Initial Data and Eccentricity Reduction Toolkit for Binary Black Hole Numerical Relativity Waveforms

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Abstract. The production of numerical relativity waveforms that describe quasi-circular binary black hole mergers requires high-quality initial data, and an algorithm to iteratively reduce residual eccentricity. To date, these tools remain closed source, or in commercial software that prevents their use in high performance computing platforms. To address these limitations, and to ensure that the broader numerical relativity community has access to these tools, herein we provide all the required elements to produce high-quality numerical relativity simulations in supercomputer platforms, namely: open source parameter files to numerical simulate spinning black hole binaries with asymmetric mass-ratios; open source `Python` tools to produce high-quality initial data for numerical relativity simulations of spinning black hole binaries on quasi-circular orbits; open source `Python` tools for eccentricity reduction, both as stand-alone software and deployed in the `Einstein Toolkit`'s software infrastructure. This open source toolkit fills in a critical void in the literature at a time when numerical relativity has an ever increasing role in the study and interpretation of gravitational wave sources. As part of our community building efforts, and to streamline and accelerate the use of these resources, we provide tutorials that describe, step by step, how to obtain and use these open source numerical relativity tools.

1. Introduction

Numerical relativity [1, 2, 3, 4] plays a central role in contemporary gravitational wave astrophysics [5, 6, 7, 8]. The use of numerical relativity waveforms has been essential to develop approximate waveform models that are extensively used for gravitational wave detection and parameter estimation [9, 10, 11, 12, 13]. The construction of numerical relativity waveforms catalogs [14, 15, 16, 17, 18, 19] has enabled in-depth analyses of the astrophysical properties of gravitational wave sources [20, 21, 22, 23].

As gravitational wave astrophysics continues to probe the gravitational wave spectrum [24, 25, 26, 27, 28], numerical relativity will be essential to enable and interpret new discoveries, enlighten our understanding of the physics of these sources, and provide constraints that may further establish general relativity or favor alternatives theories of gravity [29, 30, 31].

Advancing our understanding of gravitational wave sources depends critically on the production of high quality numerical relativity waveforms that, in the case of binary black hole mergers, span an 8-D parameter space that includes mass-ratio, two 3-D vectors that define the individual spin of the binary components, and orbital eccentricity, $(q, \mathbf{s}_1, \mathbf{s}_2, e)$, respectively. It is then apparent that despite the existence of thousands of numerical relativity waveforms, we need to be creative about how to combine them to densely sample these high dimensional signal manifold [32, 33]. It is also clear that we need to continue producing numerical relativity waveforms to describe sources whose parameters are not accurately captured by existing approximate waveform models or available numerical relativity waveforms.

In order to empower the broader numerical relativity community to participate in the construction of numerical relativity waveform catalogs, we introduce open source `Python` libraries that have been tested and deployed within the `Einstein Toolkit` [34] to streamline and accelerate these activities. This approach builds upon our previous software development that consisted of open source `Python` libraries to post-process numerical relativity data to extract the waveform strain at future null infinity [35]. These combined tools provide the required end-to-end software infrastructure to utilize the `Einstein Toolkit` for the construction of high-quality numerical relativity waveform catalogs.

This manuscript is organized as follows. Section 2 describes our approach to construct high-quality initial data, and to post-process the data products of numerical relativity simulations to remove residual eccentricity. We put these tools at work in Section 3, where we show that we can produce nearly circular initial data, and that our method for eccentricity reduction produces waveforms with eccentricities of order $\mathcal{O} \sim 10^{-4}$ after just one iteration. We summarize this work, and outline future research directions in Section 4. We present a tutorial that describes how to obtain and use these libraries in Appendix A.

2. Methods

In this section we describe the approach followed to produce high-quality initial data for binary black hole simulations. Thereafter, we briefly introduce the method used for eccentricity reduction.

2.1. Initial Data Production

The first guess for the tangential, p_t, and radial, p_r, components of momenta for the black hole binary system are generated using techniques presented in [36] and [37]. They extract momenta components from Hamilton's equations of motion in post-Newtonian (PN) theory, combined with high-PN-order expressions for the gravitational-wave flux, dE_{GW}/dt, and the tidal energy injected into the black holes, dM/dt.

The Hamiltonian contains orbital [38], spin-orbit [38, 39, 40], spin-spin [38, 41, 42], and spin-spin-spin [43] terms up to and including 3.5PN order.

The high-PN-order expressions for dE_{GW}/dt incorporate nonspinning and precessing-spin terms [44, 45], and are adjusted to account for the tidal energy injected into the black holes dM/dt [46].

The above expressions were implemented in the open-source, Python-based `NRPyPN` software, which is part of `NRPy+` ("Python-based code generation for numerical relativity... and beyond!") [47]. A tutorial for using the software is given in Appendix A.1 below. In short, the expression for tangential momentum $p_t(r)$ up to and including 3.5PN order is taken from [36] and validated up to 3PN order against [37], and up to 3.5PN order against the original `Mathematica` notebooks used by [36].

Meanwhile, the expression for radial momentum p_r up to and including 3.5PN order is derived in `NRPyPN` as follows. First, Hamilton's equations of motion imply that

$$\frac{dr}{dt} = \frac{\partial H}{\partial p_r}. \tag{1}$$

Next we Taylor expand $\partial H/\partial p_r$ in powers of p_r, about $p_r = 0$, to obtain (to first order in p_r):

$$p_r \approx \left(\frac{dr}{dt} - \frac{\partial H}{\partial p_r}\bigg|_{p_r=0} \right) \left(\frac{\partial^2 H}{\partial p_r^2}\bigg|_{p_r=0} \right)^{-1}, \tag{2}$$

where

$$\frac{dr}{dt} = \left(\frac{dE_{\mathrm{GW}}}{dt} + \frac{dM}{dt} \right) \left[\frac{dH_{\mathrm{circ}}}{dr} \right]^{-1}, \tag{3}$$

and

$$\frac{dH_{\mathrm{circ}}(r, p_t(r))}{dr} = \frac{\partial H(p_r = 0)}{\partial r} + \frac{\partial H(p_r = 0)}{\partial p_t}\frac{\partial p_t}{\partial r}, \tag{4}$$

are given explicitly in terms of binary input parameters and $M\Omega$ (as given to 3.5PN order by [36]).

2.2. Eccentricity Reduction

The algorithm we describe in this section was introduced in [36], and was originally developed as a `Mathematica` notebook. As part of this work, we have re-written this eccentricity reduction method using `Python` libraries, optimized it, and tailored it to conduct automated, large-scale, numerical relativity campaigns on high performance computing platforms.

This eccentricity reduction method is applied to remove eccentricity from a numerical relativity simulation whose initial data were produced with the method described in the previous section. Once the numerical simulation has progressed enough, typically between 500M to 600M of evolution, we process the relevant data files, as described in Appendix A, to compute correction factors, λ_t, λ_r, of the initial components of the momenta p_t^0 and p_r^0.

To compute λ_t, λ_r we assume that oscillations induced as a result of eccentricity in the orbital frequency, Ω, take the form

$$\mathcal{R}_\Omega = A + B\cos(\Omega_r t + \Psi),\tag{5}$$

where A, B and Ψ are coefficients to be determined, and Ω_r is the frequency of the radial oscillations. Using the 1PN order quasi-Keplerian parametrization [48], we can obtain closed form expressions for these correction factors

$$\lambda_t = 1 + \left[\frac{B}{4\Omega_0} - \gamma\frac{B(3\eta+1)}{8r_0\Omega_0}\right]\cos\Psi,\tag{6}$$

$$\lambda_r = 1 + \frac{B\eta}{2r_0^{1/2}\Omega_0\,|p_r^0|}\left[1 + \gamma\frac{1}{r_0}\right]\sin\Psi,\tag{7}$$

where $\eta = m_1 m_2/(m_1+m_2)$ is the symmetric mass ratio, (m_1, m_2) represent the masses of the binary components, r_0 is the initial orbital separation, and Ω_0 is the quasi-circular initial orbital frequency calculated at 3.5PN order [36]. In Appendix A.1 we describe how to use simulation data and analytical approximations to compute the correction factors λ_t and λ_r.
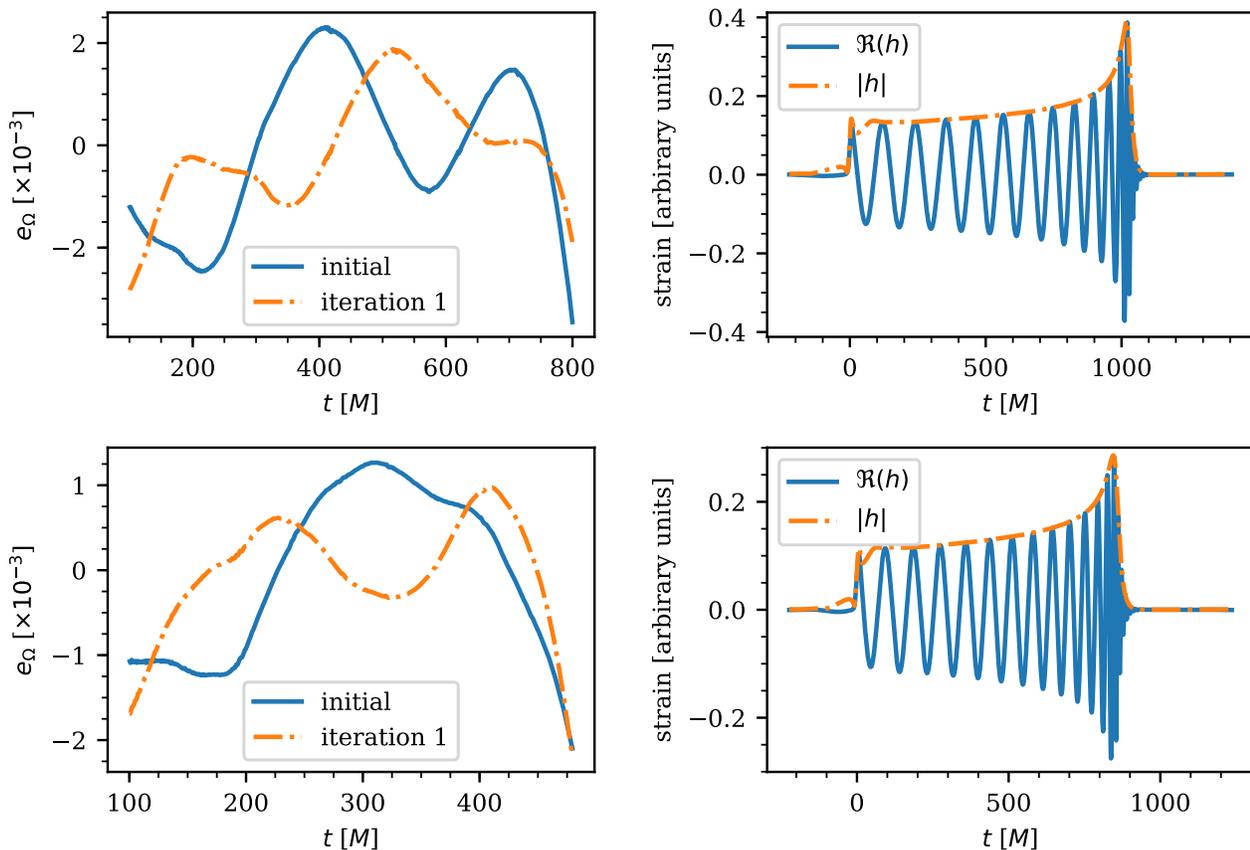
3. Results

In this section we combine the tools described above for initial data production and eccentricity reduction. We selected three binary black hole systems whose properties are described in Table 3. Notice that these systems span three different mass-ratios, $q \in \{1, 3\}$, and several spinning, non-precessing configurations.

The results presented in Table 3 show that for all the binary systems under consideration, our toolkit produces systems whose initial eccentricities are $e_0 \sim 10^{-3}$. Furthermore, these eccentricity values are reduced to $e_0 \sim 10^{-4}$ after just one iteration. In other words, these ready-to-use tools produce high-quality numerical relativity waveforms after a minimal number of iterations.

D [M]	$(\chi_1^x, \chi_1^y, \chi_1^z)$	$(\chi_2^x, \chi_2^y, \chi_2^z)$	q	**Iter** #	p_r [$\times 10^{-4}$]	p_ϕ [$\times 10^{-2}$]	e_0 [$\times 10^{-3}$]
11.0	(0.0, 0.0, -0.4)	(0.0, 0.0, -0.5)	1.0	**0**	-8.60	9.293	2.43
S_q_1				1	-7.70	9.284	0.74
9.0	(0.0, 0.0, 0.4)	(0.0, 0.0, -0.5)	3.0	**0**	-7.50	7.652	1.70
S_q_3				1	-6.60	7.650	0.71

Table 1. Summary of the astrophysical and orbital parameters of three binary black hole systems used to test our open source toolkit for the production of initial data and eccentricity reduction. Notice that in all cases the initial eccentricity is of order $e_0 \sim \mathcal{O}\left(10^{-3}\right)$ for the zeroth iteration, and it is reduced to $e_0 \sim \mathcal{O}\left(10^{-4}\right)$ after just one iteration.

Figure 1. Top panels: results of the eccentricity reduction procedure for simulation **S_q_1**. Top-left panel: eccentricity estimator results for the zeroth and first iterations. Top-right panel: gravitational wave signal, extracted at future null infinity, after eccentricity reduction. Bottom panels: as top panels, but for simulation **S_q_3**.

Figure 1 present two types of results. The left panels present results for the eccentricity estimator e_Ω of the orbital frequency ad defined in Eq. (3.13) of [36]. These results show, as discussed previously, that even the zeroth iteration is already nearly circular. The right panels present waveforms of the first iteration extracted at future null infinity [35]. These results indicate that the open source tools presented in this article, along with the tutorials and configuration files released with this work, will provide the required

building blocks to engage a broader cross section of the numerical relativity community in the production of large scale numerical relativity waveform catalogs.

4. Conclusions

Numerical relativity simulations [1, 2, 3] of binary black hole mergers were produced a decade before the first gravitational wave detection of these astrophysical events was realized by the advanced LIGO detectors [24]. Over the last decade, numerical relativity software stacks have matured to the point of automating and streamlining the production of large-scale numerical relativity catalogs [14, 15, 16, 17, 18, 19]. Nonetheless, the available number of numerical relativity waveforms is not sufficient to densely cover the high dimensional signal manifold spanned by these astrophysical events.

Furthermore, essential tools to produce initial data and to automate eccentricity reduction continue to be kept as closed source software or licensed software. Neither of these solutions is adequate if we aim to enable a larger cross section of the numerical relativity community to participate in the production of numerical relativity waveforms to accurately infer the astrophysical properties of compact binary sources. This need will become a pressing issue as advanced gravitational wave detectors gradually reach design sensitivity, and the number of detections reaches the expected number of one event for every fifteen minutes of searched data.

The deployment of these tools as stand-alone software and within the Einstein Toolkit is aligned with our community building efforts, and marks another milestone in our program for the production of an end-to-end software framework that enables users to produce high quality initial data, automate eccentricity reduction, and post-process numerical relativity data products to extract numerical relativity waveforms at future null infinity [35]. These user-friendly tools will allow new users to engage in the development of open source numerical relativity software, using the `Einstein Toolkit` as the driver for such community building activities.

5. Acknowledgements

Appendix A. Step by Step Tutorial

Appendix A.1. Initial Data Production

NRPyPN is part of the open-source NRPy+ ("Python-based code generation for numerical relativity... and beyond!") [47], and provides the zeroth estimate for low-eccentricity initial data in this paper. To obtain this estimate from NRPyPN, first clone the NRPy+ github repository:

    git clone https://github.com/zachetienne/nrpytutorial.git

    Then (assuming that Python 2 or 3 is installed with pip), install SymPy [49] and Jupyter:

```
pip install -U sympy jupyter
```

    Next, from within the nrpytutorial/NRPyPN/ directory, run

```
jupyter notebook NRPyPN.ipynb
```

    A Jupyter notebook will open, in which the binary black hole initial data parameters for initial separation, spins, and mass ratio can be specified in the code cell at the bottom of the notebook. When the code cell is run (Shift+Enter), the radial p_r and tangential momenta p_t to 3.5 post-Newtonian order (largely following [36] but fully documented in the linked Jupyter notebooks) will be output. These momenta can be directly inserted into a Bowen-York binary black hole initial data solver (the TwoPunctures thorn was used in this work). For example, for a binary orbiting in the xy-plane with black holes initially located on the x-axis at $x = \pm a$ (with the center of mass at the origin, $x = y = z = 0$), $\pm|p_r|$ will correspond to the x-component of momentum p_x for the puncture at $x = \mp a$, respectively. Also one may choose $\pm|p_t|$ to correspond to the y-component of momentum p_y for the puncture at $x = \pm a$ or $\mp a$, depending on whether a clockwise or a counterclockwise orbit is desired.

Appendix A.2. Eccentricity Reduction

Our eccentricity code is a available as a Python3 module on GitHub [50]

```
git clone https://github.com/ncsagravity/eccred
```

    Then, assuming Python 3 is installed, run:

```
python
>>> import EccRed
>>> EccRed.ComputeCorrections("output_glob", MinTime=X, MaxTime=Y)
```

where *output_ glob* is a shell pattern (glob) that matches all directories containing output files, MinTime and MaxTime are time bounds. For best results, MinTime should be shortly after any "junk" radiation has pass from the vicinity of the black holes and any initial gauge transition has settled, MaxTime should be close to the time plunge occurs.

    Four correction values as well as the estimated eccentricity will be returned from EccRed.ComputeCorrections. In order, they are λ_r, λ_t computed using two different

methods (from PN expansion and from an eccentricity estimator respectively), and δR, the correction factors to radial and tangential momentum components and (additive) correction to initial orbital separation respectively. These corrections can then be applied to the respective initial values.

The code expects to two sets of files in the output directories: (i) a file *TwoPunctures.bbh* as produced by the `TwoPunctures` thorn that describes the parameters of the initial black holes, and (ii) a set of puncture location files *puncturetracker-pt_loc..asc* as produced by the `PunctureTracker` thorn. Columns *pt_loc_x[0]*, *pt_loc_x[1]* etc., are expected to contain the location of the original `plus` and `minus` punctures. This matches the setup in [51].

In the event that `EccRed.ComputeCorrections` throws a runtime error, a likely solution is to adjust `MinTime` or `MaxTime` to better characterize the time domain of inspiral.

Appendix A.3. Automated Eccentricity Reduction

To simplify automation the process of eccentricity reduction using `Simulation Factory` [52] the `Python` module can be called as a command line script

```
./EccRed.py --tmin X --tmax X --input-parfile "input_parameter_file" \
            --output-parfile "output_parameter_file" "output_glob"
```

which automatically applies the correction factors to `TwoPunctures`' radial and tangential momentum parameters found in *input_parameter_file* and produces a new parameter file in *output_parameter_file*.

We provide a fragment of code in `RunScript.part` that can be inserted into `Simulation Factory`'s run script files to automate the process of extending a simulation until sufficiently much data has been produced, estimating eccentricity, computing correction factors, applying them to the parameter file and submitting a new round of eccentricity reduction.

The fragment contains placeholders *@ECC_TARGET@* and *@ECC_TIME@* for the estimated eccentricity at which to stop the iteration and the time for which to simulate before applying the correction algorithm:

```
sim create --define ECC_TARGET "ECC_TARGET" --define ECC_TIME "ECC_TIME" ...
```

which starts the automated process.

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