}^2 \Rightarrow \text{SE } X_5 = S^5/\mathbb{Z}_3$
- 3d case: Candelas-de la Ossa [1990] = anticanonical cone over  $\mathbf{CP}^1 \times \mathbf{CP}^1$  (resolved conifold)  
 $\Rightarrow \text{SE } X_5 := T^{1,1} = \frac{SU(2) \times SU(2)}{U(1)}$



# Other cones?

- One can only build Ricci-flat cones over complex manifolds of ‘positive curvature’ (i.e. with ample anticanonical class)
- For the cone to be of  $\dim_{\mathbb{C}} = 3$ , we take the underlying base to be of  $\dim_{\mathbb{C}} = 2$
- Apart from  $\mathbb{C}P^2$  and  $\mathbb{C}P^1 \times \mathbb{C}P^1$ , there are only 8 other positively curved complex surfaces – the del Pezzo surfaces

$dP_1, \dots, dP_8$

# The del Pezzo surface $dP_1$

- $dP_n$  can be seen as  $CP^2$ , blown-up in  $n$  sufficiently generic points
- We will consider the simplest non-homogeneous case, i.e. the cone over  $dP_1$
- Any metric on  $dP_1$  should have at least two parameters – the sizes of  $CP^2$  and of the blown-up  $CP^1$
- Do these parameters persist in the cone over  $dP_1$ ?

# Isometries

- Whereas the automorphism group of  $\mathbf{CP}^2$  is  $PGL(3, \mathbf{C})$ , the automorphism group of the del Pezzo surface is reduced to

$$Aut(dP_1) = P \begin{pmatrix} \bullet & \bullet & 0 \\ \bullet & \bullet & 0 \\ \bullet & \bullet & \bullet \end{pmatrix} \quad (1)$$

- The isometry group of the metric on the *cone* is the maximal compact subgroup of the parabolic subgroup shown above, i.e.

$$\mathbf{Isom} = U(1) \times U(2)$$

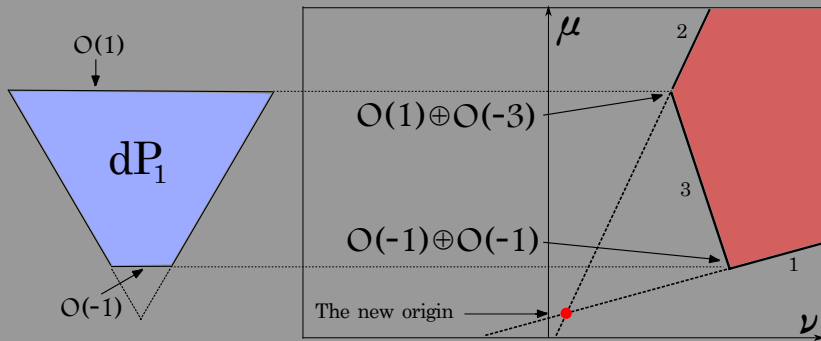
# The main equation

- We will look for a Kähler potential of the form  $K = K(|u|^2, |z_1|^2 + |z_2|^2) := K(e^t, e^s)$
- Just as in the case of the Eguchi-Hanson metric, we can write out a Ricci-flatness equation
- More convenient to perform a Legendre transform w.r.t.  $t, s$ , introducing the dual momentum maps  $\mu = \frac{\partial K}{\partial t}$ ,  $\nu = \frac{\partial K}{\partial s}$  and a dual potential  $G = t\mu + s\nu - K$

# The equation

$$e^{G_\mu + G_\nu} (G_{\mu\mu} G_{\nu\nu} - G_{\mu\nu}^2) = \mu$$

- The domain – the moment polygon



# The expansion at $\infty$

- We can solve the equation exactly at large  $\mu, \nu$  with fixed ‘angle’  $\xi = \frac{\mu}{\nu}$ , assuming the conical form of the metric
- This gives  $G = 3\nu(\log \nu - 1) + \nu P_0(\xi)$
- $P_0(\xi)$  satisfies an ODE and can be found exactly. It provides a Sasaki-Einstein metric, which in the  $dP_1$  case is the  $Y^{2,1}$  manifold  
[Gauntlett, Martelli, Sparks, Waldram \[2004\]](#)

# $M^{\text{th}}$ order and the Heun equation

- We can build a systematic perturbation theory

$$G = 3\nu(\log \nu - 1) + \nu P_0(\xi) + \log \nu + \sum_{k=0}^{\infty} \nu^{-k} P_{k+1}(\xi)$$

- In order  $\nu^{-M}$  we obtain the equation

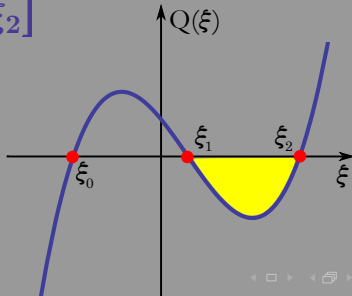
$$\frac{d}{d\xi} \left( Q(\xi) \frac{dP_M}{d\xi} \right) - \left( (M-2)^2 - 1 \right) \xi P_M = \text{r.h.s.},$$

$$\text{where } Q(\xi) = \xi^3 - \frac{3}{2}\xi^2 + d$$

- This is a Heun equation – an analogue of hypergeometric equation with 4 Fuchsian singularities on  $\mathbb{CP}^1$

# Resolution parameters

- All resolution parameters should arise as (coefficients in front of) the solutions to the homogeneous equation in some order of perturbation theory
- The equation is solved in a ‘physical’ interval  $\xi \in [\xi_1, \xi_2]$



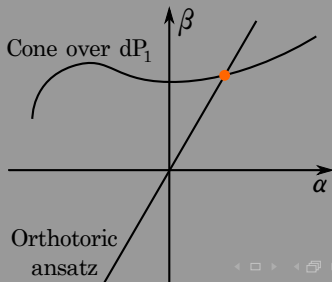


## Resolution parameters. 2.

- Regularity of the metric at the boundaries of the moment polytope requires that the solutions should be regular at  $\xi = \xi_1, \xi_2$   
 $\Rightarrow$  Eigenvalue problem
- Solutions exist for  $M = 3, 4$ :  
 $P_3 = \alpha, P_4 = \beta (\xi - 1)$
- Conjecture:  
For other  $M$  solutions do not exist

# Resolution parameters. 3.

- When  $\beta = -\frac{\alpha}{2\xi_0}$ , the resolved metric is known [Calderbank, Gauduchon \[2006\]](#), [Chen, Lu, Pope \[2006\]](#)
- In general, topology imposes one more relation between  $\beta$  and  $\alpha$  [Martelli, Sparks \[2007\]](#)
- Hence the general situation is as follows:



# Questions / Answers

- Can one obtain an exact formula with both parameters  $\alpha, \beta$ ?
- As we discussed, there is an exact formula when  $\beta = -\frac{\alpha}{2\xi_0}$ . Is there a generalization?
- Dual field theories for  $AdS_5 \times X^5$  have been conjectured Feng, Hanany, He, 2000
- Interpret the new parameter in terms of these  $\mathcal{N} = 1$  gauge theories

Thank you!