

GW170608: OBSERVATION OF A 19-SOLAR-MASS BINARY BLACK HOLE COALESCENCE

B. P. ABBOTT,¹ R. ABBOTT,¹ T. D. ABBOTT,² F. ACERNESE,^{3,4} K. ACKLEY,^{5,6} C. ADAMS,⁷ T. ADAMS,⁸ P. ADDESSO,⁹
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arXiv:1711.05578v1 [astro-ph.HE] 15 Nov 2017

V. GAYATHRI,¹¹⁶ N. GEHRELS,^{79,*} G. GEMME,⁵⁹ E. GENIN,³⁰ A. GENNAI,²⁴ D. GEORGE,¹² J. GEORGE,⁶⁰ L. GERGELY,¹¹⁷
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M. MANSKE,^{21,25} M. MANTOVANI,³⁰ F. MARCHESONI,^{52,43} F. MARION,⁸ S. MÁRKA,⁵⁰ Z. MÁRKA,⁵⁰ C. MARKAKIS,¹²
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R. ORMISTON,⁴⁵ L. F. ORTEGA,⁵ R. O'SHAUGHNESSY,⁵⁷ S. OSSOKINE,³⁸ D. J. OTTAWAY,⁷² H. OVERMIER,⁷ B. J. OWEN,⁸³
A. E. PACE,⁶³ J. PAGE,¹³⁷ M. A. PAGE,⁶⁴ A. PAI,^{116,145} S. A. PAI,⁶⁰ J. R. PALAMOS,⁶⁹ O. PALASHOV,¹²⁸ C. PALOMBA,³⁵
A. PAL-SINGH,³⁴ HOWARD PAN,⁸⁷ HUANG-WEI PAN,⁸⁷ B. PANG,⁴⁸ P. T. H. PANG,⁹³ C. PANKOW,⁸⁹ F. PANNARALE,³⁶
B. C. PANT,⁶⁰ F. PAOLETTI,²⁴ A. PAOLI,³⁰ M. A. PAPA,^{38,21,10} A. PARIDA,¹⁹ W. PARKER,⁷ D. PASCUCCI,⁴⁶
A. PASQUALETTI,³⁰ R. PASSAQUIETI,^{23,24} D. PASSUELLO,²⁴ M. PATIL,¹³³ B. PATRICELLI,^{146,24} B. L. PEARLSTONE,⁴⁶
M. PEDRAZA,¹ R. PEDURAND,^{26,147} L. PEKOWSKY,⁴⁴ A. PELE,⁷ S. PENN,¹⁴⁸ C. J. PEREZ,⁴⁷ A. PERRECA,^{1,110,95}
L. M. PERRI,⁸⁹ H. P. PFEIFFER,^{90,38} M. PHELPS,⁴⁶ O. J. PICCINI,^{97,35} M. PICHOT,⁶⁶ F. PIERGIOVANNI,^{121,122} V. PIERRO,⁹

G. PILLANT,³⁰ L. PINARD,²⁶ I. M. PINTO,⁹ M. PIRELLO,⁴⁷ M. PITKIN,⁴⁶ M. POE,²¹ R. POGGIANI,^{23,24} P. POPOLIZIO,³⁰ E. K. PORTER,³⁹ A. POST,¹⁰ J. POWELL,^{46,149} J. PRASAD,¹⁹ J. W. W. PRATT,³⁷ G. PRATTEN,¹⁰² V. PREDOI,³⁶ T. PRESTEGARD,²¹ M. PRIJATELJ,¹⁰ M. PRINCIPE,⁹ S. PRIVITERA,³⁸ G. A. PRODI,^{110,95} L. G. PROKHOROV,⁶¹ O. PUNCKEN,¹⁰ M. PUNTURO,⁴³ P. PUPPO,³⁵ M. PÜRRER,³⁸ H. QI,²¹ V. QUETSCHKE,¹⁰³ E. A. QUINTERO,¹ R. QUITZOW-JAMES,⁶⁹ F. J. RAAB,⁴⁷ D. S. RABELING,²⁵ H. RADKINS,⁴⁷ P. RAFFAI,⁵⁵ S. RAJA,⁶⁰ C. RAJAN,⁶⁰ B. RAJBHANDARI,⁸³ M. RAKHMANOV,¹⁰³ K. E. RAMIREZ,¹⁰³ A. RAMOS-BUADES,¹⁰² P. RAPAGNANI,^{97,35} V. RAYMOND,³⁸ M. RAZZANO,^{23,24} J. READ,²⁹ T. REGIMBAU,⁶⁶ L. REI,⁵⁹ S. REID,⁶² D. H. REITZE,^{1,5} W. REN,¹² S. D. REYES,⁴⁴ F. RICCI,^{97,35} P. M. RICKER,¹² S. RIEGER,¹⁰ K. RILES,¹¹⁸ M. RIZZO,⁵⁷ N. A. ROBERTSON,^{1,46} R. ROBBIE,⁴⁶ F. ROBINET,²⁸ A. ROCCHI,³³ L. ROLLAND,⁸ J. G. ROLLINS,¹ V. J. ROMA,⁶⁹ R. ROMANO,^{3,4} C. L. ROMEL,⁴⁷ J. H. ROMIE,⁷ D. ROSIŃSKA,^{150,56} M. P. ROSS,¹⁵¹ S. ROWAN,⁴⁶ A. RÜDIGER,¹⁰ P. RUGGI,³⁰ G. RUTINS,²⁷ K. RYAN,⁴⁷ S. SACHDEV,¹ T. SADECKI,⁴⁷ L. SADEGHIAN,²¹ M. SAKELLARIADOU,¹⁵² L. SALCONI,³⁰ M. SALEEM,¹¹⁶ F. SALEMI,¹⁰ A. SAMAJDAR,¹⁴³ L. SAMMUT,⁶ L. M. SAMPSON,⁸⁹ E. J. SANCHEZ,¹ L. E. SANCHEZ,¹ N. SANCHIS-GUAL,⁸⁵ V. SANDBERG,⁴⁷ J. R. SANDERS,⁴⁴ B. SASSOLAS,²⁶ B. S. SATHYAPRAKASH,^{63,36} P. R. SAULSON,⁴⁴ O. SAUTER,¹¹⁸ R. L. SAVAGE,⁴⁷ A. SAWADSKY,³⁴ P. SCHALE,⁶⁹ M. SCHEEL,⁸⁹ J. SCHEUER,⁸⁹ J. SCHMIDT,¹⁰ P. SCHMIDT,^{1,65} R. SCHNABEL,³⁴ R. M. S. SCHOFIELD,⁶⁹ A. SCHÖNBECK,³⁴ E. SCHREIBER,¹⁰ D. SCHUETTE,^{10,22} B. W. SCHULTE,¹⁰ B. F. SCHUTZ,^{36,10} S. G. SCHWALBE,³⁷ J. SCOTT,⁴⁶ S. M. SCOTT,²⁵ E. SEIDEL,¹² D. SELLERS,⁷ A. S. SENGUPTA,¹⁵³ D. SENTENAC,³⁰ V. SEQUINO,^{32,33,17} A. SERGEEV,¹²⁸ D. A. SHADDOCK,²⁵ T. J. SHAFFER,⁴⁷ A. A. SHAH,¹³⁷ M. S. SHAHRIAR,⁸⁹ M. B. SHANER,¹⁰⁹ L. SHAO,³⁸ B. SHAPIRO,⁵¹ P. SHAWHAN,⁷⁵ A. SHEPERD,²¹ D. H. SHOEMAKER,¹⁵ D. M. SHOEMAKER,⁷⁶ K. SIELLEZ,⁷⁶ X. SIEMENS,²¹ M. SIENIAWSKA,⁵⁶ D. SIGG,⁴⁷ A. D. SILVA,¹⁶ L. P. SINGER,⁷⁹ A. SINGH,^{38,10,22} A. SINGHAL,^{17,35} A. M. SINTES,¹⁰² B. J. J. SLAGMOLEN,²⁵ B. SMITH,⁷ J. R. SMITH,²⁹ R. J. E. SMITH,^{1,6} S. SOMALA,¹⁵⁴ E. J. SON,¹³¹ J. A. SONNENBERG,²¹ B. SORAZU,⁴⁶ F. SORRENTINO,⁵⁹ T. SOURADEEP,¹⁹ A. P. SPENCER,⁴⁶ A. K. SRIVASTAVA,¹⁰⁵ K. STAATS,³⁷ A. STALEY,⁵⁰ M. STEINKE,¹⁰ J. STEINLECHNER,^{34,46} S. STEINLECHNER,³⁴ D. STEINMEYER,¹⁰ S. P. STEVENSON,^{58,149} R. STONE,¹⁰³ D. J. STOPS,⁵⁸ K. A. STRAIN,⁴⁶ G. STRATTA,^{121,122} S. E. STRIGIN,⁶¹ A. STRUNK,⁴⁷ R. STURANI,¹⁵⁵ A. L. STUVER,⁷ T. Z. SUMMERSCALES,¹⁵⁶ L. SUN,⁹⁶ S. SUNIL,¹⁰⁵ J. SURESH,¹⁹ P. J. SUTTON,³⁶ B. L. SWINKELS,³⁰ M. J. SZCZEPAŃCZYK,³⁷ M. TACCA,¹⁴ S. C. TAIT,⁴⁶ C. TALBOT,⁶ D. TALUKDER,⁶⁹ D. B. TANNER,⁵ M. TÁPAI,¹¹⁷ A. TARACCHINI,³⁸ J. D. TASSON,⁷ J. A. TAYLOR,¹³⁷ R. TAYLOR,¹ S. V. TEWARI,¹⁴⁸ T. THEEG,¹⁰ F. THIES,¹⁰ E. G. THOMAS,⁵⁸ M. THOMAS,⁷ P. THOMAS,⁴⁷ K. A. THORNE,⁷ E. THRANE,⁶ S. TIWARI,^{17,95} V. TIWARI,³⁶ K. V. TOKMAKOV,⁶² K. TOLAND,⁴⁶ M. TONELLI,^{23,24} Z. TORNASI,⁴⁶ A. TORRES-FORNÉ,⁸⁵ C. I. TORRIE,¹ D. TÖYRÄ,⁵⁸ F. TRAVASSO,^{30,43} G. TRAYLOR,⁷ J. TRINASTIC,⁵ M. C. TRINGALI,^{110,95} L. TROZZO,^{157,24} K. W. TSANG,¹⁴ M. TSE,¹⁵ R. TSO,¹ L. TSUKADA,⁸¹ D. TSUNA,⁸¹ D. TUYENBAYEV,¹⁰³ K. UENO,²¹ D. UGOLINI,¹⁵⁸ C. S. UNNIKRIISHNAN,¹¹⁹ A. L. URBAN,¹ S. A. USMAN,³⁶ H. VAHLBRUCH,²² G. VAJENTE,¹ G. VALDES,² N. VAN BAKEL,¹⁴ M. VAN BEUZEKOM,¹⁴ J. F. J. VAN DEN BRAND,^{74,14} C. VAN DEN BROECK,^{14,159} D. C. VANDER-HYDE,⁴⁴ L. VAN DER SCHAAF,¹⁴ J. V. VAN HEIJNINGEN,¹⁴ A. A. VAN VEGGEL,⁴⁶ M. VARDARO,^{53,54} V. VARMA,⁴⁸ S. VASS,¹ M. VASÚTH,⁴⁹ A. VECCHIO,⁵⁸ G. VEDOVATO,⁵⁴ J. VEITCH,⁴⁶ P. J. VEITCH,⁷² K. VENKATESWARA,¹⁵¹ G. VENUGOPALAN,¹ D. VERKINDT,⁸ F. VETRANO,^{121,122} A. VICERÉ,^{121,122} A. D. VIETS,²¹ S. VINCIGUERRA,⁵⁸ D. J. VINE,²⁷ J.-Y. VINET,⁶⁶ S. VITALE,¹⁵ T. VO,⁴⁴ H. VOCCA,^{42,43} C. VORVICK,⁴⁷ S. P. VYATCHANIN,⁶¹ A. R. WADE,¹ L. E. WADE,⁸⁴ M. WADE,⁸⁴ R. WALET,¹⁴ M. WALKER,²⁹ L. WALLACE,¹ S. WALSH,^{38,10,21} G. WANG,^{17,122} H. WANG,⁵⁸ J. Z. WANG,⁶³ W. H. WANG,¹⁰³ Y. F. WANG,⁹³ R. L. WARD,²⁵ J. WARNER,⁴⁷ M. WAS,⁸ J. WATCHI,⁹⁸ B. WEAVER,⁴⁷ L.-W. WEI,^{10,22} M. WEINERT,¹⁰ A. J. WEINSTEIN,¹ R. WEISS,¹⁵ L. WEN,⁶⁴ E. K. WESSEL,¹² P. WESSELS,¹⁰ J. WESTERWECK,¹⁰ T. WESTPHAL,¹⁰ K. WETTE,²⁵ J. T. WHELAN,⁵⁷ B. F. WHITING,⁵ C. WHITTLE,⁶ D. WILKEN,¹⁰ D. WILLIAMS,⁴⁶ R. D. WILLIAMS,¹ A. R. WILLIAMSON,⁶⁵ J. L. WILLIS,^{1,160} B. WILLKE,^{22,10} M. H. WIMMER,¹⁰ W. WINKLER,¹⁰ C. C. WIPF,¹ H. WITTEL,^{10,22} G. WOAN,⁴⁶ J. WOEHLENER,¹⁰ J. WOFFORD,⁵⁷ K. W. K. WONG,⁹³ J. WORDEN,⁴⁷ J. L. WRIGHT,⁴⁶ D. S. WU,¹⁰ D. M. WYSOCKI,⁵⁷ S. XIAO,¹ H. YAMAMOTO,¹ C. C. YANCEY,⁷⁵ L. YANG,¹⁶¹ M. J. YAP,²⁵ M. YAZBACK,⁵ HANG YU,¹⁵ HAOCUN YU,¹⁵ M. YVERT,⁸ A. ZADROŻNY,¹³² M. ZANOLIN,³⁷ T. ZELENKOVA,³⁰ J.-P. ZENDRI,⁵⁴ M. ZEVIN,⁸⁹ L. ZHANG,¹ M. ZHANG,¹⁴⁰ T. ZHANG,⁴⁶ Y.-H. ZHANG,⁵⁷ C. ZHAO,⁶⁴ M. ZHOU,⁸⁹ Z. ZHOU,⁸⁹ S. J. ZHU,^{38,10} X. J. ZHU,⁶ A. B. ZIMMERMAN,⁹⁰ M. E. ZUCKER,^{1,15} AND J. ZWEIZIG¹

LIGO SCIENTIFIC COLLABORATION AND VIRGO COLLABORATION

¹LIGO, California Institute of Technology, Pasadena, CA 91125, USA

²Louisiana State University, Baton Rouge, LA 70803, USA

³Università di Salerno, Fisciano, I-84084 Salerno, Italy

⁴INFN, Sezione di Napoli, Complesso Universitario di Monte S. Angelo, I-80126 Napoli, Italy

⁵University of Florida, Gainesville, FL 32611, USA

⁶OzGrav, School of Physics & Astronomy, Monash University, Clayton 3800, Victoria, Australia

⁷LIGO Livingston Observatory, Livingston, LA 70754, USA

⁸Laboratoire d'Annecy-le-Vieux de Physique des Particules (LAPP), Université Savoie Mont Blanc, CNRS/IN2P3, F-74941 Annecy, France

-
- ⁹ *University of Sannio at Benevento, I-82100 Benevento, Italy and INFN, Sezione di Napoli, I-80100 Napoli, Italy*
- ¹⁰ *Max Planck Institute for Gravitational Physics (Albert Einstein Institute), D-30167 Hannover, Germany*
- ¹¹ *The University of Mississippi, University, MS 38677, USA*
- ¹² *NCSA, University of Illinois at Urbana-Champaign, Urbana, IL 61801, USA*
- ¹³ *University of Cambridge, Cambridge CB2 1TN, United Kingdom*
- ¹⁴ *Nikhef, Science Park, 1098 XG Amsterdam, The Netherlands*
- ¹⁵ *LIGO, Massachusetts Institute of Technology, Cambridge, MA 02139, USA*
- ¹⁶ *Instituto Nacional de Pesquisas Espaciais, 12227-010 São José dos Campos, São Paulo, Brazil*
- ¹⁷ *Gran Sasso Science Institute (GSSI), I-67100 L'Aquila, Italy*
- ¹⁸ *INFN, Laboratori Nazionali del Gran Sasso, I-67100 Assergi, Italy*
- ¹⁹ *Inter-University Centre for Astronomy and Astrophysics, Pune 411007, India*
- ²⁰ *International Centre for Theoretical Sciences, Tata Institute of Fundamental Research, Bengaluru 560089, India*
- ²¹ *University of Wisconsin-Milwaukee, Milwaukee, WI 53201, USA*
- ²² *Leibniz Universität Hannover, D-30167 Hannover, Germany*
- ²³ *Università di Pisa, I-56127 Pisa, Italy*
- ²⁴ *INFN, Sezione di Pisa, I-56127 Pisa, Italy*
- ²⁵ *OzGrav, Australian National University, Canberra, Australian Capital Territory 0200, Australia*
- ²⁶ *Laboratoire des Matériaux Avancés (LMA), CNRS/IN2P3, F-69622 Villeurbanne, France*
- ²⁷ *SUPA, University of the West of Scotland, Paisley PA1 2BE, United Kingdom*
- ²⁸ *LAL, Univ. Paris-Sud, CNRS/IN2P3, Université Paris-Saclay, F-91898 Orsay, France*
- ²⁹ *California State University Fullerton, Fullerton, CA 92831, USA*
- ³⁰ *European Gravitational Observatory (EGO), I-56021 Cascina, Pisa, Italy*
- ³¹ *Chennai Mathematical Institute, Chennai 603103, India*
- ³² *Università di Roma Tor Vergata, I-00133 Roma, Italy*
- ³³ *INFN, Sezione di Roma Tor Vergata, I-00133 Roma, Italy*
- ³⁴ *Universität Hamburg, D-22761 Hamburg, Germany*
- ³⁵ *INFN, Sezione di Roma, I-00185 Roma, Italy*
- ³⁶ *Cardiff University, Cardiff CF24 3AA, United Kingdom*
- ³⁷ *Embry-Riddle Aeronautical University, Prescott, AZ 86301, USA*
- ³⁸ *Max Planck Institute for Gravitational Physics (Albert Einstein Institute), D-14476 Potsdam-Golm, Germany*
- ³⁹ *APC, AstroParticule et Cosmologie, Université Paris Diderot, CNRS/IN2P3, CEA/Irfu, Observatoire de Paris, Sorbonne Paris Cité, F-75205 Paris Cedex 13, France*
- ⁴⁰ *Korea Institute of Science and Technology Information, Daejeon 34141, Korea*
- ⁴¹ *West Virginia University, Morgantown, WV 26506, USA*
- ⁴² *Università di Perugia, I-06123 Perugia, Italy*
- ⁴³ *INFN, Sezione di Perugia, I-06123 Perugia, Italy*
- ⁴⁴ *Syracuse University, Syracuse, NY 13244, USA*
- ⁴⁵ *University of Minnesota, Minneapolis, MN 55455, USA*
- ⁴⁶ *SUPA, University of Glasgow, Glasgow G12 8QQ, United Kingdom*
- ⁴⁷ *LIGO Hanford Observatory, Richland, WA 99352, USA*
- ⁴⁸ *Caltech CaRT, Pasadena, CA 91125, USA*
- ⁴⁹ *Wigner RCP, RMKI, H-1121 Budapest, Konkoly Thege Miklós út 29-33, Hungary*
- ⁵⁰ *Columbia University, New York, NY 10027, USA*
- ⁵¹ *Stanford University, Stanford, CA 94305, USA*
- ⁵² *Università di Camerino, Dipartimento di Fisica, I-62032 Camerino, Italy*
- ⁵³ *Università di Padova, Dipartimento di Fisica e Astronomia, I-35131 Padova, Italy*
- ⁵⁴ *INFN, Sezione di Padova, I-35131 Padova, Italy*
- ⁵⁵ *Institute of Physics, Eötvös University, Pázmány P. s. 1/A, Budapest 1117, Hungary*
- ⁵⁶ *Nicolaus Copernicus Astronomical Center, Polish Academy of Sciences, 00-716, Warsaw, Poland*
- ⁵⁷ *Rochester Institute of Technology, Rochester, NY 14623, USA*
- ⁵⁸ *University of Birmingham, Birmingham B15 2TT, United Kingdom*
- ⁵⁹ *INFN, Sezione di Genova, I-16146 Genova, Italy*
- ⁶⁰ *RRCAT, Indore MP 452013, India*
- ⁶¹ *Faculty of Physics, Lomonosov Moscow State University, Moscow 119991, Russia*

-
- ⁶² SUPA, University of Strathclyde, Glasgow G1 1XQ, United Kingdom
- ⁶³ The Pennsylvania State University, University Park, PA 16802, USA
- ⁶⁴ OzGrav, University of Western Australia, Crawley, Western Australia 6009, Australia
- ⁶⁵ Department of Astrophysics/IMAPP, Radboud University Nijmegen, P.O. Box 9010, 6500 GL Nijmegen, The Netherlands
- ⁶⁶ Artemis, Université Côte d'Azur, Observatoire Côte d'Azur, CNRS, CS 34229, F-06304 Nice Cedex 4, France
- ⁶⁷ Institut FOTON, CNRS, Université de Rennes 1, F-35042 Rennes, France
- ⁶⁸ Washington State University, Pullman, WA 99164, USA
- ⁶⁹ University of Oregon, Eugene, OR 97403, USA
- ⁷⁰ Laboratoire Kastler Brossel, UPMC-Sorbonne Universités, CNRS, ENS-PSL Research University, Collège de France, F-75005 Paris, France
- ⁷¹ Carleton College, Northfield, MN 55057, USA
- ⁷² OzGrav, University of Adelaide, Adelaide, South Australia 5005, Australia
- ⁷³ Astronomical Observatory Warsaw University, 00-478 Warsaw, Poland
- ⁷⁴ VU University Amsterdam, 1081 HV Amsterdam, The Netherlands
- ⁷⁵ University of Maryland, College Park, MD 20742, USA
- ⁷⁶ School of Physics, Georgia Institute of Technology, Atlanta, GA 30332, USA
- ⁷⁷ Université Claude Bernard Lyon 1, F-69622 Villeurbanne, France
- ⁷⁸ Università di Napoli 'Federico II,' Complesso Universitario di Monte S. Angelo, I-80126 Napoli, Italy
- ⁷⁹ NASA Goddard Space Flight Center, Greenbelt, MD 20771, USA
- ⁸⁰ Dipartimento di Fisica, Università degli Studi di Genova, I-16146 Genova, Italy
- ⁸¹ RESCEU, University of Tokyo, Tokyo, 113-0033, Japan.
- ⁸² Tsinghua University, Beijing 100084, China
- ⁸³ Texas Tech University, Lubbock, TX 79409, USA
- ⁸⁴ Kenyon College, Gambier, OH 43022, USA
- ⁸⁵ Departamento de Astronomía y Astrofísica, Universitat de València, E-46100 Burjassot, València, Spain
- ⁸⁶ Museo Storico della Fisica e Centro Studi e Ricerche Enrico Fermi, I-00184 Roma, Italy
- ⁸⁷ National Tsing Hua University, Hsinchu City, 30013 Taiwan, Republic of China
- ⁸⁸ Charles Sturt University, Wagga Wagga, New South Wales 2678, Australia
- ⁸⁹ Center for Interdisciplinary Exploration & Research in Astrophysics (CIERA), Northwestern University, Evanston, IL 60208, USA
- ⁹⁰ Canadian Institute for Theoretical Astrophysics, University of Toronto, Toronto, Ontario M5S 3H8, Canada
- ⁹¹ University of Chicago, Chicago, IL 60637, USA
- ⁹² Pusan National University, Busan 46241, Korea
- ⁹³ The Chinese University of Hong Kong, Shatin, NT, Hong Kong
- ⁹⁴ INAF, Osservatorio Astronomico di Padova, I-35122 Padova, Italy
- ⁹⁵ INFN, Trento Institute for Fundamental Physics and Applications, I-38123 Povo, Trento, Italy
- ⁹⁶ OzGrav, University of Melbourne, Parkville, Victoria 3010, Australia
- ⁹⁷ Università di Roma 'La Sapienza,' I-00185 Roma, Italy
- ⁹⁸ Université Libre de Bruxelles, Brussels 1050, Belgium
- ⁹⁹ Sonoma State University, Rohnert Park, CA 94928, USA
- ¹⁰⁰ Departamento de Matemáticas, Universitat de València, E-46100 Burjassot, València, Spain
- ¹⁰¹ Montana State University, Bozeman, MT 59717, USA
- ¹⁰² Universitat de les Illes Balears, IAC3—IEEC, E-07122 Palma de Mallorca, Spain
- ¹⁰³ The University of Texas Rio Grande Valley, Brownsville, TX 78520, USA
- ¹⁰⁴ Bellevue College, Bellevue, WA 98007, USA
- ¹⁰⁵ Institute for Plasma Research, Bhat, Gandhinagar 382428, India
- ¹⁰⁶ The University of Sheffield, Sheffield S10 2TN, United Kingdom
- ¹⁰⁷ Dipartimento di Scienze Matematiche, Fisiche e Informatiche, Università di Parma, I-43124 Parma, Italy
- ¹⁰⁸ INFN, Sezione di Milano Bicocca, Gruppo Collegato di Parma, I-43124 Parma, Italy
- ¹⁰⁹ California State University, Los Angeles, 5151 State University Dr, Los Angeles, CA 90032, USA
- ¹¹⁰ Università di Trento, Dipartimento di Fisica, I-38123 Povo, Trento, Italy
- ¹¹¹ Montclair State University, Montclair, NJ 07043, USA
- ¹¹² National Astronomical Observatory of Japan, 2-21-1 Osawa, Mitaka, Tokyo 181-8588, Japan
- ¹¹³ Observatori Astronòmic, Universitat de València, E-46980 Paterna, València, Spain
- ¹¹⁴ School of Mathematics, University of Edinburgh, Edinburgh EH9 3FD, United Kingdom

-
- ¹¹⁵ *University and Institute of Advanced Research, Koba Institutional Area, Gandhinagar Gujarat 382007, India*
¹¹⁶ *IISER-TVM, CET Campus, Trivandrum Kerala 695016, India*
¹¹⁷ *University of Szeged, Dóm tér 9, Szeged 6720, Hungary*
¹¹⁸ *University of Michigan, Ann Arbor, MI 48109, USA*
¹¹⁹ *Tata Institute of Fundamental Research, Mumbai 400005, India*
¹²⁰ *INAF, Osservatorio Astronomico di Capodimonte, I-80131, Napoli, Italy*
¹²¹ *Università degli Studi di Urbino ‘Carlo Bo,’ I-61029 Urbino, Italy*
¹²² *INFN, Sezione di Firenze, I-50019 Sesto Fiorentino, Firenze, Italy*
¹²³ *Physik-Institut, University of Zurich, Winterthurerstrasse 190, 8057 Zurich, Switzerland*
¹²⁴ *American University, Washington, D.C. 20016, USA*
¹²⁵ *University of Białystok, 15-424 Białystok, Poland*
¹²⁶ *University of Southampton, Southampton SO17 1BJ, United Kingdom*
¹²⁷ *University of Washington Bothell, 18115 Campus Way NE, Bothell, WA 98011, USA*
¹²⁸ *Institute of Applied Physics, Nizhny Novgorod, 603950, Russia*
¹²⁹ *Korea Astronomy and Space Science Institute, Daejeon 34055, Korea*
¹³⁰ *Inje University Gimhae, South Gyeongsang 50834, Korea*
¹³¹ *National Institute for Mathematical Sciences, Daejeon 34047, Korea*
¹³² *NCBJ, 05-400 Świerk-Otwock, Poland*
¹³³ *Institute of Mathematics, Polish Academy of Sciences, 00656 Warsaw, Poland*
¹³⁴ *Hillsdale College, Hillsdale, MI 49242, USA*
¹³⁵ *Hanyang University, Seoul 04763, Korea*
¹³⁶ *Seoul National University, Seoul 08826, Korea*
¹³⁷ *NASA Marshall Space Flight Center, Huntsville, AL 35811, USA*
¹³⁸ *ESPCI, CNRS, F-75005 Paris, France*
¹³⁹ *Southern University and A&M College, Baton Rouge, LA 70813, USA*
¹⁴⁰ *College of William and Mary, Williamsburg, VA 23187, USA*
¹⁴¹ *Centre Scientifique de Monaco, 8 quai Antoine 1er, MC-98000, Monaco*
¹⁴² *Indian Institute of Technology Madras, Chennai 600036, India*
¹⁴³ *IISER-Kolkata, Mohanpur, West Bengal 741252, India*
¹⁴⁴ *Whitman College, 345 Boyer Avenue, Walla Walla, WA 99362 USA*
¹⁴⁵ *Indian Institute of Technology Bombay, Powai, Mumbai, Maharashtra 400076, India*
¹⁴⁶ *Scuola Normale Superiore, Piazza dei Cavalieri 7, I-56126 Pisa, Italy*
¹⁴⁷ *Université de Lyon, F-69361 Lyon, France*
¹⁴⁸ *Hobart and William Smith Colleges, Geneva, NY 14456, USA*
¹⁴⁹ *OzGrav, Swinburne University of Technology, Hawthorn VIC 3122, Australia*
¹⁵⁰ *Janusz Gil Institute of Astronomy, University of Zielona Góra, 65-265 Zielona Góra, Poland*
¹⁵¹ *University of Washington, Seattle, WA 98195, USA*
¹⁵² *King’s College London, University of London, London WC2R 2LS, United Kingdom*
¹⁵³ *Indian Institute of Technology, Gandhinagar Ahmedabad Gujarat 382424, India*
¹⁵⁴ *Indian Institute of Technology Hyderabad, Sangareddy, Khandi, Telangana 502285, India*
¹⁵⁵ *International Institute of Physics, Universidade Federal do Rio Grande do Norte, Natal RN 59078-970, Brazil*
¹⁵⁶ *Andrews University, Berrien Springs, MI 49104, USA*
¹⁵⁷ *Università di Siena, I-53100 Siena, Italy*
¹⁵⁸ *Trinity University, San Antonio, TX 78212, USA*
¹⁵⁹ *Van Swinderen Institute for Particle Physics and Gravity, University of Groningen, Nijenborgh 4, 9747 AG Groningen, The Netherlands*
¹⁶⁰ *Abilene Christian University, Abilene, TX 79699, USA*
¹⁶¹ *Colorado State University, Fort Collins, CO 80523, USA*

ABSTRACT

On June 8, 2017 at 02:01:16.49 UTC, a gravitational-wave signal from the merger of two stellar-mass black holes was observed by the two Advanced LIGO detectors with a network signal-to-noise ratio of 13. This system is the lightest black hole binary so far observed, with component masses $12^{+7}_{-2} M_{\odot}$ and $7^{+2}_{-2} M_{\odot}$ (90% credible intervals). These lie in the range of measured black hole masses in low-mass X-ray binaries, thus allowing us to compare black holes detected through gravitational waves with electromagnetic observations. The source’s luminosity distance is 340^{+140}_{-140} Mpc,

corresponding to redshift $0.07_{-0.03}^{+0.03}$. We verify that the signal waveform is consistent with the predictions of general relativity.

* Deceased, February 2017.

† Deceased, December 2016.

1. INTRODUCTION

The first detections of binary black hole mergers were made by the Advanced Laser Interferometer Gravitational-Wave Observatory (LIGO) (Aasi et al. 2015; Abbott et al. 2016a) during its first observing run (O1) in 2015 (Abbott et al. 2016b,c,d). Following a commissioning break, LIGO undertook a second observing run (O2) from November 30, 2016 to August 25, 2017, with the Advanced Virgo detector (Acernese et al. 2015) joining the run on August 1, 2017. Two binary black hole mergers (Abbott et al. 2017a,b) and one binary neutron star merger (Abbott et al. 2017c) have been reported in O2 data. Here we describe GW170608, a binary black hole merger with likely the lowest mass of any so far observed by LIGO.

GW170608 was first identified in data from the LIGO Livingston Observatory (LLO), which was in normal observing mode. The LIGO Hanford Observatory (LHO) was operating stably with sensitivity typical for O2, but its data were not analyzed automatically as the detector was undergoing a routine angular control procedure (Section 2 and Appendix A). Matched-filter analysis of a segment of data around this time revealed a candidate with source parameters consistent between both LIGO detectors; further offline analyses of a longer period of data confirmed the presence of a gravitational wave (GW) signal from the coalescence of a binary black hole system, with high statistical significance (Section 3).

The source’s parameters were estimated via coherent Bayesian analysis (Veitch et al. 2015; Abbott et al. 2016e). A degeneracy between the component masses m_1, m_2 prevents precise determination of their individual values, but the chirp mass $\mathcal{M} = (m_1 m_2)^{3/5} (m_1 + m_2)^{-1/5}$ is well measured and is the smallest so far observed for a merging black hole binary system, with the total mass $M = m_1 + m_2$ also likely the lowest so far observed (Section 4). Individual black hole spins are poorly constrained, however we find a slight preference for a small positive net component of spin in the direction of the binary orbital angular momentum.

In combination with GW151226 (Abbott et al. 2016c), this system points to a population of black hole binaries with component masses comparable to those of black holes found in X-ray binaries (Section 5) and significantly below those seen in other LIGO-Virgo black hole binaries.

We also test the consistency of the observed GW signal with the predictions of general relativity (GR); we find no deviations from those predictions.

2. DETECTOR OPERATION

The LIGO detectors measure gravitational-wave strain using dual-recycled Michelson interferometers with Fabry-Perot arm cavities (Aasi et al. 2015; Abbott et al. 2016a). During O2, the horizon distance for systems with component masses similar to GW170608—the distance at which a binary merger optimally oriented with respect to a detector has an expected signal-to-noise ratio (SNR) of 8 (Allen et al. 2012; Chen et al. 2017)—peaked at ~ 1 Gpc for LLO, and at ~ 750 Mpc for LHO.

At the time of GW170608, LLO was observing with a sensitivity close to its peak. LHO was operating in a stable configuration with a sensitivity of ~ 650 Mpc; a routine procedure to minimize angular noise coupling to the strain measurement was being performed (Kasprzack & Yu 2016). Although such times are in general not included in searches, it was determined that LHO strain data were unaffected by the procedure at frequencies above 30 Hz, and may thus be used to identify a GW source and measure its properties. More details on LHO data are given in Appendix A.

Similar procedures to those used in verifying previous GW detections (Abbott et al. 2017b) were followed and indicate that no disturbance registered by LIGO instrumental or environmental sensors (Effler et al. 2015) was strong enough to have caused the GW170608 signal.

Calibration of the LIGO detectors is performed by inducing test-mass motion using photon pressure from modulated auxiliary lasers (Karki et al. 2016; Abbott et al. 2017d; Cahillane et al. in preparation). The maximum $1\text{-}\sigma$ calibration uncertainties for strain data used in this analysis are 5% in amplitude and 3° in phase over the frequency range 20–1024 Hz.

The Advanced Virgo detector was, at the time of the event, in observation mode with a horizon distance for signals comparable to GW170608 of 60–70 Mpc. This was however during an early commissioning phase with still limited sensitivity, therefore Virgo data are not included in the analyses presented here.

3. SEARCH FOR BINARY MERGER SIGNALS

3.1. Low latency identification of a candidate event

GW170608 was first identified as a loud (SNR ~ 9) event in LLO data, via visual inspection of single-detector events from a low-latency compact binary matched filter (‘template’) analysis (Usman et al. 2016; Nitz et al. 2017b,a). Such events are displayed automatically to diagnose changes in detector operation and in populations of non-Gaussian transient noise artifacts (glitches) (Abbott et al. 2016f). Low-latency templated searches (Cannon et al. 2015; Messick et al. 2017; Adams et al. 2016; Nitz et al. 2017b) did not detect the event

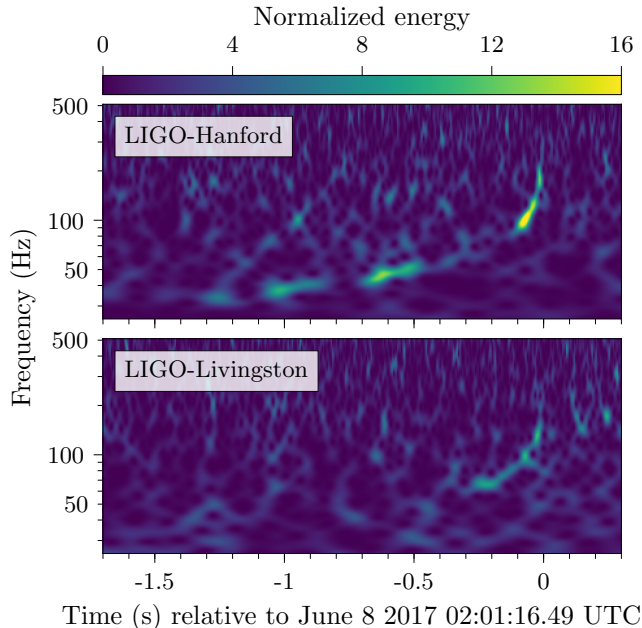


Figure 1. Time-frequency power maps of LIGO strain data at the time of GW170608. The characteristic upwards-chirping morphology of a binary inspiral driven by GW emission is visible in both detectors, with a higher signal amplitude in LHO. This figure, and all others in this Letter, were produced from noise-subtracted data (Section 4).

with high significance because LHO data were not analysed automatically. The morphology of the LLO event is consistent with a compact binary merger signal, as shown in Figure 1 (lower panel), but a noise origin could not be ruled out using LLO data alone.

Consequently, LHO data were investigated and were determined to be stable at frequencies above 30 Hz (see Appendix A). A segment of LHO data around the event time was then searched with a filter starting frequency of 30 Hz, using templates approximating the waveforms from compact binary systems with component spins aligned with the orbital angular momentum (Bohé et al. 2017; Pürrer 2016). The fraction of SNR expected to be lost due to imposing the 30 Hz cutoff, as compared to the lower starting frequencies typically used in O2 data (Dal Canton & Harry 2017), is $\sim 1\%$ or less. An event was found having consistent template binary masses and spins, times of arrival and SNRs in LHO and LLO. Based on this 2-detector coincident event an alert was issued to electromagnetic observing partners 13.5 hours after the event time, with a sky localization (Singer & Price 2016) covering 860 deg^2 (90% credible region). GRB Coordinates Network Circulars related to this event are archived at <https://gcn.gsfc.nasa.gov/other/G288732.gcn3>.

3.2. Offline search

To establish the significance of this coincident event, a period between June 7, 2017 and June 9, 2017 was identified for analysis during which both LIGO interferometers were operating in the same configuration as at the event time. Times at which commissioning activities at LHO produced severe or broad-band disturbances in the strain data were excluded from the analysis. Standard offline data quality vetoes for known environmental or instrumental artifacts were also applied, resulting overall in 1.2 days of coincident LHO–LLO data searched.

Two matched filter pipelines identified GW170608, with a network SNR of 13. An candidate event is assigned a ranking statistic value, in each pipeline, that represents its relative likelihood of originating from a GW signal *vs.* from noise. One pipeline estimates the noise background using time-shifted data (Usman et al. 2016), finding a rate of occurrence of noise events ranked higher than GW170608 of *less* than 1 in 3,000 years. This limit corresponds to the maximum background analysis time available from time shifts separated by 0.1 s. The other pipeline uses different methods for ranking candidate events and for estimating the background (Cannon et al. 2015; Messick et al. 2017) and assigns the event a false alarm rate of 1 in 160,000 years.

A search for transient GW signals coherent between LHO and LLO with frequency increasing over time, without using waveform templates (Klimenko et al. 2016), also identified GW170608 with a false alarm rate of 1 in ~ 30 years; the lower significance is expected as this analysis is typically less sensitive to lower-mass compact binary signals than matched filter searches.

4. SOURCE PROPERTIES

4.1. Binary Parameters

The parameters of the GW source are inferred from a coherent Bayesian analysis (Veitch et al. 2015; Abbott et al. 2016e) using noise-subtracted data from the two LIGO observatories. Noise subtraction is a data processing step that removes several instrumental noise sources from the GW strain measurements (Abbott et al. (2017b) and references therein), thus increasing the expected SNR of compact binary signals in LHO data by typically 25% (Driggers et al. 2017). The likelihood integration is performed starting at 30 Hz in LHO and 20 Hz in LLO, includes marginalization over strain calibration uncertainties (Farr et al. 2015), and uses the noise power spectral densities (Littenberg & Cornish 2015) at the time of the event.

Two different gravitational-wave signal models calibrated to numerical relativity simulations of general relativistic binary black hole mergers (Mroue et al. 2013;

Table 1. Source properties for GW170608: given are median values with 90% credible intervals. Source-frame masses are quoted; to convert to detector frame, multiply by $(1+z)$ (Krolak & Schutz 1987). The redshift assumes a flat cosmology with Hubble parameter $H_0 = 67.9 \text{ km s}^{-1} \text{ Mpc}^{-1}$ and matter density parameter $\Omega_m = 0.3065$ (Ade et al. 2016).

Chirp mass \mathcal{M}	$7.9^{+0.2}_{-0.2} M_\odot$
Total mass M	$19^{+5}_{-1} M_\odot$
Primary black hole mass m_1	$12^{+7}_{-2} M_\odot$
Secondary black hole mass m_2	$7^{+2}_{-2} M_\odot$
Mass ratio m_2/m_1	$0.6^{+0.3}_{-0.4}$
Effective inspiral spin parameter χ_{eff}	$0.07^{+0.23}_{-0.09}$
Final black hole mass M_f	$18.0^{+4.8}_{-0.9} M_\odot$
Final black hole spin a_f	$0.69^{+0.04}_{-0.05}$
Radiated energy E_{rad}	$0.85^{+0.07}_{-0.17} M_\odot c^2$
Peak luminosity ℓ_{peak}	$3.4^{+0.5}_{-1.6} \times 10^{56} \text{ erg s}^{-1}$
Luminosity distance D_L	$340^{+140}_{-140} \text{ Mpc}$
Source redshift z	$0.07^{+0.03}_{-0.03}$

Chu et al. 2016; Husa et al. 2016), building on the breakthrough reported in (Pretorius 2005; Campanelli et al. 2006; Baker et al. 2006), are used. One waveform family models the inspiral-merger-ringdown signal of precessing binary black holes (Hannam et al. 2014), which includes spin-induced orbital precession through a transformation of the aligned-spin waveform model of (Husa et al. 2016; Khan et al. 2016); we refer to this model as the *effective precession* model. The other waveform model describes binaries with spin angular momenta aligned with the orbital angular momentum (Bohé et al. 2017; Pürrer 2016), henceforth referred to as *non-precessing*. For their common parameters, both waveform models yield consistent parameter ranges.

A selection of inferred source parameters for GW170608 is given in Table 1; unless otherwise noted, we report median values and symmetric 90% credible intervals. The quoted parameter uncertainties include statistical and systematic errors from averaging posterior probability samples over the two waveform models. As in Abbott et al. (2017a), our estimates of the mass and spin of the final black hole, the total energy radiated in GWs as well as the peak luminosity are computed from fits to numerical relativity simulations (Hofmann et al. 2016; Keitel et al. 2017; Healy & Lousto 2017; Jiménez-Forteza et al. 2017).

The posterior probability distributions for the source-frame mass parameters of GW170608 are shown in Figure 2, together with those for GW151226 (Abbott et al.

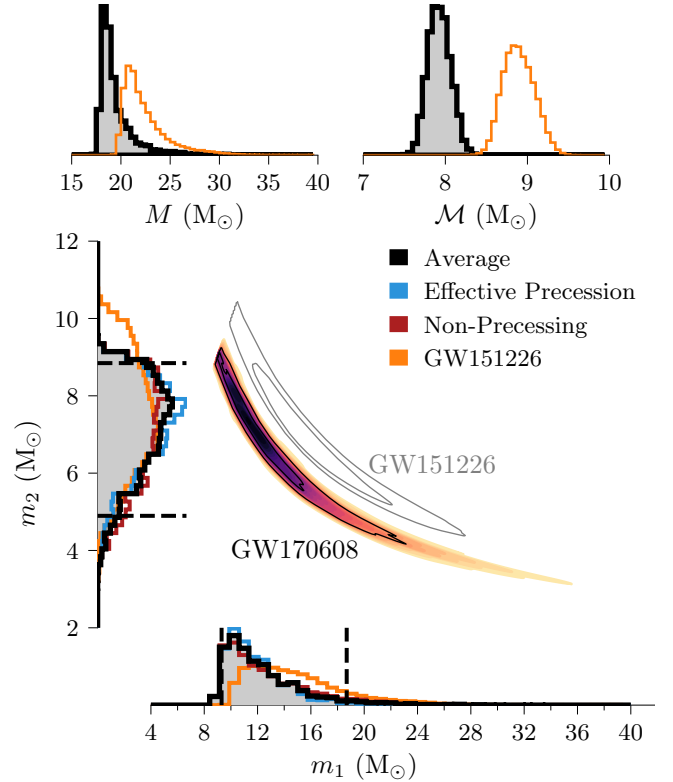


Figure 2. Posterior probability densities for binary component masses (m_1 , m_2), total mass (M), and chirp mass (\mathcal{M}) in the source frame. One-dimensional component mass distributions include posteriors for the effective precession (blue) and the non-precessing (red) waveform model, as well as their average (black). The dashed lines demarcate the 90% credible intervals for the average posterior. The two-dimensional plot shows contours of the 50% and 90% credible regions overlaid on a color-coded posterior density function. For comparison, we show both one- and two-dimensional distributions of averaged component mass posterior samples for GW151226 (orange) (Abbott et al. 2016c). In the top panel, we further compare GW170608 and GW151226's source-frame total mass (left) and source-frame chirp mass (right). All other known binary black holes lie at higher chirp masses than GW170608 and GW151226.

2016c). The initial binary of GW170608 consisted of two compact objects with source-frame component masses $m_1 = 12^{+7}_{-2} M_\odot$ and $m_2 = 7^{+2}_{-2} M_\odot$, with the mass ratio loosely constrained to $m_2/m_1 = 0.6^{+0.3}_{-0.4}$. Since neutron stars are expected to have masses below $\sim 4M_\odot$ (Lattimer & Prakash 2016), both objects are most likely black holes. Notably, we find this binary black hole system to be the least massive yet observed through gravitational waves. The next lightest, GW151226 (Abbott et al. 2016c), has a chirp mass $\mathcal{M} = 8.9^{+0.3}_{-0.3}$ and a total mass $M = 21.8^{+5.9}_{-1.7}$, compared with values of $\mathcal{M} = 7.9^{+0.2}_{-0.2} M_\odot$ and $M = 19^{+5}_{-1} M_\odot$ for GW170608.

The probability that GW170608’s total mass is smaller than GW151226’s is 0.89.

While the chirp mass is tightly constrained, spins have a more subtle effect on the GW signal. The effective inspiral spin χ_{eff} , a mass-weighted combination of the spin components (anti-)aligned with the orbital angular momentum (Racine 2008; Ajith et al. 2011), predominantly affects the inspiral rate of the binary but also influences the merger. We infer that $\chi_{\text{eff}} = 0.07^{+0.23}_{-0.09}$ disfavoring large, anti-aligned spins on both black holes.

An independent parameter estimation method comparing LIGO strain data to hybridized numerical relativity simulations of binary black hole systems with non-precessing spins (Abbott et al. 2016g) yields estimates of component masses and χ_{eff} consistent with our model-waveform analysis.

Spin components orthogonal to the orbital angular momentum are the source of precession (Apostolatos et al. 1994; Kidder 1995), and may be parameterized by a single effective precession spin χ_p (Schmidt et al. 2015). For precessing binaries, component spin orientations evolve over time; we report results evolved to a reference GW frequency of 20 Hz. The spin prior assumed in this analysis is uniform in dimensionless spin magnitudes $\chi_i \equiv c|\mathbf{S}_i|/(Gm_i^2)$ with $i = 1, 2$ between 0 and 0.89, and isotropic in their orientation; this prior on component spins maps to priors for the effective parameters χ_{eff} and χ_p . The top panel of Figure 3 shows the prior and posterior probability distributions of χ_{eff} and χ_p obtained for the effective-precession waveform model. While we gain some information about χ_{eff} , the χ_p posterior is dominated by its prior, as for previous GW events (Abbott et al. 2016b,c, 2017a), indicating that we cannot draw any strong conclusion on the size of spin components in the orbital plane (Vitale et al. 2017). The inferred component spin magnitudes and orientations are shown in the bottom panel of Figure 3. We find the dimensionless spin magnitude of the primary black hole, χ_1 , to be less than 0.75 (90% credible limit); this limit is robust to extending the prior range of spin magnitudes and to using different waveform models.

The measurability of precession depends on the intrinsic source properties as well as the angle of the binary orbital angular momentum to the line of sight (i.e. inclination). The inclination of GW170608’s orbit is likely close to either 0° or 180° , due to a selection effect: the distance inside which a given binary merger would be detectable at a fixed SNR threshold is largest for these inclination values (Schutz 2011). For such values, the waveform carries little information on precession.

The distance of GW170608 is extracted from the observed signal amplitude given the binary’s inclina-

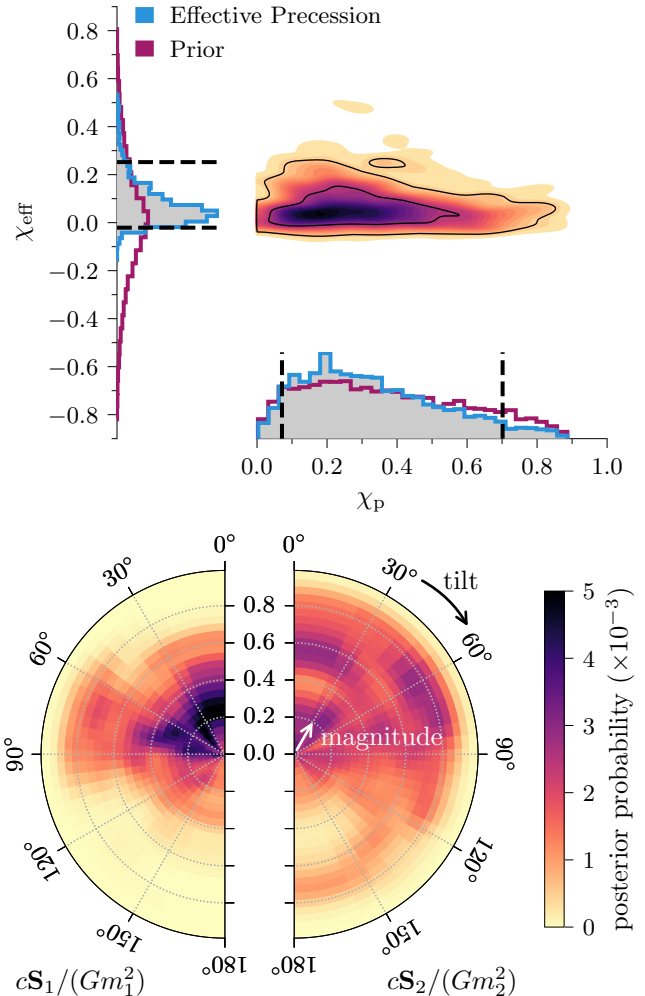


Figure 3. *Top panel:* Marginalized one-dimensional posterior density functions for the spin parameters χ_p and χ_{eff} (blue) in comparison to their prior distributions (pink) as obtained from the effective-precession model. The dashed lines indicate the 90% credible interval. The two-dimensional plot shows the 50% and 90% credible regions plotted over the posterior density function. *Bottom panel:* Posterior probabilities for the dimensionless component spins χ_i with $i = 1, 2$ relative to the Newtonian orbital angular momentum $\hat{\mathbf{L}}$, i.e. the normal of the orbital plane. The tilt angles are 0° for spins parallel to $\hat{\mathbf{L}}$, and 180° for spins anti-parallel to $\hat{\mathbf{L}}$. The posterior density functions are marginalized over the azimuthal angles. Each pixel has a prior probability of $\sim 1.8 \times 10^{-3}$; they are spaced linearly in spin magnitudes and the cosine of the tilt angles.

tion (Abbott et al. 2016e). With the network of two nearly co-aligned LIGO detectors, the uncertainty on inclination translates into a large distance uncertainty: we infer a luminosity distance of $D_L = 340^{+140}_{-140}$ Mpc, corresponding to a redshift of $z = 0.07^{+0.03}_{-0.03}$ assuming a flat Λ CDM cosmology (Ade et al. 2016).

GW170608 is localized to a sky area of ~ 520 deg² in the Northern hemisphere (90% credible region), determined largely by the signal’s measured arrival time at LLO ~ 7 ms later than at LHO. This reduction in area relative to the low-latency map is partly attributable to the use of noise-subtracted data with offline calibration (Abbott et al. 2017b).

4.2. Consistency with General Relativity

To test whether GW170608 is consistent with the predictions of GR, we consider possible deviations of coefficients describing the binary inspiral part of the signal waveform from the values expected in GR, as was done for previous detections (Abbott et al. 2016h,d, 2017a). Tests involving parameters describing the merger and ringdown do not yield informative results, since the merger happens at relatively high frequency where the LIGO detectors are less sensitive. As in Abbott et al. (2017b), we also allow a sub-leading phase contribution at effective -1 PN order, i.e. with a frequency dependence of $f^{-7/3}$, which is absent in GR. The GR predicted value is contained within the 90% credible interval of the posterior distribution for all parameters tested.

Assuming that gravitons are dispersed in vacuum like massive particles, we also obtained an upper bound on the mass of the graviton comparable to the constraints previously obtained (Abbott et al. 2016b,h, 2017a). Possible violations of local Lorentz invariance, manifested via modifications to the GW dispersion relation, were investigated (Abbott et al. 2017a), again finding upper bounds comparable to previous results.

5. ASTROPHYSICAL IMPLICATIONS

The low mass of GW170608’s source binary, in comparison to other binary black hole systems observed by LIGO and Virgo, has potential implications for the binary’s progenitor environment. High-metallicity progenitors are expected to experience substantial mass loss through strong stellar winds (Spera et al. 2015), implying that high-mass black hole binaries are not formed in such environments. GW170608’s low mass suggests formation in a higher metallicity environment; however, formation at lower metallicity with comparatively lower mass progenitors is not excluded. Further discussion of the relationship between black hole masses and metallicity can be found in Abbott et al. (2016i).

At this lower boundary of the observed distribution of binary black hole masses, we can compare component black holes with those found in X-ray binaries. X-ray binary systems contain either a black hole or neutron star which accretes matter from a companion donor star. Low-mass X-ray binaries (LMXBs) are X-ray binaries with a low-mass donor star which transfer mass

through Roche lobe overflow (Charles & Coe 2003). The inferred component masses of GW170608 are consistent with dynamically-measured masses of black holes found in LMXBs, typically less than $10M_{\odot}$ (Özel et al. 2010; Farr et al. 2011; Corral-Santana et al. 2016).

Black holes in LMXBs are believed to form with near-zero spin and acquire spin as a byproduct of mass accretion (Fragos & McClintock 2015). For GW170608, we infer an effective spin probability distribution that is concentrated around zero with the 90% credible interval extending to small positive spin. Thus we can exclude highly negative anti-parallel spin components, while remaining broadly consistent with expected LMXB spins (Miller & Miller 2015; Fragos & McClintock 2015).

Binary black holes may form through many different channels, including, but not limited to, dynamical interaction (Rodríguez et al. 2016; O’Leary et al. 2016; Mapelli 2016) and isolated binary evolution (Belczynski et al. 2016; Eldridge & Stanway 2016; Lipunov et al. 2017; Stevenson et al. 2017b). While the inferred masses and tilt measurements of GW170608 are not sufficiently constrained to favor a formation channel, future measurements of binary black hole systems may hint at the formation histories of such systems (see Abbott et al. (2017a) and references therein). It may be possible to determine the relative proportion of binaries originating in each canonical formation channel following $\mathcal{O}(100)$ binary black hole detections (Vitale et al. 2017; Stevenson et al. 2017a; Zevin et al. 2017; Talbot & Thrane 2017; Farr et al. 2017a,b).

The detection of GW170608 is consistent with the merger populations considered in Abbott et al. (2016j,d) for which a rate of $12\text{--}213$ Gpc⁻³ yr⁻¹ was estimated in Abbott et al. (2017a).

6. OUTLOOK

LIGO’s detection of GW170608 extends the range of known stellar-mass binary black hole systems at the low-mass boundary, and hints at connections with other known astrophysical systems containing black holes. The O2 run ended on August 25th, 2017; a full catalog of binary merger gravitational-wave events for this run is in preparation, including candidate signals with lower significance and systems other than stellar-mass black hole binaries (Abbott et al. 2017c). Estimates of the merger rate and mass distribution for the emerging compact binary population will also be updated.

With expected increases in detector sensitivity in the third advanced detector network observing run, projected for late 2018 (Abbott et al. 2016k), detection of black hole binaries will be a routine occurrence; studying

this population will eventually answer many questions about these systems' origins and evolution.

The authors gratefully acknowledge the support of the United States National Science Foundation (NSF) for the construction and operation of the LIGO Laboratory and Advanced LIGO as well as the Science and Technology Facilities Council (STFC) of the United Kingdom, the Max-Planck-Society (MPS), and the State of Niedersachsen/Germany for support of the construction of Advanced LIGO and construction and operation of the GEO600 detector. Additional support for Advanced LIGO was provided by the Australian Research Council. The authors gratefully acknowledge the Italian Istituto Nazionale di Fisica Nucleare (INFN), the French Centre National de la Recherche Scientifique (CNRS) and the Foundation for Fundamental Research on Matter supported by the Netherlands Organisation for Scientific Research, for the construction and operation of the Virgo detector and the creation and support of the EGO consortium. The authors also gratefully acknowledge research support from these agencies as well as by the Council of Scientific and Industrial Research of India, the Department of Science and Technology, India, the Science & Engineering Research Board (SERB), India, the Ministry of Human Resource Development, India, the Spanish Agencia Estatal de Investigación, the Vicepresidència i Conselleria d'Innovació, Recerca i Tur-

isme and the Conselleria d'Educació i Universitat del Govern de les Illes Balears, the Conselleria d'Educació, Investigació, Cultura i Esport de la Generalitat Valenciana, the National Science Centre of Poland, the Swiss National Science Foundation (SNSF), the Russian Foundation for Basic Research, the Russian Science Foundation, the European Commission, the European Regional Development Funds (ERDF), the Royal Society, the Scottish Funding Council, the Scottish Universities Physics Alliance, the Hungarian Scientific Research Fund (OTKA), the Lyon Institute of Origins (LIO), the National Research, Development and Innovation Office Hungary (NKFI), the National Research Foundation of Korea, Industry Canada and the Province of Ontario through the Ministry of Economic Development and Innovation, the Natural Science and Engineering Research Council Canada, the Canadian Institute for Advanced Research, the Brazilian Ministry of Science, Technology, Innovations, and Communications, the International Center for Theoretical Physics South American Institute for Fundamental Research (ICTP-SAIFR), the Research Grants Council of Hong Kong, the National Natural Science Foundation of China (NSFC), the Leverhulme Trust, the Research Corporation, the Ministry of Science and Technology (MOST), Taiwan and the Kavli Foundation. The authors gratefully acknowledge the support of the NSF, STFC, MPS, INFN, CNRS and the State of Niedersachsen/Germany for provision of computational resources.

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APPENDIX

A. ANGULAR COUPLING MINIMIZATION

GW170608 was observed during a routine instrumental procedure at LHO that minimises the coupling of angular control of the test masses to noise in the GW strain measurement. To maintain resonant power in the arms, the pitch and yaw angular degrees of freedom of the four suspended cavity test masses at each detector (Abbott et al. 2016a) must be controlled. This is achieved by actuating on the second stage of the LIGO quadruple suspensions. A feed-forward control is employed in order to leave the beam position of the main laser on the test mass unchanged while this actuation is applied. However, if this position differs from the actuation point, the angular control can affect the differential arm length, thus introducing additional noise in the strain measurement (Kasprzack & Yu 2016). As the beam position can drift over periods of hours or days, the angular feed-forward control must be periodically adjusted in order to minimize the coupling to strain.

During this procedure, high amplitude pitch and yaw excitations are applied to the test masses via actuation of the suspensions. Each of the 8 angular degrees of freedom is excited at a distinct frequency; the resulting length signals are observed via demodulation at each excitation frequency, revealing how strongly the corresponding degree of freedom couples to differential arm length. The feed-forward gain settings are stepped at intervals of approximately 45 s and the global minimum of angular control coupling to strain is determined from the resulting measurements. The frequencies of angular excitations are equally spaced between ~ 19 Hz and ~ 23 Hz, generating excess power in the differential arm motion, and thus in the measured strain around these frequencies. This procedure covers from ~ 2 minutes before to ~ 14 minutes after GW170608, shown in Figure 4 (left). During the period from -2 to 2 minutes substantial excess noise is visible at frequencies around 20 Hz. To characterize this noise we show amplitude spectral densities derived from 240 s of data both before the onset of the angular excitations and during the excitations around the event time in Figure 4 (right). No effect on the spectrum is visible above 30 Hz.

During the procedure, angular control gain settings are stepped abruptly; inspection of all such transition times shows no evidence for transient excess noise in the strain data outside the 19–23 Hz excitation band. The closest transition to the event time was 10 s before the binary merger, thus any transient noise associated with this transition

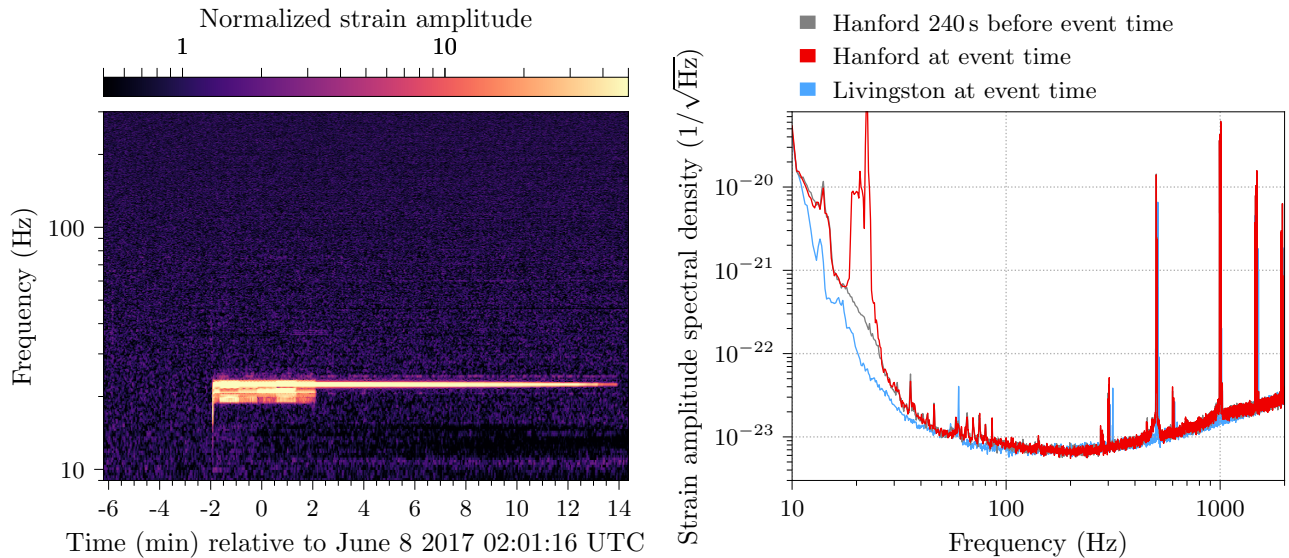


Figure 4. Left: Spectrogram of strain data from LHO around the time of GW170608. This plot shows variations in the noise spectrum of the detector over periods on the scale of minutes; unlike Figure 1, it is not designed to show short-duration transient events. The strain amplitude is normalized to the interval between -6 and -2 minutes relative to the event time. See Appendix A for discussion of the feature around 20 Hz due to an angular control procedure. Right: Amplitude spectral density of strain data at both LIGO observatories for 240 s around the event time, $(-2, 2)$ minutes on left panel, and for data before the start of the angular coupling minimization at LHO, $(-6, -2)$ minutes. Excess noise is clearly visible around 20 Hz but data above 30 Hz are unaffected.

could not have affected the matched-filter output at the event time (template waveforms for GW170608-like signals have a duration between 2 and 3 s.) Furthermore, events from a single-detector matched-filter search covering other periods at LHO when this procedure was performed shows no anomalous features compared to other times. Thus, we find no evidence that the angular coupling minimization affected the recorded strain data at LHO around the event time at frequencies above 30 Hz.