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Spatially and Temporally Resolved Heterogeneities in a Miscible Polymer Blend

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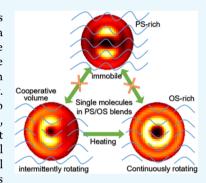
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ABSTRACT: Mapping the spatial and temporal heterogeneities in miscible polymer blends is critical for understanding and further improving their material properties. However, a complete picture on the heterogeneous dynamics is often obscured in ensemble measurements. Herein, the spatial and temporal heterogeneities in fully miscible polystyrene/oligostyrene blend films are investigated by monitoring the rotational diffusion of embedded individual probe molecules using defocused wide-field fluorescence microscopy. In the same blend film, three significantly different types of dynamical behaviors (referred to as modes) of the probe molecules can be observed at the same time, namely, immobile, continuously rotating, and intermittently rotating probe molecules. This reveals a prominent spatial heterogeneity in local dynamics at the nanometer scale. In addition to that, temporal heterogeneity is uncovered by the nonexponential characteristic of the rotational autocorrelation functions of single-molecule probes. Moreover, the occurrence probabilities



of these different modes strongly depend on the polystyrene: oligostyrene ratios in the blend films. Remarkably, some probe molecules switch between the continuous and intermittent rotational modes at elevated temperature, indicating a possible alteration in local dynamics that is triggered by the dynamic heterogeneity in the blends. Although some of these findings can be discussed by the self-concentration model and the results provided by ensemble averaging techniques (e.g., dielectric spectroscopy), there are implications that go beyond current models of blend dynamics.

■ INTRODUCTION

Mixing two or more polymers usually results in macro-phase separation. There exists, however, a small class of polymers, with not too different polarizabilities, that upon mixing, form thermodynamically miscible blends. The latter have received much attention in academia and, not unexpectantly, less attention in industry. 1-12 The physical properties of these blends are often linked to those of each individual component; however, it is hardly a linear sum-up, resulting in extra complexity in predicting the blend properties compared with their pure polymer counterparts.^{6–12} For instance, the glass temperature (T_g) is an important parameter for evaluating the dynamics of polymers or polymer blends. It was traditionally believed that a miscible polymer blend possesses a single $T_{\rm g}$. However, both a broadened $T_{\rm g}^{13,14}$ and two calorimetric $T_{\varrho} {\rm s}^{15-18}$ or two separate segmental processes that reflect the component's segmental dynamics have been reported in miscible polymer blends. The presence of two distinct segmental dynamics in a fully miscible polymer blend is known by the term "dynamic heterogeneity". 1-9

The dynamic heterogeneity in a miscible polymer blend is controlled by the dynamic asymmetry, *i.e.*, the difference in the glass temperatures of the parent homopolymers. ^{1-9,13-21} Along these lines, Harmandaris et al. have reported a *perfectly* miscible blend of polystyrene and oligostyrene that possesses

large dynamic asymmetry as revealed both in dielectric spectroscopy measurements and in the molecular dynamics simulation. The blend shows a bimodal segmental relaxation process, implying dynamic heterogeneity. 5

Bulk measurements of the dynamics in homopolymers reveal the heterogeneous nature of the segmental dynamics as evidenced by the broadening of the structural relaxation on approaching $T_{\rm g}$ from higher temperatures. Nonetheless, with the exception of a few systems (poly(n-butylmethacrylate) (PnBMA)), it is unknown how the heterogeneity can be resolved because of spatial and temporal averaging in bulk measurements. Single-molecule spectroscopy and microscopy are able to provide information on structures and dynamics with high spatial and temporal resolutions while avoid ensemble averaging. Local dynamics of a polymer have been studied by following the rotational diffusion of single molecules embedded in a polymer matrix. The reorientation of these embedded single molecules is largely

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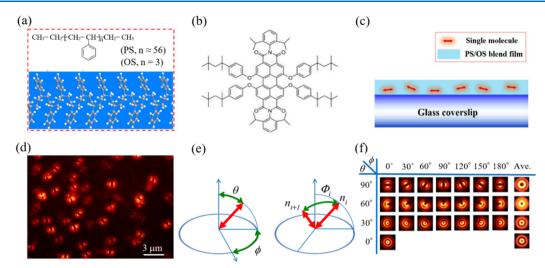


Figure 1. (a) Chemical structure of polystyrene (PS, $n \approx 56$) and oligostyrene (OS, n = 3). (b) Chemical structure of the perylene diimide derivative (PDI) dye molecule. (c) Schematic view of experimental samples. (d) Typical defocused image of single PDI molecules embedded in a PS/OS blend film with $\sim 1~\mu$ m defocusing toward the sample. (e) Transition dipole moment of a single PDI molecule presented by using in-plane (ϕ) and out-of-plane (θ) angles (left) and 3D angular displacement (Φ) (right). (f) Defocused patterns of single molecule for various ϕ and θ angles as well as an averaged image for each row.

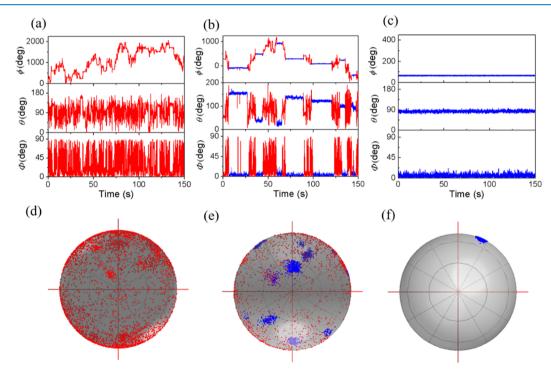


Figure 2. Three typical rotational diffusion behaviors of single PDI molecules coexisting in all PS/OS blend films with different compositions at 296–315 K. (a–c) Trajectories of ϕ , θ , and Φ as a function of time and (d–f) the corresponding 3D projection maps. These trajectories and 3D projection maps indicate the diffusion behaviors of a continuously rotating molecule (panels (a, d)), an intermittently rotating molecule (panels (b, e)), and an immobile molecule (panels (c, f)). The blue color is used to highlight the orientation states with a longer dwell time.

dependent on the dynamics of the surrounding polymer matrix and thus can be used to probe the structural (segmental) relaxation of the surrounding matrix at the nanometer scale. Particularly, the three-dimensional (3D) evolution of single-molecule orientations can be visualized using single-molecule defocused wide-field fluorescence microscopy (SMDWM) with a sub-micrometer spatial resolution, 30-34,36-38 providing a powerful tool for investigating spatial and temporal heterogeneities related to the glass temperature. 31,32

Returning to the issue of dynamic heterogeneity in miscible polymer blends, nothing is known about the extent of spatial and temporal heterogeneities related to the individual component dynamics. Given the mixing of segments within the relevant length scale (the Kuhn length according to the self-concentration model of Lodge and McLeish), ³⁹ the expectation is that spatial and temporal heterogeneities of the individual components will be further enhanced.

In this work, we investigated the spatial and temporal heterogeneities in thin blend films of polystyrene (PS) and

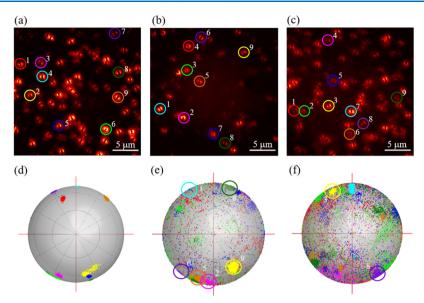


Figure 3. Typical defocused wide-field fluorescence images of single molecules (a) in the pure PS films, (b) in the blend films of 75 wt % PS, and (c) 25 wt % PS at a temperature of 296 K. These images are formed by accumulating 100 frames of respective defocused imaging sequences. (d-f) 3D projection maps of the orientation for the molecules circled in panels (a-c) over time. Orientations of each molecule every 50 ms are represented by dots of a single color. These dots circled on the 3D projection maps in panels (e, f) indicate immobile molecules, while all other dots indicate continuously rotating and intermittently rotating molecules.

oligostyrene (OS) that are dynamically heterogeneous by means of SMDWM. By acquiring time-sequenced defocused images, it was found that the single molecules embedded in the blend film exhibit clearly different dynamical behaviors, revealing spatial heterogeneity. In addition, the non-exponential characteristic of the rotational autocorrelation function of these single molecules from this study indicates temporal heterogeneity. To further examine the evolvement of the heterogeneity, a temperature-dependent investigation was carried out, revealing the temporal transitions between different types of dynamical behaviors of the probe molecules.

■ RESULTS AND DISCUSSION

PS/OS Blends and Defocused Imaging. The PS/OS blends used in this work are composed of PS ($M_{\rm w}=6140~{\rm g}~{\rm mol}^{-1}$, D=1.04) and oligomer OS ($M_{\rm w}=370~{\rm g}~{\rm mol}^{-1}$, D=1.00). The chemical structure of PS and OS is presented in Figure 1a. The two components are perfectly miscible but with very different $T_{\rm g}$ s (difference in glass temperatures, $\Delta T_{\rm g}=122~{\rm K}$). This induces a strong dynamical asymmetry in the blends. Blends were prepared in toluene with a PS to OS weight ratio of 75:25, 50:50, or 25:75. The perylene diimide derivative (PDI, chemical structure shown in Figure 1b) molecules were then added as optical nanoprobes for long-period single-molecule imaging. $^{30-32,40-42}$ This mixture was subsequently spin-coated onto a cleaned glass coverslip, obtaining a ~300 nm-thick PS/OS blend film (schematic image shown in Figure 1c). The details of the sample preparation are presented in the Experimental Section.

3D orientation information of single PDI molecules can be extracted from the defocused fluorescence images (a typical defocused image shown in Figure 1d). As illustrated in the left part of Figure 1e, the 3D orientation of a transition dipole can be described using in-plane (φ) and out-of-plane (θ) angles in a polar coordinate system. The 3D angular displacement (Φ) of a PDI molecule could be illustrated in this coordinate as shown in the right part of Figure 1e. Defocused patterns of a

molecular transition dipole with various ϕ and θ angles are numerically obtained according to the reported model (Figure 1f).³⁶ By comparing an acquired defocused pattern of a PDI molecule with the simulated ones (pattern matching), the 3D orientation of this molecule (ϕ and θ angles) can be determined (see the Experimental Section for details of defocused imaging analysis).

Dynamical Behavior of Single Molecules in PS/OS Blend Films. The 3D rotational motions of the embedded single PDI molecules can thus be monitored as function of time by analyzing the time sequenced defocused images, as illustrated in Figure 2. Three types of rotational diffusion behaviors were found in the blend films: continuous rotating, intermittent rotating, and immobilized, which coexisted in all blend films with different compositions and in a temperature range of 296-315 K. The rotational behavior can be well illustrated by the trajectories of ϕ , θ , and Φ as a function of time, as well as the corresponding 3D projection maps, in Figure 2. For a continuously rotating molecule, fast consecutive reorientation of this probe molecule took place with relatively large angular displacements and is barely interrupted (the lowest part in Figure 2a). By projecting all these molecular orientations on a 3D projection map, an almost complete coverage of the angular space was observed (Figure 2d). For an intermittently rotating molecule, it switched between the "slow" (locked orientation, blue traces) and the "fast" (large 3D angular displacement, red traces) dynamical states intermittently, as shown in Figure 2b,e. For an immobile molecule, it remained static in a single fixed orientation throughout the whole measurement of several hours, as shown in Figure 2c,f.

These immobile and continuously rotating dynamics observed in the blend film are similar to those in pure PS film and pure OS film under the same measurement conditions, respectively. As shown in Figure S1, all probe molecules in the bulk PS film kept immobile, while almost all probe molecules in pure OS film kept rotating continuously.

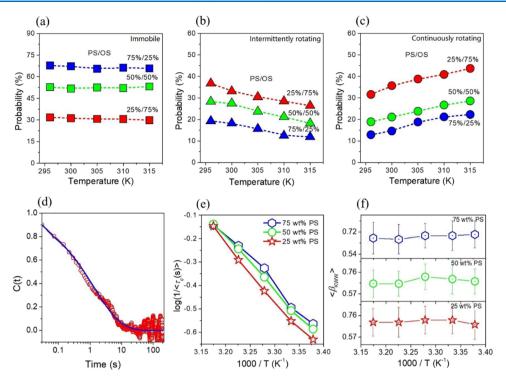


Figure 4. Statistics for the probabilities of (a) immobile molecules, (b) intermittently rotating molecules, and (c) continuously rotating molecules in the blend films of 75, 50, and 25 wt % PS at 296, 300, 305, 310, and 315 K. (d) Rotational autocorrelation function of a single molecule in the PS/OS blend films and a fitting with KWW stretched exponential function. (e) Arrhenius diagram for the PS/OS bend films with 75, 50, and 25 wt % PS. The average rotational correlation times, $\langle \tau_c \rangle$, were obtained from \sim 100 molecules in each blend composition at each temperature, respectively. (f) Calculated $\langle \beta_{\text{KWW}} \rangle$ parameters obtained from the fitting of KWW stretched exponential functions.

Note that the experiment temperature (296-315 K) is far below the $T_{\rm g}$ (~358 K) of PS and is far above the $T_{\rm g}$ (~236 K) of OS. Therefore, no segmental relaxation is present in the pure PS film while a homogeneous segmental relaxation is present in the pure OS film.³² The presence of bimodal relaxation times in the PS/OS miscible blend can be explained with the self-concentration model. 39,43-46 According to this model, the slow component forms a vitreous matrix with intricate "cavities" in which the fast component is still mobile. In addition, the local dynamics of the fast component in polymer blends resembles that in pure polymers.³⁹ It is reasonable to argue that the PS component (slow component) in the blend is frozen at the temperature of the measurement, whereas the OS component (fast component) remains highly mobile. In previous reports, the segmental relaxation of the fast component under the constraint of the slow component (more rigid) has been confirmed in a polymer blend using a combination of dielectric spectroscopy and nuclear magnetic resonance.⁴⁷ Therefore, these continuously rotating (fast dynamics) and immobile molecules (slow dynamics) are the natural consequences from the respective local segmental relaxations of OS and PS components in the surroundings of PDI molecules.

Spatial Heterogeneity in PS/OS Blend Films. The defocused images and corresponding 3D projection maps in Figure 3 demonstrate the typical spatial heterogeneity observed in single PS/OS blend films and illustrate how the reorientation behaviors of the probe molecules are dependent on the blend composition. The orientation evolutions of randomly picked probe molecules that are highlighted with circles of different colors in Figure 3a—c are plotted in the 3D orientation projection maps in Figure 3d—f with corresponding

colors. As a reference, the probe molecules embedded in a pure PS film remained mostly immobile, as shown in Figure 3a,d. Because of this immobility, the defocused patterns of the probe molecules remained almost static and resulted in closely clustered spots in the 3D-orientation projection for each of the probe molecules. In contrast, the probe molecules' reorientational behaviors in blend films are more dynamic and more diverse. Figure 3b,c shows the accumulated defocused images of randomly picked probe molecules in the blend films with 75 and 25 wt % PS. It is noteworthy that some defocused patterns (molecules 3, 4, and 7 in Figure 3b) appear with closed circular rings, similar to that in Figure 1f. Such patterns correspond to the probe molecules that are rotating continuously during the acquisition of an image frame, thus showing an averaged image of various orientations. It was also found that the spatial distribution of the continuously rotating, intermittently rotating, and immobile molecules is random across the whole blend films, indicating a prominent spatial heterogeneity. From the 3D projection maps shown in Figure 3e,f, it can also be noted that there are significantly more rotating molecules in the blend films with 25 wt % PS than that with 75 wt % PS. The statistical results on various rotating behaviors will be presented in the following paragraphs.

Figure 4a—c represents the composition- and temperature-dependencies of the appearance probabilities of immobile, intermittently rotating, and continuously rotating molecules. At a given blend composition and temperature, the statistical data are obtained by analyzing ~300 molecules. It was found that the occurrence probabilities of immobile molecules strongly depend on the blend composition rather than on temperature (Figure 4a). According to the self-concentration model,³⁹ the local composition around any chosen segment is

always richer in same species because of chain connectivity. Therefore, it is reasonable to believe that the immobile molecules reside in a local PS-rich segment, freezing the rotational mobility of probe molecules at 296-315 K. It is therefore not surprising to find a high percentage (\sim 67%) of the probe molecules being immobile in the blend film with 75 wt % PS, as the higher coverage of PS-rich regions is a logical consequence of the high PS content in the blend film. The intermittently rotating molecules can be explained with the self-concentration model of Lodge and McLeish. 39 Because the local compositions formed by OS-rich regions consist of a few PS segments, the PS segments in the local compositions affect the continuous rotation of probe molecules at the experimental temperature, which leads to the intermittently rotating behaviors.

In contrast, the occurrence probabilities of the continuously rotating and intermittently rotating molecules depend not only on the blend composition but also on the temperature (Figure 4b,c). Interestingly, these two modes can be switched when varying the temperature of the PS/OS blend. As the temperature increases, the occurrence probability of intermittently rotating molecules decreases and the occurrence probability of the continuously rotating one increases in the same sample area, indicating that some of the probe molecules switch from the intermittently rotating behavior to the continuously rotating one at the elevated temperature. The switching process can occur in the opposite direction by decreasing the temperature. This thermal-driven switch implies that these probe molecules locate in a "cooperative volume" that is formed by the segments from both components. In this cooperative volume, the segmental relaxation of OS component is more effective in driving the rotation of probe molecules at an elevated temperature, leading to a switching of probe molecules from intermittent rotating to continuous rotating. The temperature-dependent switching indicates a possible alteration in local dynamics that is triggered by the dynamic heterogeneity in the blends. For those molecules rotating continuously at 296 K, they may locate in the local compositions formed by OS-rich regions. Therefore, we can conclude that the observed spatial heterogeneity is coupled with the dynamic heterogeneity in PS/OS blends.

Temporal Heterogeneity in PS/OS Blend Films. To further investigate the local relaxation dynamic in the PS/OS blend films, the rotational correlation time (τ_c) of probe molecules was estimated by fitting rotational autocorrelation function with the Kohlrausch-Williams-Watts (KWW) stretched exponential function, $C(t) = \exp \left[- \left(t / \tau_{KWW} \right)^{\beta_{KWW}} \right]$ (see the Experimental Section for details) where τ_{KWW} is the decay constant, and $\beta_{\rm KWW}$ is the stretched exponent which describes the width relaxation time distribution. Figure 4d shows a typical autocorrelation curve of a rotating molecule, and τ_c was estimated to be 2.58 s ($\tau_{KWW} = 1.33$ s, $\beta_{KWW} = 0.51$) by using the KWW function. Tables S1-S3 summarize the averaged rotational correlation time ($\langle \tau_c \rangle$) of the molecules embedded in PS/OS blends with PS percentages of 75, 50, and 25 wt %, respectively. The $\langle \tau_c \rangle$ were determined by averaging the τ_c values obtained from nearly 100 molecules in a given blend composition and temperature (the histograms of τ_c are presented in Figure S2). All continuously rotating molecules and most intermittently rotating molecules were included in this fitting. The values of $\langle \tau_c \rangle$ show a clear temperature dependency, as depicted in the Arrhenius plot in Figure 4e. Note that the $\log(1/\langle \tau_c(s) \rangle)$ values in the blend films with

75 and 50 wt % PS are slightly longer than that with 25 wt % PS, indicating faster dynamics in the blend films with the higher PS content. The faster dynamics corresponds to the faster dynamics in the OS-rich segment (the fast component) that are confined by the high percentage PS component (slow component), in agreement with the literature. 7,44,48-5

The dependency of the average stretched exponent $(\langle \beta_{KWW} \rangle)$ on component and temperature is shown in Figure 4f. β_{KWW} describes the nonexponential behavior that is related to the temporal heterogeneity of polymer blends. A smaller β_{KWW} indicates a more pronounced temporal heterogeneity in single-molecule spectroscopy.³² Note that $\langle \beta_{KWW} \rangle$ in Figure 4f is almost independent on temperature, and the standard deviation of $\langle \beta_{\rm KWW} \rangle$ is large at all temperatures. These findings are similar to those observed in the poly(n-butyl methacrylate) homopolymer near $T_{\rm g}$. This temperatureindependent $\langle \beta_{KWW} \rangle$ suggests that the temporal heterogeneity is less pronounced than the spatial heterogeneity within the time scale of observation.³²

Discussion on the Difference between Single-Molecule Data and Dielectric Spectroscopy Data. To further understand the dynamics information in the PS/OS blends, the correlation times for the fast and slow segmental dynamics at the temperature range of 203-423 K were obtained by dielectric spectroscopy. Compared with the results measured with an ensemble averaging technique (dielectric spectroscopy, as presented in Figure S3, see Supporting Information for experimental details), the correlation times from SMDWM measurements are much longer and show weak dependency on the blend composition and temperature (Figure S4).

These differences can be explained by the following two arguments. The first relates to the fact that the rotational diffusions reflect the behaviors of all observed single molecules and the correlation times were obtained from only a small proportion of the molecules (the proportion was represented by N as shown in Tables S1–S3). For example, for the blend with 75% PS at 296 K, the correlation times were calculated from the rotating molecules, which accounts for only ~25% of all the probe molecules. The remaining molecules (\sim 75%) are immobile (67.5%, as shown in Figure 4a) or intermittently rotating (\sim 7.5%, due to that their autocorrelation function does not show a full relaxation), which were excluded in the calculation of correlation time. Therefore, the correlation times obtained with SMDWM reveal only part of the blend dynamics. The second argument relates to differences in the characteristic length scales as probed by the different techniques. In the miscible blends of PS/OS, the Kuhn lengths for PS and OS are 1.48 and ~0.65 nm, respectively. Based on the self-concentration model,³⁹ miscible blends can be envisioned as uniformly distributed rigid balls (mainly composed of PS) of 1.48 nm and fluid balls (mainly composed of OS) of ~0.65 nm. The size of PDI probe molecules (~2 nm) exceeds the characteristic lengths of both PS and OS. As we have shown earlier with respect to the PnBMA study³² with the combination of DS and SMDWM, the larger size of the probe molecule will sample differently the segmental dynamics and results to a longer time scale. This is in line with the observed longer correlation times in the SMDWM experiment.

Lastly, we note that the SMDWM results go beyond the current models predicting miscible blend dynamics. Current models of miscible blend dynamics consider fast and slow relaxing components within the Kuhn volumes corresponding to the low- and high- $T_{\rm g}$ components, respectively. Here, the switching in the dynamics between the continuously and intermittently rotating molecules at elevated temperatures and its relation to the dynamic heterogeneity requires further consideration by theory and experiment.

CONCLUSIONS

In conclusion, the spatial and temporal heterogeneities in the fully miscible blend films of PS and OS have been investigated by monitoring the rotational diffusion of embedded probe molecules with SMDWM. Three types of molecular behaviors have been observed: immobilized, continuously rotating, and intermittently rotating, and these behaviors distribute randomly across the blend films, revealing spatial heterogeneity that can be linked to differences in local compositions. The immobilized mode is most likely associated with the PS-rich regions, and the occurrence probability of it depends on the blend composition rather than temperature. In contrast, the occurrence probabilities of the continuously and the intermittently rotating molecules are dependent on both the blend composition and temperature. Remarkably, these two modes can switch with temperature, indicating a possible alteration in local dynamics that is triggered by the dynamic heterogeneity in the blends. Additionally, the rotational autocorrelation function reveals nonexponential characteristic of temporal heterogeneity. Some of these features can be qualitatively discussed with the help of the self-concentration model. Our work demonstrates that SMDWM is a powerful tool to study the complex dynamics in miscible polymer blends and allows gaining a deeper understanding of spatial and temporal heterogeneity in dynamically asymmetric yet fully miscible blends.

■ EXPERIMENTAL SECTION

Sample Preparation. The PS $(M_w = 6140 \text{ g mol}^{-1}, D =$ 1.04) and oligomer OS $(M_w = 370 \text{ g mol}^{-1}, D = 1.00)$ were purchased from Polymer Standards Service GmbH. The glass temperatures $(T_{o}s)$ of PS and OS were estimated to be ~358 and ~236 K, respectively, by differential scanning calorimetry (Pyris1, Perkin-Elmer). The PS and OS blends were prepared in toluene with a PS to OS weight ratio of 75:25, 50:50, or 25:75. The PDI dyes in toluene were added into the blend solution, and then the blend solution was further diluted to 1 wt %. The blend films were made by spin-coating 1 wt % blends with PDI dyes (~10⁻⁹ M) in toluene onto cleaned cover glasses. The films were annealed for more than 3 h at ~373 K under vacuum to remove the residual solvent. The elimination of the residual solvent is very critical. If not, the PDI molecules undergo Brownian motion (translational diffusion), making it impossible to obtain the correct dynamics information of polymer blends. The samples were measured within one day after preparation. The thickness of the PS/OS blend films was estimated by an atomic force microscope.

Experimental Setup and Procedures. An inverted microscope (IX71, Olympus) equipped with an oil immersion objective (Plan Fluorite, Olympus, 1.3-N.A. $100\times$) and an EMCCD camera (ImagEM, Hamamatsu) was used to perform the defocused wide-field fluorescence imaging. The excitation light from a 532 nm laser (gem532-smd6000, Laser Quantum) was circularly polarized by using $\lambda/2$ and $\lambda/4$ waveplates. Fluorescence of single molecules was collected by the same immersion objective and then passed through a dichroic mirror (z532rdc, Chroma Technology Co.) and a long pass filter

(HQ542LP, Chroma Technology Co.). An additional lens was used to further magnify the image by 3.3 times, and therefore an imaging view of $24.6 \times 24.6 \ \mu\text{m}^2$ is obtained. A quasi total internal reflection mode (qTIRF) has been employed to enhance the excitation of molecules oriented along the z-axis. The excitation field contains only a limited contribution of radial or z-polarized light when using a standard Köhler illumination. A mantle around the objective and microscope stage was applied to control the temperature during measurements along with a thermostat water bath. To reduce the effect of refractive index change with temperature, different types of immersion oils were used below and above 310 K (DF and type-37, Electron Microscopy Sciences).

The measurements were performed starting from 315 K and cooling down with 5 K temperature steps. At each temperature, the sample was incubated for at least 2 h before measurements. In order to reduce photo-bleaching of single molecules, the samples were measured under nitrogen. To perform defocused imaging, the wide-field system has been modified by working slightly out of focus with $\sim\!\!1~\mu\rm m$ defocused distance. Sequences of defocused images have been recorded up to 5000–10,000 frames with integration times between 30 and 100 ms. A home-built MatLab routine was used to analyze the resulting defocused images according to the theoretical model. 36,55

Details of Defocused Imaging Analysis. The defocused images were analyzed by a pattern matching routine written in MatLab. 31,32,36,55 The defocused image patterns are roughly matched with the calculated patterns every 10 degrees in- and out-of-plane, and then the orientation was further matched by varying one degree in the in-plane and out-of-plane orientations. The defocused patterns of the single molecules in the film depend on the *z*-position. For single molecules with different defocused distances, the corresponding pattern matrix is generated for the fitting (see Supporting Information for details). The in-plane (φ) and out-of-plane (θ) angles are used to express the 3D orientation of the transition dipole moment of a single molecule. Note that the *z*-axis is parallel to the optical axis and the x-y plane is the sample surface. The 3D angle displacement is defined by,

$$\Phi = \arccos\left(\mathbf{n}(t) \cdot \mathbf{n}(t + \delta t)\right) \tag{1}$$

where $\mathbf{n}(t)$ is the unit orientation vector of the transition dipole moment at time t and δt is the time interval between consecutive images.

The rotational correlation time was obtained by fitting autocorrelation functions of $\langle\cos\Phi\rangle$ with the Kohlrausch–Williams–Watts (KWW)-stretched exponential function,

$$C(t) = \exp[-(t/\tau_{\text{KWW}})^{\beta_{\text{KWW}}}]$$
 (2)

where $\tau_{\rm KWW}$ is the decay constant and $\beta_{\rm KWW}$ is the stretched exponent, which describes the width of the distribution of relaxation times. $\beta_{\rm KWW}=1$ corresponds to a single exponential decay. The correlation time $\tau_{\rm C}$ can be calculated by the following equation,

$$\tau_{c} = \int_{0}^{\infty} C(t) dt = \frac{\tau_{\text{KWW}}}{\beta} \Gamma\left(\frac{1}{\beta}\right)$$
(3)

where Γ denotes a gamma function.

ASSOCIATED CONTENT

Solution Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acsomega.0c03173.

Dielectric spectroscopy measurements and data, defocused imaging analysis for single molecules with different defocused distances, defocused images of single molecules in pure PS and OS films, histograms of rotational correlation times, fitting parameters of the KWW stretched exponential function and calculated correlation times, and Arrhenius diagram for comparing between single molecule data and dielectric spectroscopy data (PDF)

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Notes

The authors declare no competing financial interest.

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