# Quantum differential equations and helices 

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#### Abstract

These notes are a short and self-contained introduction to the isomonodromic approach to quantum cohomology, and Dubrovin's conjecture. An overview of recent results obtained in joint works with B. Dubrovin and D. Guzzetti [6], and A. Varchenko [9] is given.


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## 1. Quantum cohomology

### 1.1. Notations and conventions

Let $X$ be a smooth projective variety over $\mathbb{C}$ with vanishing odd-cohomology, i.e. $H^{2 k+1}(X, \mathbb{C})=0$, for $k \geq 0$. Fix a homogeneous basis $\left(T_{1}, \ldots, T_{n}\right)$ of the complex vector space $H^{\bullet}(X):=\bigoplus_{k} H^{2 k}(X, \mathbb{C})$, and denote by $\boldsymbol{t}:=$ $\left(t^{1}, \ldots, t^{n}\right)$ the corresponding dual coordinates. Without loss of generality, we assume that $T_{1}=1$. The Poincaré pairing on $H^{\bullet}(X)$ will be denoted by

$$
\begin{equation*}
\eta(u, v):=\int_{X} u \cup v, \quad u, v \in H^{\bullet}(X) \tag{1.1}
\end{equation*}
$$

and we put $\eta_{\alpha \beta}:=\eta\left(T_{\alpha}, T_{\beta}\right)$, for $\alpha, \beta=1, \ldots, n$, to be the Gram matrix wrt the fixed basis. The entries of the inverse matrix will be denoted by $\eta^{\alpha \beta}$, for $\alpha, \beta=1, \ldots, n$. In all the paper, the Einstein rule of summation over repeated indices is used. General references for this Section are [5, 6, 10, 11, 12, 13, 27, 29, 31.

### 1.2. Gromov-Witten invariants in genus 0

For a fixed $\beta \in H_{2}(X, \mathbb{Z}) /$ torsion, denote by $\overline{\mathcal{M}}_{0, k}(X, \beta)$ the Deligne-Mumford moduli stack of $k$-pointed stable rational maps with target $X$ of degree $\beta$ :

$$
\begin{equation*}
\overline{\mathcal{M}}_{0, k}(X, \beta):=\left\{f:(C, \boldsymbol{x}) \rightarrow X, f_{*}[C]=\beta\right\} / \text { equivalencies, } \tag{1.2}
\end{equation*}
$$

where $C$ is an algebraic curve of genus 0 with at most nodal singularities, $\boldsymbol{x}:=\left(x_{1}, \ldots, x_{k}\right)$ is a $k$-tuple of pairwise distinct marked points of $C$, and equivalencies are automorphisms of $C \rightarrow X$ identical on $X$ and the markings.

Gromov-Witten invariants ( $G W$-invariants for short) of $X$, and their descendants, are defined as intersection numbers of cycles on $\overline{\mathcal{M}}_{0, k}(X, \beta)$, by the integrals

$$
\begin{equation*}
\left\langle\tau_{d_{1}} \gamma_{1}, \ldots, \tau_{d_{k}} \gamma_{k}\right\rangle_{k, \beta}^{X}:=\int_{\left[\overline{\mathcal{M}}_{0, k}(X, \beta)\right]_{\mathrm{virt}}} \prod_{i=1}^{k} \operatorname{ev}_{i}^{*} \gamma_{i} \wedge \psi_{i}^{d_{i}} \tag{1.3}
\end{equation*}
$$

for $\gamma_{1}, \ldots, \gamma_{k} \in H^{\bullet}(X), d_{i} \in \mathbb{N}$. In formula (1.3),

$$
\begin{equation*}
\mathrm{ev}_{i}: \overline{\mathcal{M}}_{0, k}(X, \beta) \rightarrow X, \quad f \mapsto f\left(x_{i}\right), \quad i=1, \ldots, k \tag{1.4}
\end{equation*}
$$

are evaluation maps, and $\psi_{i}:=c_{1}\left(\mathcal{L}_{i}\right)$ are the first Chern classes of the universal cotangent line bundles

$$
\begin{equation*}
\mathcal{L}_{i} \rightarrow \overline{\mathcal{M}}_{0, k}(X, \beta),\left.\quad \mathcal{L}_{i}\right|_{f}=T_{x_{i}}^{*} C, \quad i=1, \ldots, k \tag{1.5}
\end{equation*}
$$

The virtual fundamental cycle $\left[\overline{\mathcal{M}}_{0, k}(X, \beta)\right]^{\text {virt }}$ is an element of the Chow $\operatorname{ring} A \bullet\left(\overline{\mathcal{M}}_{0, k}(X, \beta)\right)$, namely

$$
\left[\overline{\mathcal{M}}_{0, k}(X, \beta)\right]^{\mathrm{virt}} \in A_{D}\left(\overline{\mathcal{M}}_{0, k}(X, \beta)\right), \quad D:=\operatorname{dim}_{\mathbb{C}} X-3+k+\int_{\beta} c_{1}(X)
$$

See [1] for its construction.

### 1.3. Quantum cohomology as a Frobenius manifold

Introduce infinitely many variables $\boldsymbol{t}_{\boldsymbol{\bullet}}:=\left(t_{p}^{\alpha}\right)_{\alpha, p}$ with $\alpha=1, \ldots, n$ and $p \in \mathbb{N}$.
Definition 1.1. The genus 0 total descendant potential of $X$ is the generating function $\mathcal{F}_{0}^{X} \in \mathbb{C} \llbracket \boldsymbol{t}_{\bullet} \rrbracket$ of descendant $G W$-invariants of $X$ defined by

$$
\mathcal{F}_{0}^{X}\left(\boldsymbol{t}_{\bullet}\right):=\sum_{k=0}^{\infty} \sum_{\beta} \sum_{\alpha_{1}, \ldots, \alpha_{k}=1}^{n} \sum_{p_{1}, \ldots, p_{k}=0}^{\infty} \frac{t_{p_{1}}^{\alpha_{1}} \ldots t_{p_{k}}^{\alpha_{k}}}{k!}\left\langle\tau_{p_{1}} T_{\alpha_{1}}, \ldots, \tau_{p_{k}} T_{\alpha_{k}}\right\rangle_{k, \beta}^{X}
$$

Setting $t_{0}^{\alpha}=t^{\alpha}$ and $t_{p}^{\alpha}=0$ for $p>0$, we obtain the Gromov-Witten potential of $X$

$$
\begin{equation*}
F_{0}^{X}(\boldsymbol{t}):=\sum_{k=0}^{\infty} \sum_{\beta} \sum_{\alpha_{1}, \ldots, \alpha_{k}=1}^{n} \frac{t^{\alpha_{1}} \ldots t^{\alpha_{k}}}{k!}\left\langle T_{\alpha_{1}}, \ldots, T_{\alpha_{k}}\right\rangle_{k, \beta}^{X} . \tag{1.6}
\end{equation*}
$$

Let $\Omega \subseteq H^{\bullet}(X)$ be the domain of convergence of $F_{0}^{X}(\boldsymbol{t})$, assumed to be non-empty. We denote by $T \Omega$ and $T^{*} \Omega$ its holomorphic tangent and cotangent bundles, respectively. Each tangent space $T_{p} \Omega$, with $p \in \Omega$, is canonically identified with the space $H^{\bullet}(X)$, via the identification $\frac{\partial}{\partial t^{\alpha}} \mapsto T_{\alpha}$. The Poincaré metric $\eta$ defines a flat non-degenerate $\mathcal{O}_{\Omega}$-bilinear pseudoriemannian metric on $\Omega$. The coordinates $t$ are manifestly flat. Denote by $\nabla$ the Levi-Civita connection of $\eta$.
Definition 1.2. Define the tensor $c \in \Gamma\left(T \Omega \otimes \bigodot^{2} T^{*} \Omega\right)$ by

$$
\begin{equation*}
c_{\beta \gamma}^{\alpha}:=\eta^{\alpha \lambda} \nabla_{\lambda \beta \gamma}^{3} F_{0}^{X}, \quad \alpha, \beta, \gamma=1, \ldots, n, \tag{1.7}
\end{equation*}
$$

and let us introduce a product $*$ on vector fields on $\Omega$ by

$$
\begin{equation*}
\frac{\partial}{\partial t^{\beta}} * \frac{\partial}{\partial t^{\gamma}}:=c_{\beta \gamma}^{\alpha} \frac{\partial}{\partial t^{\alpha}}, \quad \beta, \gamma=1, \ldots, n \tag{1.8}
\end{equation*}
$$

Theorem 1.3 ([27, 31]). The Gromov-Witten potential $F_{0}^{X}(\boldsymbol{t})$ is a solution of WDVV equations

$$
\begin{equation*}
\frac{\partial^{3} F_{0}^{X}(\boldsymbol{t})}{\partial t^{\alpha} \partial t^{\beta} \partial t^{\gamma}} \eta^{\gamma \delta} \frac{\partial^{3} F_{0}^{X}(\boldsymbol{t})}{\partial t^{\delta} \partial t^{\epsilon} \partial t^{\phi}}=\frac{\partial^{3} F_{0}^{X}(\boldsymbol{t})}{\partial t^{\phi} \partial t^{\beta} \partial t^{\gamma}} \eta^{\gamma \delta} \frac{\partial^{3} F_{0}^{X}(\boldsymbol{t})}{\partial t^{\delta} \partial t^{\epsilon} \partial t^{\alpha}}, \tag{1.9}
\end{equation*}
$$

for $\alpha, \beta, \epsilon, \phi=1, \ldots, n$.
On each tangent space $T_{p} \Omega$, the product $*_{p}$ defines a structure of associative, commutative algebra with unit $\frac{\partial}{\partial t^{1}} \equiv 1$. Furthermore, the product $*$ is compatible with the Poincare metric, namely

$$
\begin{equation*}
\eta(u * v, w)=\eta(u, v * w), \quad u, v, w \in \Gamma(T \Omega) \tag{1.10}
\end{equation*}
$$

This endows $\left(T_{p} \Omega, *_{p}, \eta_{p},\left.\frac{\partial}{\partial t^{1}}\right|_{p}\right)$ with a complex Frobenius algebra structure.
Definition 1.4. The vector field

$$
\begin{equation*}
E=c_{1}(X)+\sum_{\alpha=1}^{n}\left(1-\frac{1}{2} \operatorname{deg} T_{\alpha}\right) t^{\alpha} \frac{\partial}{\partial t^{\alpha}} \tag{1.11}
\end{equation*}
$$

is called Euler vector field. Here, $\operatorname{deg} T_{\alpha}$ denotes the cohomological degree of $T_{\alpha}$, i.e. $\operatorname{deg} T_{\alpha}:=r_{\alpha}$ if and only if $T_{\alpha} \in H^{r_{\alpha}}(X, \mathbb{C})$. We denote by $\mathcal{U}$ the $(1,1)$-tensor defined by the multiplication with the Euler vector field, i.e.

$$
\begin{equation*}
\mathcal{U}: \Gamma(T \Omega) \rightarrow \Gamma(T \Omega), \quad v \mapsto E * v \tag{1.12}
\end{equation*}
$$

Proposition 1.5 ([11, [13)). The Euler vector field $E$ is a Killing conformal vector field, whose flow preserves the structure constants of the Frobenius algerbas:

$$
\begin{equation*}
\mathfrak{L}_{E} \eta=\left(2-\operatorname{dim}_{\mathbb{C}} X\right) \eta, \quad \mathfrak{L}_{E} c=c \tag{1.13}
\end{equation*}
$$

The structure $\left(\Omega, c, \eta, \frac{\partial}{\partial t^{1}}, E\right)$ gives an example of analytic Frobenius manifold, called quantum cohomology of $X$ and denoted by $Q H^{\bullet}(X)$, see [11, 12, 13, 29].

### 1.4. Extended deformed connection

Definition 1.6. The grading operator $\mu \in \operatorname{End}(T \Omega)$ is the tensor defined by

$$
\begin{equation*}
\mu(v):=\frac{2-\operatorname{dim}_{\mathbb{C}} X}{2} v-\nabla_{v} E, \quad v \in \Gamma(T \Omega) \tag{1.14}
\end{equation*}
$$

Consider the canonical projection $\pi: \mathbb{C}^{*} \times \Omega \rightarrow \Omega$, and the pull-back bundle $\pi^{*} T \Omega$. Denote by

1. $\mathscr{T}_{\Omega}$ the sheaf of sections of $T \Omega$,
2. $\pi^{*} \mathscr{T}_{\Omega}$ the pull-back sheaf, i.e. the sheaf of sections of $\pi^{*} T \Omega$
3. $\pi^{-1} \mathscr{T}_{\Omega}$ the sheaf of sections of $\pi^{*} T \Omega$ constant on the fibers of $\pi$.

All the tensors $\eta, c, E, \mathcal{U}, \mu$ can be lifted to $\pi^{*} T \Omega$, and their lifts will be denoted by the same symbols. The Levi-Civita connection $\nabla$ is lifted on $\pi^{*} T \Omega$, and it acts so that

$$
\begin{equation*}
\nabla_{\frac{\partial}{\partial z}} v=0 \quad \text { for } v \in\left(\pi^{-1} \mathscr{T}_{\Omega}\right)(\Omega) \tag{1.15}
\end{equation*}
$$

where $z$ is the coordinate on $\mathbb{C}^{*}$.

Definition 1.7. The extended deformed connection is the connection $\widehat{\nabla}$ on the bundle $\pi^{*} T \Omega$ defined by

$$
\begin{align*}
\widehat{\nabla}_{w} v & =\nabla_{w} v+z \cdot w * v  \tag{1.16}\\
\widehat{\nabla}_{\frac{\partial}{\partial z}} v & =\nabla_{\partial_{z}} v+\mathcal{U}(v)-\frac{1}{z} \mu(v) \tag{1.17}
\end{align*}
$$

for $v, w \in \Gamma\left(\pi^{*} T \Omega\right)$.
Theorem 1.8 ([11, [13]). The connection $\widehat{\nabla}$ is flat.

### 1.5. Semisimple points and orthonormalized idempotent frame

Definition 1.9. A point $p \in \Omega$ is semisimple if and only if the corresponding Frobenius algebra ( $T_{p} \Omega, *_{p}, \eta_{p},\left.\frac{\partial}{\partial t^{1}}\right|_{p}$ ) is without nilpotents. Denote by $\Omega_{s s}$ the open dense subset of $\Omega$ of semisimple points.
Theorem $1.10([24])$. The set $\Omega_{s s}$ is non-empty only if $X$ is of Hodge-Tat $\rrbracket^{1}$ type, i.e. $h^{p, q}(X)=0$ for $p \neq q$.

On $\Omega_{s s}$ there are $n$ well-defined idempotent vector fields $\pi_{1}, \ldots, \pi_{n} \in$ $\Gamma\left(T \Omega_{s s}\right)$, satisfying

$$
\begin{equation*}
\pi_{i} * \pi_{j}=\delta_{i j} \pi_{i}, \quad \eta\left(\pi_{i}, \pi_{j}\right)=\delta_{i j} \eta\left(\pi_{i}, \pi_{i}\right), \quad i, j=1, \ldots, n . \tag{1.18}
\end{equation*}
$$

Theorem 1.11 (10, 11, 13). The idempotent vector fields pairwise commute: $\left[\pi_{i}, \pi_{j}\right]=0$ for $i, j=1, \ldots, n$. Hence, there exist holomorphic local coordinates $\left(u_{1}, \ldots, u_{n}\right)$ on $\Omega_{s s}$ such that $\frac{\partial}{\partial u_{i}}=\pi_{i}$ for $i=1, \ldots, n$.
Definition 1.12. The coordinates $\left(u_{1}, \ldots, u_{n}\right)$ of Theorem 1.11 are called canonical coordinates.

Proposition 1.13 (11, 13). Canonical coordinates are uniquely defined up to ordering and shifts by constants. The eigenvalues of the tensor $\mathcal{U}$ define a system of canonical coordinates in a neighborhood of any semisimple point of $\Omega_{s s}$.

Definition 1.14. We call orthonormalized idempotent frame a frame $\left(f_{i}\right)_{i=1}^{n}$ of $T \Omega_{s s}$ defined by

$$
\begin{equation*}
f_{i}:=\eta\left(\pi_{i}, \pi_{i}\right)^{-\frac{1}{2}} \pi_{i}, \quad i=1, \ldots, n, \tag{1.19}
\end{equation*}
$$

for arbitrary choices of signs of the square roots. The $\Psi$-matrix is the matrix $\left(\Psi_{i \alpha}\right)_{i, \alpha=1}^{n}$ of change of tangent frames, defined by

$$
\begin{equation*}
\frac{\partial}{\partial t^{\alpha}}=\sum_{i=1}^{n} \Psi_{i \alpha} f_{i}, \quad \alpha=1, \ldots, n \tag{1.20}
\end{equation*}
$$

Remark 1.15. In the orthonormalized idempotent frame, the operator $\mathcal{U}$ is represented by a diagonal matrix, and the operator $\mu$ by an antisymmetric matrix:

$$
\begin{gather*}
U:=\operatorname{diag}\left(u_{1}, \ldots, u_{n}\right), \quad \Psi \mathcal{U} \Psi^{-1}=U,  \tag{1.21}\\
V:=\Psi \mu \Psi^{-1}, \quad V^{T}+V=0 \tag{1.22}
\end{gather*}
$$

[^0]
## 2. Quantum differential equation

The connection $\widehat{\nabla}$ induces a flat connection on $\pi^{*}\left(T^{*} \Omega\right)$. Let $\xi \in \Gamma\left(\pi^{*}\left(T^{*} \Omega\right)\right)$ be a flat section. Consider the corresponding vector field $\zeta \in \Gamma\left(\pi^{*}(T \Omega)\right)$ via musical isomorphism, i.e. such that $\xi(v)=\eta(\zeta, v)$ for all $v \in \Gamma\left(\pi^{*}(T \Omega)\right)$.

The vector field $\zeta$ satisfies the following system ${ }^{2}$ of equations

$$
\begin{align*}
\frac{\partial}{\partial t^{\alpha}} \zeta & =z \mathcal{C}_{\alpha} \zeta, \quad \alpha=1, \ldots, n  \tag{2.1}\\
\frac{\partial}{\partial z} \zeta & =\left(\mathcal{U}+\frac{1}{z} \mu\right) \zeta \tag{2.2}
\end{align*}
$$

Here $\mathcal{C}_{\alpha}$ is the $(1,1)$-tensor defined by $\left(\mathcal{C}_{\alpha}\right)_{\gamma}^{\beta}:=c_{\alpha \gamma}^{\beta}$.
Definition 2.1. The quantum differential equation $(q D E)$ of $X$ is the differential equation (2.2).

The $q D E$ is an ordinary differential equation with rational coefficients. It has two singularities on the Riemann sphere $\mathbb{P}^{1}(\mathbb{C})$ :

1. a Fuchsian singularity at $z=0$,
2. an irregular singularity (of Poincaré rank 1) at $z=\infty$.

Points of $\Omega$ are parameters of deformation of the coefficients of the $q D E$. Solutions $\zeta(\boldsymbol{t}, z)$ of the joint system of equations (2.1), (2.2) are "multivalued" functions wrt $z$, i.e. they are well-defined functions on $\Omega \times \widehat{\mathbb{C}^{*}}$, where $\widehat{\mathbb{C}^{*}}$ is the universal cover of $\mathbb{C}^{*}$.
2.1. Solutions in Levelt form at $z=0$ and topological-enumerative solution

Theorem 2.2 ([5, 11, [13). There exist fundamental systems of solutions $Z_{0}(\boldsymbol{t}, z)$ of the joint system (2.1), (2.2) with expansions at $z=0$ of the form

$$
\begin{equation*}
Z_{0}(\boldsymbol{t}, z)=F(\boldsymbol{t}, z) z^{\mu} z^{R}, \quad R=\sum_{k \geq 1} R_{k}, \quad F(\boldsymbol{t}, z)=I+\sum_{j=1}^{\infty} F_{j}(\boldsymbol{t}) z^{j} \tag{2.3}
\end{equation*}
$$

where $\left(R_{k}\right)_{\alpha \beta} \neq 0$ only if $\mu_{\alpha}-\mu_{\beta}=k$. The series $F(\boldsymbol{t}, z)$ is convergent and satisfies the orthogonality condition

$$
\begin{equation*}
F(\boldsymbol{t},-z)^{T} \eta F(\boldsymbol{t}, z)=\eta . \tag{2.4}
\end{equation*}
$$

Definition 2.3. A fundamental system of solutions $Z_{0}(\boldsymbol{t}, z)$ of the form described in Theorem 2.3 are said to be in Levelt form at $z=0$.

Remark 2.4. Fundamental systems of solutions in Levelt form are not unique. The exponent $R$ is not uniquely determined. Moreover, even for a fixed exponent $R$, the series $F(\boldsymbol{t}, z)$ is not uniquely determined, see [5]. It can be proved that the matrix $R$ can be chosen as the matrix of the operator $c_{1}(X) \cup(-): H^{\bullet}(X) \rightarrow H^{\bullet}(X)$ wrt the basis $\left(T_{\alpha}\right)_{\alpha=1}^{n}$ [13, Corollary 2.1].

[^1]Remark 2.5. Let $Z_{0}(\boldsymbol{t}, z)$ be a fundamental system of solutions in Levelt form (2.3). The monodromy matrix $M_{0}(\boldsymbol{t})$, defined by

$$
\begin{equation*}
Z_{0}\left(\boldsymbol{t}, e^{2 \pi \sqrt{-1}} z\right)=Z_{0}(\boldsymbol{t}, z) M_{0}(\boldsymbol{t}), \quad z \in \widehat{\mathbb{C}^{*}} \tag{2.5}
\end{equation*}
$$

is given by

$$
\begin{equation*}
M_{0}(\boldsymbol{t})=\exp (2 \pi \sqrt{-1} \mu) \exp (2 \pi \sqrt{-1} R) \tag{2.6}
\end{equation*}
$$

In particular, $M_{0}$ does not depend on $\boldsymbol{t}$.
Definition 2.6. Define the functions $\theta_{\beta, p}(\boldsymbol{t}, z), \theta_{\beta}(\boldsymbol{t}, z)$, with $\beta=1, \ldots, n$ and $p \in \mathbb{N}$, by

$$
\begin{gather*}
\theta_{\beta, p}(\boldsymbol{t}):=\left.\frac{\partial^{2} \mathcal{F}_{0}^{X}\left(\boldsymbol{t}_{\boldsymbol{\bullet}}\right)}{\partial t_{0}^{1} \partial t_{p}^{\beta}}\right|_{t_{p}^{\alpha}=0 \text { for } p>1, \quad t_{0}^{\alpha}=t^{\alpha} \text { for } \alpha=1, \ldots, n}  \tag{2.7}\\
\theta_{\beta}(\boldsymbol{t}, z):=\sum_{p=0}^{\infty} \theta_{\beta, p}(\boldsymbol{t}) z^{p} \tag{2.8}
\end{gather*}
$$

Define the matrix $\Theta(\boldsymbol{t}, z)$ by

$$
\begin{equation*}
\Theta(\boldsymbol{t}, z)_{\beta}^{\alpha}:=\eta^{\alpha \lambda} \frac{\partial \theta_{\beta}(\boldsymbol{t}, z)}{\partial t^{\lambda}}, \quad \alpha, \beta=1, \ldots, n \tag{2.9}
\end{equation*}
$$

Theorem 2.7 ([5, [13]). The matrix $Z_{\mathrm{top}}(\boldsymbol{t}, z):=\Theta(\boldsymbol{t}, z) z^{\mu} z^{c_{1}(X) \cup}$ is a fundamental system of solutions of the joint system (2.1)-(2.2) in Levelt form at $z=0$.

Definition 2.8. The solution $Z_{\mathrm{top}}(\boldsymbol{t}, z)$ is called topological-enumerative solution of the joint system (2.1), (2.2).

### 2.2. Stokes rays and $\ell$-chamber decomposition

Definition 2.9. We call Stokes rays at a point $p \in \Omega$ the oriented rays $R_{i j}(p)$ in $\mathbb{C}$ defined by

$$
\begin{equation*}
R_{i j}(p):=\left\{-\sqrt{-1}\left(\overline{u_{i}(p)}-\overline{u_{j}(p)}\right) \rho: \rho \in \mathbb{R}_{+}\right\} \tag{2.10}
\end{equation*}
$$

where $\left(u_{1}(p), \ldots, u_{n}(p)\right)$ is the spectrum of the operator $\mathcal{U}(p)$ (with a fixed arbitrary order).

Fix an oriented ray $\ell$ in the universal cover $\widehat{\mathbb{C}^{*}}$.
Definition 2.10. We say that $\ell$ is admissible at $p \in \Omega$ if the projection of the the ray $\ell$ on $\mathbb{C}^{*}$ does not coincide with any Stokes ray $R_{i j}(p)$.

Definition 2.11. Define the open subset $O_{\ell}$ of points $p \in \Omega$ by the following conditions:

1. the eigenvalues $u_{i}(p)$ are pairwise distinct,
2. $\ell$ is admissible at $p$.

We call $\ell$-chamber of $\Omega$ any connected component of $O_{\ell}$.

### 2.3. Stokes fundamental solutions at $z=\infty$

Fix an oriented ray $\ell \equiv\{\arg z=\phi\}$ in $\widehat{\mathbb{C}^{*}}$. For $m \in \mathbb{Z}$, define the sectors in $\widehat{\mathbb{C}^{*}}$

$$
\begin{align*}
& \Pi_{L, m}(\phi):=\left\{z \in \widehat{\mathbb{C}^{*}}: \phi+2 \pi m<\arg z<\phi+\pi+2 \pi m\right\},  \tag{2.11}\\
& \Pi_{R, m}(\phi):=\left\{z \in \widehat{\mathbb{C}^{*}}: \phi-\pi+2 \pi m<\arg z<\phi+2 \pi m\right\} . \tag{2.12}
\end{align*}
$$

Definition 2.12. The coalescence locus of $\Omega$ is the set

$$
\begin{equation*}
\Delta_{\Omega}:=\left\{p \in \Omega: u_{i}(p)=u_{j}(p), \quad \text { for some } i \neq j\right\} \tag{2.13}
\end{equation*}
$$

Theorem 2.13 (11, 13]). There exists a unique formal solution $Z_{\text {form }}(\boldsymbol{t}, z)$ of the joint system (2.1), (2.2) of the form

$$
\begin{align*}
Z_{\text {form }}(\boldsymbol{t}, z) & =\Psi(\boldsymbol{t})^{-1} G(\boldsymbol{t}, z) \exp (z U(\boldsymbol{t})),  \tag{2.14}\\
G(\boldsymbol{t}, z) & =I+\sum_{k=1}^{\infty} \frac{1}{z^{k}} G_{k}(\boldsymbol{t}) \tag{2.15}
\end{align*}
$$

where the matrices $G_{k}(\boldsymbol{t})$ are holomorphic on $\Omega \backslash \Delta_{\Omega}$.
Theorem 2.14 (11, 13). Let $m \in \mathbb{Z}$. There exist unique fundamental systems of solutions $Z_{L, m}(\boldsymbol{t}, z), Z_{R, m}(\boldsymbol{t}, z)$ of the joint system (2.1), (2.2) with asymptotic expansion

$$
\begin{array}{ll}
Z_{L, m}(\boldsymbol{t}, z) \sim Z_{\text {form }}(\boldsymbol{t}, z), & |z| \rightarrow \infty, \\
Z_{R, m}(\boldsymbol{t}, z) \sim Z_{\text {form }}(\boldsymbol{t}, z), & z \mid \rightarrow \Pi_{L, m}(\phi)  \tag{2.17}\\
\end{array}
$$

respectively.
Definition 2.15. The solutions $Z_{L, m}(\boldsymbol{t}, z)$ and $Z_{R, m}(\boldsymbol{t}, z)$ are called Stokes fundamental solutions of the joint system (2.1), (2.2) on the sectors $\Pi_{L, m}(\phi)$ and $\Pi_{R, m}(\phi)$ respectively.

### 2.4. Monodromy data

Let $\ell \equiv\{\arg z=\phi\}$ be an oriented ray in $\widehat{\mathbb{C}^{*}}$ and consider the corresponding Stokes fundamental systems of solutions $Z_{L, m}(\boldsymbol{t}, z), Z_{R, m}(\boldsymbol{t}, z)$, for $m \in \mathbb{Z}$.

Definition 2.16. We define the Stokes and central connection matrices $S^{(m)}(p)$, $C^{(m)}(p)$, with $m \in \mathbb{Z}$, at the point $p \in O_{\ell}$ by the identities

$$
\begin{array}{r}
Z_{L, m}(\boldsymbol{t}(p), z)=Z_{R, m}(\boldsymbol{t}(p), z) S^{(m)}(p), \\
Z_{R, m}(\boldsymbol{t}(p), z)=Z_{\mathrm{top}}(\boldsymbol{t}(p), z) C^{(m)}(p) . \tag{2.19}
\end{array}
$$

Set $S(p):=S^{(0)}(p)$ and $C(p):=C^{(0)}(p)$.
Definition 2.17. The monodromy data at the point $p \in O_{\ell}$ are defined as the 4-tuple ( $\mu, R, S(p), C(p)$ ), where

- $\mu$ is the (matrix associated to) the grading operator,
- $R$ is the (matrix associated to) the operator $c_{1}(X) \cup: H^{\bullet}(X) \rightarrow H^{\bullet}(X)$,
- $S(p), C(p)$ are the Stokes and central connection matrices at $p$, respectively.

Remark 2.18. The definition of the Stokes and central connection matrices is subordinate to several non-canonical choices:

1. the choice of an oriented ray $\ell$ in $\widehat{\mathbb{C}^{*}}$,
2. the choice of an ordering of canonical coordinates $u_{1}, \ldots, u_{n}$ on each $\ell$-chamber,
3. the choice of signs in (1.19), and hence of the branch of the $\Psi$-matrix on each $\ell$-chamber.
Different choices affect the numerical values of the data $(S, C)$, see [5]. In particular, for different choices of ordering of canonical coordinates, the Stokes and central connection matrices transform as follows:

$$
\begin{equation*}
S \mapsto \Pi S \Pi^{-1}, \quad C \mapsto C \Pi^{-1}, \quad \Pi \text { permutation matrix. } \tag{2.20}
\end{equation*}
$$

Definition 2.19. Fix a point $p \in O_{\ell}$ with canonical coordinates $\left(u_{i}(p)\right)_{i=1}^{n}$. Define the oriented rays $L_{j}(p, \phi), j=1, \ldots, n$, in the complex plane by the equations

$$
\begin{equation*}
L_{j}(p, \phi):=\left\{u_{j}(p)+\rho e^{\sqrt{-1}\left(\frac{\pi}{2}-\phi\right)}: \rho \in \mathbb{R}_{+}\right\} \tag{2.21}
\end{equation*}
$$

The ray $L_{j}(p, \phi)$ is oriented from $u_{j}(p)$ to $\infty$. We say that $\left(u_{i}(p)\right)_{i=1}^{n}$ are in $\ell$-lexicographical order if $L_{j}(p, \phi)$ is on the left of $L_{k}(p, \phi)$ for $1 \leq j<k \leq n$.

In what follows, it is assumed that the $\ell$-lexicographical order of canonical coordinates is fixed at all points of $\ell$-chambers.

Lemma 2.20 (5, 13]). If the canonical coordinates $\left(u_{i}(p)\right)_{i=1}^{n}$ are in $\ell$-lexicographical order at $p \in O_{\ell}$, then the Stokes matrices $S^{(m)}(p), m \in \mathbb{Z}$, are upper triangular with 1 's along the diagonal.

By Remarks 2.4 and 2.5 the matrices $\mu$ and $R$ determine the monodromy of solutions of the $q D E$,

$$
\begin{equation*}
M_{0}:=\exp (2 \pi \sqrt{-1} \mu) \exp (2 \pi \sqrt{-1} R) \tag{2.22}
\end{equation*}
$$

Moreover, $\mu$ and $R$ do not depend on the point $p$. The following theorem furnishes a refinement of this property.

Theorem 2.21 ([5, 11, 13]). The monodromy data ( $\mu, R, S, C$ ) are constant in each $\ell$-chamber. Moreover, they satisfy the following identities:

$$
\begin{gather*}
C S^{T} S^{-1} C^{-1}=M_{0}  \tag{2.23}\\
S=C^{-1} \exp (-\pi \sqrt{-1} R) \exp (-\pi \sqrt{-1} \mu) \eta^{-1}\left(C^{T}\right)^{-1}  \tag{2.24}\\
S^{T}=C^{-1} \exp (\pi \sqrt{-1} R) \exp (\pi \sqrt{-1} \mu) \eta^{-1}\left(C^{T}\right)^{-1} \tag{2.25}
\end{gather*}
$$

Theorem 2.22 ([5). The Stokes and central connection matrices $S_{m}, C_{m}$, with $m \in \mathbb{Z}$, can be reconstructed from the monodromy data $(\mu, R, S, C)$ :

$$
\begin{equation*}
S^{(m)}=S, \quad C^{(m)}=M_{0}^{-m} C, \quad m \in \mathbb{Z} \tag{2.26}
\end{equation*}
$$

Remark 2.23. Points of $O_{\ell}$ are semisimple. The results of 4, 5, 7, 8] imply that the monodromy data ( $\mu, R, S, C$ ) are well defined also at points $p \in \Omega_{s s} \cap \Delta_{\Omega}$, and that Theorem 2.21 still holds true.

Remark 2.24. From the knowledge of the monodromy data $(\mu, R, S, C)$ the Gromov-Witten potential $F_{0}^{X}(\boldsymbol{t})$ can be recostructed via a Riemann-Hilbert boundary value problem, see [5, 6, 13, 23]. Hence, the monodromy data may be interpreted as a system of coordinates in the space of solutions of $W D V V$ equations.

### 2.5. Action of the braid group $\mathcal{B}_{n}$

Consider the braid group $\mathcal{B}_{n}$ with generators $\beta_{1}, \ldots, \beta_{n-1}$ satisfying the relations

$$
\begin{gather*}
\beta_{i} \beta_{j}=\beta_{j} \beta_{i}, \quad|i-j|>1,  \tag{2.27}\\
\beta_{i} \beta_{i+1} \beta_{i}=\beta_{i+1} \beta_{i} \beta_{i+1} . \tag{2.28}
\end{gather*}
$$

Let $\mathcal{U}_{n}$ be the set of upper triangular $(n \times n)$-matrices with 1 's along the diagonal.

Definition 2.25. Given $U \in \mathcal{U}_{n}$ define the matrices $A^{\beta_{i}}(U)$, with $i=1, \ldots, n-$ 1 , as follows

$$
\begin{align*}
\left(A^{\beta_{i}}(U)\right)_{h h}:=1, \quad h & =1, \ldots, n, \quad h \neq i, i+1  \tag{2.29}\\
\left(A^{\beta_{i}}(U)\right)_{i+1, i+1} & =-U_{i, i+1},  \tag{2.30}\\
\left(A^{\beta_{i}}(U)\right)_{i, i+1} & =\left(A^{\beta_{i}}(U)\right)_{i+1, i}=1 \tag{2.31}
\end{align*}
$$

and all other entries of $A^{\beta_{i}}(U)$ are equal to zero.
Lemma 2.26 ([5, 11, 13]). The braid group $\mathcal{B}_{n}$ acts on $\mathcal{U}_{n} \times G L(n, \mathbb{C})$ as follows:

$$
\begin{aligned}
\mathcal{B}_{n} \times \mathcal{U}_{n} \times G L(n, \mathbb{C}) & \longrightarrow \mathcal{U}_{n} \times G L(n, \mathbb{C}) \\
\left(\beta_{i}, U, C\right) & \longmapsto\left(A^{\beta_{i}}(U) \cdot U \cdot A^{\beta_{i}}(U), C \cdot A^{\beta_{i}}(U)^{-1}\right)
\end{aligned}
$$

We denote by $(U, C)^{\beta_{i}}$ the action of $\beta_{i}$ on $(U, C)$.
Fix an oriented ray $\ell \equiv\{\arg z=\phi\}$ in $\widehat{\mathbb{C}^{*}}$, and denote by $\bar{\ell}$ its projection on $\mathbb{C}^{*}$. Let $\Omega_{\ell, 1}, \Omega_{\ell, 2}$ be two $\ell$-chambers and let $p_{i} \in \Omega_{\ell, i}$ for $i=1,2$. The difference of values of the Stokes and central connection matrices ( $S_{1}, C_{1}$ ) and $\left(S_{2}, C_{2}\right)$, at $p_{1}$ and $p_{2}$ respectively, can be described by the action of the braid group $\mathcal{B}_{n}$ of Lemma 2.26 .

Theorem 2.27 ([5, 11, 13). Consider a continuous path $\gamma:[0,1] \rightarrow \Omega$ such that

- $\gamma(0)=p_{1}$ and $\gamma(1)=p_{2}$,
- there exists a unique $t_{o} \in[0,1]$ such that $\ell$ is not admissible at $\gamma\left(t_{o}\right)$,
- there exist $i_{1}, \ldots, i_{k} \in\{1, \ldots, n\}$, with $\left|i_{a}-i_{b}\right|>1$ for $a \neq b$, such that the ray $\}^{3}\left(R_{i_{j}, i_{j}+1}(t)\right)_{j=1}^{r}\left(\right.$ resp. $\left.\left(R_{i_{j}, i_{j}+1}(t)\right)_{j=r+1}^{k}\right)$ cross the ray $\bar{\ell}$ in the clockwise (resp. counterclockwise) direction, as $t \rightarrow t_{o}^{-}$.

[^2]Then, we have

$$
\begin{equation*}
\left(S_{2}, C_{2}\right)=\left(S_{1}, C_{1}\right)^{\beta}, \quad \beta:=\left(\prod_{j=1}^{r} \beta_{i_{j}}\right) \cdot\left(\prod_{h=r+1}^{k} \beta_{i_{h}}\right)^{-1} \tag{2.32}
\end{equation*}
$$

Remark 2.28. In the general case, the points $p_{1}$ and $p_{2}$ can be connected by concatenations of paths $\gamma$ satisfying the assumptions of Theorem 2.27,

Remark 2.29. The action of $\mathcal{B}_{n}$ on $(S, C)$ also describes the analytic continuation of the Frobenius manifold structure on $\Omega$, see [13, Lecture 4].

## 3. Derived category, exceptional collections, helices

### 3.1. Notations and basic notions

Denote by $\operatorname{Coh}(X)$ the abelian category of coherent sheaves on $X$, and by $\mathcal{D}^{b}(X)$ its bounded derived category. Objects of $\mathcal{D}^{b}(X)$ are bounded complexes $A^{\bullet}$ of coherent sheaves on $X$. Morphisms are given by roofs: if $A^{\bullet}, B^{\bullet}$ are two bounded complexes, a morphism $f: A^{\bullet} \rightarrow B^{\bullet}$ in $\mathcal{D}^{b}(X)$ is the datum of

- a third object $C^{\bullet}$ in $\mathcal{D}^{b}(X)$,
- two homotopy classes of morphisms of complexes $q: C^{\bullet} \rightarrow A^{\bullet}$ and $g: C^{\bullet} \rightarrow B^{\bullet}$,
- the morphism $q$ is required to be a quasi-isomorphism, i.e. it induces isomorphism in cohomology.


The derived category $\mathcal{D}^{b}(X)$ admits a triangulated structure, the shift functor [1]: $\mathcal{D}^{b}(X) \rightarrow \mathcal{D}^{b}(X)$ being defined by

$$
\begin{equation*}
A^{\bullet}[1]:=A^{\bullet+1}, \quad A^{\bullet} \in \mathcal{D}^{b}(X) \tag{3.2}
\end{equation*}
$$

Denote by $\operatorname{Hom}^{\bullet}\left(A^{\bullet}, B^{\bullet}\right):=\bigoplus_{k \in \mathbb{Z}} \operatorname{Hom}\left(A^{\bullet}, B^{\bullet}[k]\right)$. General references for this Section are [17, 20, 21, 32].

### 3.2. Exceptional collections

Definition 3.1. An object $E \in \mathcal{D}^{b}(X)$ is called exceptional iff

$$
\begin{equation*}
\operatorname{Hom}^{\bullet}(E, E) \cong \mathbb{C} \tag{3.3}
\end{equation*}
$$

Definition 3.2. An exceptional collection is an ordered family $\left(E_{1}, \ldots, E_{n}\right)$ of exceptional objects of $\mathcal{D}^{b}(X)$ such that

$$
\begin{equation*}
\operatorname{Hom}^{\bullet}\left(E_{j}, E_{i}\right) \cong 0 \quad \text { for } j>i \tag{3.4}
\end{equation*}
$$

An exceptional collection is full if it generates $\mathcal{D}^{b}(X)$ as a triangulated category, i.e. if any full triangulated subcategory of $\mathcal{D}^{b}(X)$ containing all the objects $E_{i}$ 's is equivalent to $\mathcal{D}^{b}(X)$ via the inclusion functor.

Example. In [2] A. Beilinson showed that the collection of line bundles

$$
\begin{equation*}
\mathfrak{B}:=(\mathcal{O}, \mathcal{O}(1), \ldots, \mathcal{O}(n)) \tag{3.5}
\end{equation*}
$$

on $\mathbb{P}^{n}$ is a full exceptional collection. M. Kapranov generalized this result in [25], where full exceptional collections on Grassmannians, flag varieties of group $S L_{n}$, and smooth quadrics are constructed.

Denote by $\mathbb{G}(k, n)$ the Grassmannian of $k$-dimensional subspaces in $\mathbb{C}^{n}$, by $\mathcal{S}^{\vee}$ the dual of its tautological bundle. Let $\mathbb{S}^{\lambda}$ be the Schur functor (see [15]) labelled by a Young diagram $\lambda$ inside a rectangle $k \times(n-k)$. The collection $\mathfrak{K}:=\left(\mathbb{S}^{\lambda} \mathcal{S}^{\vee}\right)_{\lambda}$ is full and exceptional in $\mathcal{D}^{b}(\mathbb{G}(k, n))$. The order of the objects of the collection is the partial order defined by inclusion of Young diagrams.

### 3.3. Mutations and helices

Let $E$ be an exceptional object in $\mathcal{D}^{b}(X)$. For any $X \in \mathcal{D}^{b}(X)$, we have natural evaluation and co-evaluation morphisms

$$
\begin{equation*}
j^{*}: \operatorname{Hom}^{\bullet}(E, X) \otimes E \rightarrow X, \quad j_{*}: X \rightarrow \operatorname{Hom}^{\bullet}(X, E)^{*} \otimes E \tag{3.6}
\end{equation*}
$$

Definition 3.3. The left and right mutations of $X$ with respect to $E$ are the objects $\mathbb{L}_{E} X$ and $\mathbb{R}_{E} X$ uniquely defined by the distinguished triangles

$$
\begin{align*}
& \mathbb{L}_{E} X[-1] \longrightarrow \operatorname{Hom}^{\bullet}(E, X) \otimes E \xrightarrow{j^{*}} X \longrightarrow \mathbb{L}_{E} X,  \tag{3.7}\\
& \mathbb{R}_{E} X \longrightarrow X \xrightarrow{j_{*}} \operatorname{Hom}^{\bullet}(X, E)^{*} \otimes E \longrightarrow \mathbb{R}_{E} X[1], \tag{3.8}
\end{align*}
$$

respectively.
Remark 3.4. In general, the third object of a distinguished triangle is not canonically defined by the other two terms. Nevertheless, the objects $\mathbb{L}_{X} E$ and $\mathbb{R}_{E} X$ are uniquely defined up to unique isomorphism, because of the exceptionality of $E$, see [6, Section 3.3].

Definition 3.5. Let $\mathfrak{E}=\left(E_{1}, \ldots, E_{n}\right)$ be an exceptional collection. For any $i=1, \ldots, n-1$ define the left and right mutations

$$
\begin{align*}
& \mathbb{L}_{i} \mathfrak{E}:=\left(E_{1}, \ldots, \mathbb{L}_{E_{i}} E_{i+1}, E_{i}, \ldots, E_{n}\right),  \tag{3.9}\\
& \mathbb{R}_{i} \mathfrak{E}:=\left(E_{1}, \ldots, E_{i+1}, \mathbb{R}_{E_{i+1}} E_{i}, \ldots, E_{n}\right) . \tag{3.10}
\end{align*}
$$

Theorem 3.6 ( $[20, ~(32])$. For all $i=1, \ldots, n-1$ the collections $\mathbb{L}_{i} \mathfrak{E}$ and $\mathbb{R}_{i} \mathfrak{E}$ are exceptional. Moreover, we have that

$$
\begin{array}{ll}
\mathbb{L}_{i} \mathbb{R}_{i}=\mathbb{R}_{i} \mathbb{L}_{i}=\mathrm{Id}, & \mathbb{L}_{i+1} \mathbb{L}_{i} \mathbb{L}_{i+1}=\mathbb{L}_{i} \mathbb{L}_{i+1} \mathbb{L}_{i}, \quad i=1, \ldots, n \\
& \mathbb{L}_{i} \mathbb{L}_{j}=\mathbb{L}_{j} \mathbb{L}_{i}, \quad|i-j|>1
\end{array}
$$

According to Theorem [3.6] we have a well-defined action of $\mathcal{B}_{n}$ on the set of exceptional collections of length $n$ in $\mathcal{D}^{b}(X)$ : the action of the generator $\beta_{i}$ is identified with the action of the mutation $\mathbb{L}_{i}$ for $i=1, \ldots, n-1$.

Definition 3.7. Let $\mathfrak{E}=\left(E_{1}, \ldots, E_{n}\right)$ be a full exceptional collection. We define the helix generated by $\mathfrak{E}$ to be the infinite family $\left(E_{i}\right)_{i \in \mathbb{Z}}$ of exceptional objects obtained by iterated mutations

$$
E_{n+i}:=\mathbb{R}_{E_{n+i-1}} \ldots \mathbb{R}_{E_{i+1}} E_{i}, \quad E_{i-n}:=\mathbb{L}_{E_{i-n+1}} \ldots \mathbb{L}_{E_{i-1}} E_{i}, \quad i \in \mathbb{Z}
$$

Any family of $n$ consecutive exceptional objects $\left(E_{i+k}\right)_{k=1}^{n}$ is called a foundation of the helix.

Lemma 3.8 ([20]). For $i, j \in \mathbb{Z}$, we have $\operatorname{Hom}^{\bullet}\left(E_{i}, E_{j}\right) \cong \operatorname{Hom}^{\bullet}\left(E_{i-n}, E_{j-n}\right)$.

### 3.4. Exceptional bases in $K$-theory

Consider the Grothendieck group $K_{0}(X) \equiv K_{0}\left(\mathcal{D}^{b}(X)\right)$, equipped with the Grothendieck-Euler-Poincaré bilinear form

$$
\begin{equation*}
\chi([V],[F]):=\sum_{k}(-1)^{k} \operatorname{dim}_{\mathbb{C}} \operatorname{Hom}(V, F[i]), \quad V, F \in \mathcal{D}^{b}(X) \tag{3.11}
\end{equation*}
$$

Definition 3.9. A basis $\left(e_{i}\right)_{i=1}^{n}$ of $K_{0}(X)_{\mathbb{C}}$ is called exceptional if $\chi\left(e_{i}, e_{i}\right)=1$ for $i=1, \ldots, n$, and $\chi\left(e_{j}, e_{i}\right)=0$ for $1 \leq i<j \leq n$.

Lemma 3.10. Let $\left(E_{i}\right)_{i=1}^{n}$ be a full exceptional collection in $\mathcal{D}^{b}(X)$. The $K-$ classes $\left(\left[E_{i}\right]\right)_{i=1}^{n}$ form an exceptional basis of $K_{0}(X)_{\mathbb{C}}$.

The action of the braid group on the set of exceptional collections in $\mathcal{D}^{b}(X)$ admits a $K$-theoretical analogue on the set of exceptional bases of $K_{0}(X)_{\mathbb{C}}$, see [6, 20].

## 4. Dubrovin's conjecture

## 4.1. $\Gamma$-classes and graded Chern character

Let $V$ be a complex vector bundle on $X$ of rank $r$, and let $\delta_{1}, \ldots, \delta_{r}$ be its Chern roots, so that $c_{j}(V)=s_{j}\left(\delta_{1}, \ldots, \delta_{r}\right)$, where $s_{j}$ is the $j$-th elementary symmetric polynomial.

Definition 4.1. Let $Q$ be an indeterminate, and $F \in \mathbb{C} \llbracket Q \rrbracket$ be of the form $F(Q)=1+\sum_{n \geq 1} \alpha_{n} Q^{n}$. The $F$-class of $V$ is the charcateristic class $\widehat{F}_{V} \in$ $H^{\bullet}(X)$ defined by $\widehat{F}_{V}:=\prod_{j=1}^{r} F\left(\delta_{j}\right)$.

Definition 4.2. The $\Gamma^{ \pm}$-classes of $V$ are the characteristic classes associated with the Taylor expansions

$$
\begin{equation*}
\Gamma(1 \pm Q)=\exp \left(\mp \gamma Q+\sum_{m=2}^{\infty}(\mp 1)^{m} \frac{\zeta(m)}{m} Q^{n}\right) \in \mathbb{C} \llbracket Q \rrbracket, \tag{4.1}
\end{equation*}
$$

where $\gamma$ is the Euler-Mascheroni constant and $\zeta$ is the Riemann zeta function.
If $V=T X$, then we denote $\widehat{\Gamma}_{X}^{ \pm}$its $\Gamma$-classes.
Definition 4.3. The graded Chern character of $V$ is the characteristic class $\operatorname{Ch}(V) \in H^{\bullet}(X)$ defined by $\operatorname{Ch}(V):=\sum_{j=1}^{r} \exp \left(2 \pi \sqrt{-1} \delta_{j}\right)$.

### 4.2. Statement of the conjecture

Let $X$ be a Fano variety. In [12] Dubrovin conjectured that many properties of the $q D E$ of $X$, in particular its monodromy, Stokes and central connection matrices, are encoded in the geometry of exceptional collections in $\mathcal{D}^{b}(X)$. The following conjecture is a refinement of the original version in [12].

Conjecture 4.4 ([6]). Let $X$ be a smooth Fano variety of Hodge-Tate type.

1. The quantum cohomology $Q H^{\bullet}(X)$ has semisimple points if and only if there exists a full exceptional collection in $\mathcal{D}^{b}(X)$.
2. If $Q H^{\bullet}(X)$ is generically semisimple, for any oriented ray $\ell$ of slope $\phi \in[0,2 \pi[$ there is a correspondence between $\ell$-chambers and helices with a marked foundation.
3. Let $\Omega_{\ell}$ be an $\ell$-chamber and $\mathfrak{E}_{\ell}=\left(E_{1}, \ldots, E_{n}\right)$ the corresponding exceptional collection (the marked foundation). Denote by $S$ and $C$ Stokes and central connection matrices computed in $\Omega_{\ell}$.
(a) The matrix $S$ is the inverse of the Gram matrix of the $\chi$-pairing in $K_{0}(X)_{\mathbb{C}}$ wrt the exceptional basis $\left[\mathfrak{E}_{\ell}\right]$,

$$
\begin{equation*}
\left(S^{-1}\right)_{i j}=\chi\left(E_{i}, E_{j}\right) \tag{4.2}
\end{equation*}
$$

(b) The matrix $C$ coincides with the matrix associated with the $\mathbb{C}$ linear morphism

$$
\begin{align*}
& Д_{X}^{-}: K_{0}(X)_{\mathbb{C}} \longrightarrow H^{\bullet}(X)  \tag{4.3}\\
& F \longmapsto \frac{(\sqrt{-1})^{\bar{d}}}{(2 \pi)^{\frac{d}{2}}} \widehat{\Gamma}_{X}^{-} \exp \left(-\pi \sqrt{-1} c_{1}(X)\right) \operatorname{Ch}(F), \tag{4.4}
\end{align*}
$$

where $d:=\operatorname{dim}_{\mathbb{C}} X$, and $\bar{d}$ is the residue class $d(\bmod 2)$. The matrix is computed wrt the exceptional basis $\left[\mathfrak{E}_{\ell}\right]$ and the pre-fixed basis $\left(T_{\alpha}\right)_{\alpha=1}^{n}$ of $H^{\bullet}(X)$.

Remark 4.5. Conjecture 4.4 relates two different aspects of the geometry of $X$, namely its symptectic structure ( $G W$-theory) and its complex structure (the derived category $\mathcal{D}^{b}(X)$ ). Heuristically, Conjecture 4.4 follows from Homological Mirror Symmetry Conjecture of M. Kontsevich, see [6, Section 5.5].

Remark 4.6. In the paper [26] it was underlined the role of $\Gamma$-classes for refining the original version of Dubrovin's conjecture [12]. Subsequently, in [14] and [16, $\Gamma$-conjecture II] two equivalent versions of point (3.b) above were given. However, in both these versions, different choices of solutions in Levelt form of the $q D E$ at $z=0$ are chosen wrt the natural ones in the theory of Frobenius manifolds, see Remark 2.4, and [6, Section 5.6].

Remark 4.7. If point (3.b) holds true, then automatically also point (3.a) holds true. This follows from the identity (2.24) and Hirzebruch-RiemannRoch Theorem, see [6, Corollary 5.8].

Remark 4.8. Assume the validity of points (3.a) and (3.b) of Conjecture 4.4. The action of the braid group $\mathcal{B}_{n}$ on the Stokes and central connection matrices (Lemma 2.26) is compatible with the action of $\mathcal{B}_{n}$ on the marked foundations attached at each $\ell$-chambers. Different choices of the branch of the $\Psi$ matrix correspond to shifts of objects of the marked foundation. The matrix $M_{0}^{-1}$ is identified with the canonical operator $\kappa: K_{0}(X)_{\mathbb{C}} \rightarrow K_{0}(X)_{\mathbb{C}},[F] \mapsto$ $(-1)^{d}\left[F \otimes \omega_{X}\right]$. Equations (2.26) imply that the connection matrices $C^{(m)}$, with $m \in \mathbb{Z}$, correspond to the matrices of the morphism $Д_{X}^{-}$wrt the foundations $\left(\mathfrak{E}_{\ell} \otimes \omega_{X}^{\otimes m}\right)[m d]$. The statement $S^{(m)}=S$ coincides with the periodicity described in Lemma 3.8, see [6, Theorem 5.9].
Remark 4.9. Point (3.b) of Conjecture 4.4 allows to identify $K$-classes with solutions of the joint system of equations (2.1), (2.2). Under this identification, Stokes fundamental solutions correspond to exceptional bases of K theory. In the approach of [9, 33, where the equivariant case is addressed, such an identification is more fundamental and a priori, see Section 6 .

## 5. Results for Grassmannians

Conjecture 4.4 has been proved for complex Grassmannians $\mathbb{G}(k, n)$ in [6, 16]. See also [22, 34]. The proof is based on direct computation of the monodromy data of the $q D E$ at points of the small quantum cohomology, namely the subset $H^{2}(\mathbb{G}(k, n), \mathbb{C})$ of $\Omega$. Here we summarize the main results obtained.
Remark 5.1. $\mathbb{I}^{4} \pi_{1}(n) \leq k \leq n-\pi_{1}(n)$, the small quantum locus of $\mathbb{G}(k, n)$ is contained in the coalescence locus $\Delta_{\Omega}$, see [3]. In these cases, the computation of the monodromy data is justified by the results of [4, 5, 7, 8, See also Remark 2.23.

### 5.1. The case of projective spaces

Denote by $\sigma \in H^{2}\left(\mathbb{P}^{n-1}, \mathbb{C}\right)$ the hyperplane class and fix the basis $\left(\sigma^{k}\right)_{k=0}^{n-1}$ of $H^{\bullet}\left(\mathbb{P}^{n-1}\right)$. The joint system (2.1), (2.2) for $\mathbb{P}^{n-1}$, restricted at the point $t \sigma \in H^{2}\left(\mathbb{P}^{n-1}, \mathbb{C}\right)$, with $t \in \mathbb{C}$, is

$$
\begin{align*}
& \frac{\partial Z}{\partial t}=z \mathcal{C}(t) Z  \tag{5.1}\\
& \frac{\partial Z}{\partial z}=\left(\mathcal{U}(t)+\frac{1}{z} \mu\right) Z \tag{5.2}
\end{align*}
$$

with

$$
\begin{gather*}
\mathcal{U}(t)=\left(\begin{array}{ccccc}
0 & & & & n q \\
n & 0 & & & \\
& n & 0 & & \\
& & \ddots & \ddots & \\
& & & n & 0
\end{array}\right), \quad q:=e^{t}, \quad \mathcal{C}(t)=\frac{1}{n} \mathcal{U}(t),  \tag{5.3}\\
\end{gather*} \quad \begin{array}{ll} 
&  \tag{5.4}\\
& =\operatorname{diag}\left(-\frac{n-1}{2},-\frac{n-3}{2}, \ldots, \frac{n-3}{2}, \frac{n-1}{2}\right) .
\end{array}
$$

[^3]The canonical coordinates are given by the eigenvalues of the matrix $\mathcal{U}(t)$,

$$
\begin{equation*}
u_{h}(t)=n e^{\frac{2 \pi i(h-1)}{n}} q^{\frac{1}{n}} \quad h=1, \ldots, n . \tag{5.5}
\end{equation*}
$$

Fix the orthonormalized idempotent vector fields, $f_{1}(t), \ldots, f_{n}(t)$, given by $f_{h}(t):=\sum_{\ell=1}^{n} f_{h}^{\ell}(t) \sigma^{\ell-1}, \quad f_{h}^{\ell}(t):=n^{-\frac{1}{2}} q^{\frac{n+1-2 \ell}{2 n}} e^{(1-2 \ell) i \pi \frac{(h-1)}{n}} \quad h, \ell=1, \ldots, n$, and consider the following branch of the $\Psi$-matrix,

$$
\Psi(t):=\left(\begin{array}{c|c|c}
f_{1}^{1}(t) & \ldots & f_{n}^{1}(t)  \tag{5.6}\\
\vdots & & \vdots \\
f_{1}^{n}(t) & \ldots & f_{n}^{n}(t)
\end{array}\right)^{-1} .
$$

Theorem 5.2 (6]). Fix the oriented ray $\ell$ in $\widehat{\mathbb{C}^{*}}$ of slope $\phi \in\left[0, \frac{\pi}{n}[\right.$. For suitable choices of the signs of the columns of the $\Psi$-matrix (5.6), the central connection matrix computed at $0 \in H^{\bullet}\left(\mathbb{P}^{n-1}\right)$ coincides with the matrix attached to the morphism

$$
\text { Д}_{\mathbb{P}^{n-1}}^{-}: K_{0}\left(\mathbb{P}^{n-1}\right)_{\mathbb{C}} \rightarrow H^{\bullet}\left(\mathbb{P}^{n-1}\right)
$$

computed wrt the exceptional bases

$$
\begin{equation*}
\mathcal{O}\left(\frac{n}{2}\right), \bigwedge^{1} \mathcal{T}\left(\frac{n}{2}-1\right), \mathcal{O}\left(\frac{n}{2}+1\right), \bigwedge^{3} \mathcal{T}\left(\frac{n}{2}-2\right), \ldots, \mathcal{O}(n-1), \bigwedge^{n-1} \mathcal{T} \tag{5.7}
\end{equation*}
$$

for $n$ even, and

$$
\begin{align*}
& \mathcal{O}\left(\frac{n-1}{2}\right), \mathcal{O}\left(\frac{n+1}{2}\right), \bigwedge^{2} \mathcal{T}\left(\frac{n-3}{2}\right)  \tag{5.8}\\
& \mathcal{O}\left(\frac{n+3}{2}\right), \bigwedge^{4} \mathcal{T}\left(\frac{n-5}{2}\right), \ldots, \mathcal{O}(n-1), \bigwedge^{n-1} \mathcal{T}
\end{align*}
$$

for $n$ odd. In particular, Conjecture 4.4 holds true for $\mathbb{P}^{n-1}$.
Remark 5.3. Exceptional collections (5.7) and (5.8) are related to Beilinson's exceptional collection (3.5) by mutations and shifts. For different choices of the ray $\ell$, the exceptional collections attached to the monodromy data computed at $0 \in H^{\bullet}\left(\mathbb{P}^{n-1}\right)$ are given (up to shifts) by the following list, see [6, 9].

1. Case $n$ odd: an exceptional collection either of the form

$$
\begin{gathered}
\mathcal{O}\left(-k-\frac{n-1}{2}\right), \mathcal{T}\left(-k-\frac{n-1}{2}-1\right), \mathcal{O}\left(-k-\frac{n-1}{2}+1\right), \\
\bigwedge^{3} \mathcal{T}\left(-k-\frac{n-1}{2}-2\right), \mathcal{O}\left(-k-\frac{n-1}{2}+2\right), \ldots, \bigwedge^{n-4} \mathcal{T}(-k-n+2), \\
\mathcal{O}(-k-1), \bigwedge^{n-2} \mathcal{T}(-k-n+1), \mathcal{O}(-k),
\end{gathered}
$$

or of the form

$$
\begin{aligned}
& \quad \mathcal{O}\left(-k-\frac{n-1}{2}\right), \mathcal{O}\left(-k-\frac{n-1}{2}+1\right), \bigwedge^{2} \mathcal{T}\left(-k-\frac{n-1}{2}-1\right), \\
& \mathcal{O}\left(-k-\frac{n-1}{2}+2\right), \bigwedge^{3} \mathcal{T}\left(-k-\frac{n-1}{2}-2\right) \ldots, \mathcal{O}(-k-1), \\
& \bigwedge^{n-3} \mathcal{T}(-k-n+2), \mathcal{O}(-k), \bigwedge^{n-1} \mathcal{T}(-k-n+1),
\end{aligned}
$$

for some $k \in \mathbb{Z}$
2. Case $n$ even: an exceptional collection either of the form

$$
\begin{aligned}
\mathcal{O}\left(-k-\frac{n}{2}\right), & \mathcal{O}\left(-k-\frac{n}{2}+1\right), \bigwedge^{2} \mathcal{T}\left(-k-\frac{n}{2}-1\right), \mathcal{O}\left(-k-\frac{n}{2}+2\right), \ldots, \\
& \ldots, \bigwedge^{n-4} \mathcal{T}(-k-n+2), \mathcal{O}(-k-1), \bigwedge^{n-2} \mathcal{T}(-k-n+1), \mathcal{O}(-k)
\end{aligned}
$$

or of the form

$$
\begin{aligned}
& \mathcal{O}\left(-k-\frac{n}{2}+1\right), \mathcal{T}\left(-k-\frac{n}{2}\right), \mathcal{O}\left(-k-\frac{n}{2}+2\right), \bigwedge^{3} \mathcal{T}\left(-k-\frac{n}{2}-1\right), \ldots \\
& \ldots, \mathcal{O}(-k-1), \bigwedge^{n-3} \mathcal{T}(-k-n+2), \mathcal{O}(-k), \bigwedge^{n-1} \mathcal{T}(-k-n+1)
\end{aligned}
$$

for some $k \in \mathbb{Z}$.

### 5.2. The case of Grassmannians

Denote by $\mathbb{G}$ the Grassmannian $\mathbb{G}(k, n)$ parametrizing $k$-dimensional subspaces in $\mathbb{C}^{n}$, and by $\mathbb{P}$ the projective space $\mathbb{P}^{n-1}$. Let $\xi_{1}, \ldots, \xi_{k}$ be the Chern roots of the dual of the tautological bundle $\mathcal{S}$ on $\mathbb{G}$, and denote by $h_{j}(\boldsymbol{\xi})$ the $j$-th complete symmetric polynomial in $\xi_{1}, \ldots, \xi_{k}$. An additive basis of the cohomology ring

$$
\begin{equation*}
H^{\bullet}(\mathbb{G}) \cong \mathbb{C}\left[\xi_{1}, \ldots, \xi_{k}\right]^{\mathfrak{G}_{k}} /\left\langle h_{n-k+1}, \ldots, h_{n}\right\rangle \tag{5.9}
\end{equation*}
$$

is given by the Schubert classes $\left(\sigma_{\lambda}\right)_{\lambda \subseteq k \times(n-k)}$, labelled by partitions $\lambda$ with Young diagram inside a $k \times(n-k)$ rectangle. Under the presentation (5.9), the Schubert classes are given by Schur polynomials in $\boldsymbol{\xi}$,

$$
\begin{equation*}
\sigma_{\lambda}:=\frac{\operatorname{det}\left(\xi_{i}^{\lambda_{j}+k-j}\right)_{1 \leq i, j \leq k}}{\prod_{i<j}\left(\xi_{i}-\xi_{j}\right)} . \tag{5.10}
\end{equation*}
$$

Denote by $\eta_{\mathbb{P}}$ and $\eta_{\mathbb{G}}$ the Poincaré metrics on $H^{\bullet}(\mathbb{P})$ and $H^{\bullet}(\mathbb{G})$ respectively. The metric $\eta_{\mathbb{P}}$ induces a metric $\eta_{\mathbb{P}}^{\wedge^{k}}$ on the exterior power $\Lambda^{k} H^{\bullet}(\mathbb{P})$ :

$$
\begin{equation*}
\eta_{\mathbb{P}}^{\wedge^{k}}\left(\alpha_{1} \wedge \ldots, \wedge \alpha_{k}, \beta_{1} \wedge \ldots, \wedge \beta_{k}\right):=\operatorname{det}\left(\eta_{\mathbb{P}}\left(\alpha_{i}, \beta_{j}\right)\right)_{1 \leq i, j \leq k} \tag{5.11}
\end{equation*}
$$

Theorem 5.4 ([6, 16]). We have a $\mathbb{C}$-linear isometry

$$
\mathcal{I}:\left(\bigwedge^{k} H^{\bullet}(\mathbb{P}),(-1)^{\binom{k}{2}} \eta_{\mathbb{P}}^{\wedge^{k}}\right) \rightarrow\left(H^{\bullet}(\mathbb{G}), \eta_{\mathbb{G}}\right), \quad \sigma^{\nu_{1}} \wedge \cdots \wedge \sigma^{\nu_{k}} \mapsto \sigma_{\tilde{\nu}}
$$

where $n-1 \geq \nu_{1}>\nu_{2}>\cdots>\nu_{k} \geq 0$ and $\tilde{\nu}:=\left(\nu_{1}-k+1, \nu_{2}-k+2, \ldots, \nu_{k}\right)$.
Consider the domain $\Omega_{\mathbb{G}} \subset H^{\bullet}(\mathbb{G})\left(\right.$ resp. $\Omega_{\mathbb{P}} \subset H^{\bullet}(\mathbb{P})$ ) where the $G W$ potential $F_{0}^{\mathbb{G}}$ (resp. $\left.F_{0}^{\mathbb{P}}\right)$ converges. Let $t \in \mathbb{C}$ and consider the points

$$
\begin{equation*}
p:=t \sigma_{1} \in H^{2}(\mathbb{G}, \mathbb{C}), \quad \hat{p}:=(t+\pi \sqrt{-1}(k-1)) \sigma \in H^{2}(\mathbb{P}, \mathbb{C}) \tag{5.12}
\end{equation*}
$$

in the small quantum cohomology of $\mathbb{G}$ and $\mathbb{P}$ respectively. Theorem 5.4allow us to identify ${ }^{5}$ the tangent spaces $T_{p} \Omega_{\mathbb{G}}$ and $\bigwedge^{k} T_{\hat{p}} \Omega_{\mathbb{P}}$.
Lemma 5.5 ([6, [16]). Let $\Psi^{\mathbb{P}}(t)$ be the $\Psi$-matrix defined by (5.6). Then the matrix $\Psi^{\mathbb{G}}(t):=(\sqrt{-1})^{\binom{k}{2}} \bigwedge^{k} \Psi^{\mathbb{P}}(t+\pi \sqrt{-1}(k-1))$ defines a branch of the $\Psi$-matrix for $\mathbb{G}$.

[^4]The following results show that under the identification of Theorem 5.4. solutions and monodromy data of the joint system (2.1), (2.2) for $\mathbb{G}$ can be reconstructed from solutions for the joint system for $\mathbb{P}$.

Theorem 5.6 (6]). Let $Z^{\mathbb{P}}(t, z)$ be a solution of the joint system (5.1), (5.2). The function

$$
\begin{equation*}
Z^{\mathbb{G}}(t, z):=\bigwedge^{k}\left(Z^{\mathbb{P}}(t+\pi \sqrt{-1}(k-1), z)\right) \tag{5.13}
\end{equation*}
$$

is a solution for the joint system for $\mathbb{G}$, namely

$$
\begin{align*}
\frac{\partial Z^{\mathbb{G}}}{\partial t} & =z \mathcal{C}_{\mathbb{G}}(t) Z^{\mathbb{G}}  \tag{5.14}\\
\frac{\partial Z^{\mathbb{G}}}{\partial z} & =\left(\mathcal{U}_{\mathbb{G}}(t)+\frac{1}{z} \mu_{\mathbb{G}}\right) Z^{\mathbb{G}} . \tag{5.15}
\end{align*}
$$

Corollary 5.7 ([6]). Fix an oriented ray $\ell$ in $\widehat{\mathbb{C}^{*}}$ admissible at both points $p, \hat{p}$ in (5.12). Denote by $S^{\mathbb{P}}(\hat{p}), S^{\mathbb{G}}(p)$ and $C^{\mathbb{P}}(\hat{p}), C^{\mathbb{G}}(p)$ the Stokes and central connection matrices at $\hat{p}$ and $p$, respectively. We have

$$
\begin{align*}
S^{\mathbb{G}}(p) & =\bigwedge^{k} S^{\mathbb{P}}(\hat{p}),  \tag{5.16}\\
C^{\mathbb{G}}(p) & =(\sqrt{-1})^{-\binom{k}{2}}\left(\bigwedge^{k} C^{\mathbb{P}}(\hat{p})\right) \exp \left(\pi \sqrt{-1}(k-1) \sigma_{1} \cup\right) \tag{5.17}
\end{align*}
$$

Proof. Denote by

- $Z_{\text {top }}^{\mathbb{P}}(t, z)$ and $Z_{\text {top }}^{\mathbb{G}}(t, z)$ the topological-enumerative solutions for $\mathbb{P}$ and $\mathbb{G}$ respectively, restricted at their small quantum cohomologies;
- $Z_{L / R, m}^{\mathbb{P} / \mathbb{G}}(t, z)$, with $m \in \mathbb{Z}$, the Stokes fundamental solutions of the joint systems (2.1), (2.2) for $\mathbb{P}$ and $\mathbb{G}$ respectively.
We have

$$
\begin{aligned}
& Z_{\mathrm{top}}^{\mathbb{G}}(t, z)=\left(\bigwedge^{k} Z_{\mathrm{top}}^{\mathbb{P}}(t+\pi \sqrt{-1}(k-1), z)\right) \cdot \exp \left(-\pi \sqrt{-1}(k-1) \sigma_{1} \cup\right) \\
& \quad Z_{L / R, m}^{\mathbb{G}}(t, z)=(\sqrt{-1})^{-\binom{k}{2}} \bigwedge^{k} Z_{L / R, m}^{\mathbb{P}}(t+\pi \sqrt{-1}(k-1), z)
\end{aligned}
$$

See [6] for proofs of these identities.
Corollary 5.8 ([6]). The central connection matrix computed at $0 \in H^{\bullet}(\mathbb{G})$ coincides with the matrix attached to the morphism

$$
\text { Д }_{\mathbb{G}}^{-}: K_{0}(\mathbb{G})_{\mathbb{C}} \rightarrow H^{\bullet}(\mathbb{G})
$$

computed wrt an exceptional basis of $K_{0}(\mathbb{G})_{\mathbb{C}}$. Such a basis is the projection in $K$-theory of an exceptional collection of $\mathcal{D}^{b}(\mathbb{G})$ related by mutations and shifts to the twisted Kapranov excptional collection

$$
\begin{equation*}
\left(\mathbb{S}^{\lambda} \mathcal{S}^{\vee} \otimes \mathcal{L}\right), \quad \mathcal{L}:=\operatorname{det}\left(\bigwedge^{2} \mathcal{S}^{\vee}\right) \tag{5.18}
\end{equation*}
$$

In particular, Conjecture 4.4 holds true for $\mathbb{G}$.

## 6. Results on the equivariant $q D E$ of $\mathbb{P}^{n-1}$

Gromov-Witten theory, as described in Section 1.2, can be suitably adapted to the equivariant case [18]. Given a variety $X$ equipped with the action of a group $G$, a quantum deformation of the equivariant cohomology algebra $H_{G}^{\bullet}(X, \mathbb{C})$ can be defined.

Consider the projective space $\mathbb{P}^{n-1}$ equipped with the diagonal action of the torus $\mathbb{T}:=\left(\mathbb{C}^{*}\right)^{n}$. Although the isomonodromic system (5.1), (5.2) does not admit an equivariant analog, the differential equation (5.1) only can be easily modified. By change of coordinates $q:=\exp (t)$, setting $z=1$, and replacing the quantum multiplication $*_{q}$ by the corresponding equivariant one $*_{q, \boldsymbol{z}}$, equation (5.1) takes the form

$$
\begin{equation*}
q \frac{d}{d q} Z=\sigma *_{q, \boldsymbol{z}} Z \tag{6.1}
\end{equation*}
$$

Here the equivariant parameters $\boldsymbol{z}=\left(z_{1}, \ldots, z_{n}\right)$ correspond to the factors of $\mathbb{T}$, and $Z(q, \boldsymbol{z})$ takes values in $H_{\mathbb{T}}^{\bullet}\left(\mathbb{P}^{n-1}, \mathbb{C}\right)$. Equation (6.1) admits a compatible system of difference equations, called $q K Z$ difference equations

$$
\begin{equation*}
Z\left(q, z_{1}, \ldots, z_{i}-1, \ldots, z_{n}\right)=K_{i}(q, \boldsymbol{z}) Z(q, \boldsymbol{z}), \quad i=1, \ldots, n \tag{6.2}
\end{equation*}
$$

for suitable linear operators $K_{i}$ 's, introduced in [33]. The joint system (6.1), (6.2) is a suitable limit of an analogue one for the cotangent bundle $T^{*} \mathbb{P}^{n-1}$, see [19, 30]. The existence and compatibility of such a joint system for more general Nakajima quiver varieties is justified by the general theory of D. Maulik and A. Okounkov [28].

In [33], the study of the monodromy and Stokes phenomenon at $q=$ $\infty$ of solutions of the joint system (6.1), (6.2) is addressed. Furthermore, elements of $K_{0}^{\mathbb{T}}\left(\mathbb{P}^{n-1}\right)_{\mathbb{C}}$ are identified with solutions of the joint system (6.1), (6.2): Stokes bases of solutions correspond to exceptional bases.

In (9), the authors describe relations between the monodromy data of the joint system of the equivariant $q D E$ (6.1) and $q K Z$ equations (6.2) and characteristic classes of objects of the derived category $\mathcal{D}_{\mathbb{T}}^{b}\left(\mathbb{P}^{n-1}\right)$ of equivariant coherent sheaves on $\mathbb{P}^{n-1}$. Equivariant analogs of results of [6, Section 6] are obtained.

The Б-Theorem of [9] is the equivariant analog of Theorem 5.2 Moreover, in 9 the Stokes bases of solutions of the joint system (6.1), (6.2) are identified with explicit $\mathbb{T}$-full exceptional collections in $\mathcal{D}_{\mathbb{T}}^{b}\left(\mathbb{P}^{n-1}\right)$, which project to those listed in Remark 5.3 via the forgetful functor $\mathcal{D}_{\mathbb{T}}^{b}\left(\mathbb{P}^{n-1}\right) \rightarrow$ $\mathcal{D}^{b}\left(\mathbb{P}^{n-1}\right)$. This refines results of 33. Finally, in 9 it is proved that the Stokes matrices of the joint system (6.1), (6.2) equal the Gram matrices of the equivariant Grothendieck-Euler-Poincaré pairing on $K_{0}^{\mathbb{T}}\left(\mathbb{P}^{n-1}\right)_{\mathbb{C}}$ wrt the same exceptional bases.

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[^0]:    ${ }^{1}$ Here $h^{p, q}(X):=\operatorname{dim}_{\mathbb{C}} H^{q}\left(X, \Omega_{X}^{p}\right)$, with $\Omega_{X}^{p}$ the sheaf of holomorphic $p$-forms on $X$, denotes the $(p, q)$-Hodge number of $X$.

[^1]:    ${ }^{2}$ We consider the joint system (2.1), (2.2) in matrix notations ( $\zeta$ a column vector whose entries are the components $\zeta^{\alpha}(\boldsymbol{t}, z)$ wrt $\left.\frac{\partial}{\partial t^{\alpha}}\right)$. Bases of solutions are arranged in invertible $n \times n$-matrices, called fundamental systems of solutions.

[^2]:    ${ }^{3}$ Here the labeling of Stokes rays is the one prolonged from the initial point $t=0$.

[^3]:    ${ }^{4}$ Here $\pi_{1}(n)$ denotes the smallest prime number which divides $n$.

[^4]:    ${ }^{5}$ In what follows, if $A$ is a $n \times n$-matrix, we denote by $\bigwedge^{k} A$ the matrix of $k \times k$-minors of $A$, ordered in lexicographical order.

