

Neutron Production during Deuterium Injection in Deuterium Plasmas in ASDEX

K. Hübner, R. Bätzner, H. Hinsch

Institut für Angewandte Physik, Universität Heidelberg, D-6900 Heidelberg

H. Rapp, H. Wurz and A. Eberhagen, O. Gehre, V. Mertens

ASDEX Team<sup>†</sup>, Neutral Injection Team<sup>†</sup>

Max-Planck-Institut für Plasmaphysik, EURATOM Association, D-8046 Garching

Measurements of neutron energy spectra and time dependent neutron rate are presented. They agree well with the results of a simple classical deuteron relaxation model. Determination of the plasma deuterium temperature seems possible from the energy spectra as well as from the neutron rate.

Energy spectra: The energy of neutrons produced by reactions between fast injected deuterons and plasma deuterons is shifted from 2.45 MeV towards higher energies for forward emission and towards lower energies for backward emission. On ASDEX with 45 keV injection energy, neutron energies of 2.70 and 2.24 MeV, respectively, are expected (broken lines in Fig. 1). Spectral measurements by means of nuclear emulsions along tangential lines of sight in the co- and counter direction of the injection beam show shifted peaks with maxima at 2.65 and 2.32 MeV (Fig. 1). Due to integration over a finite

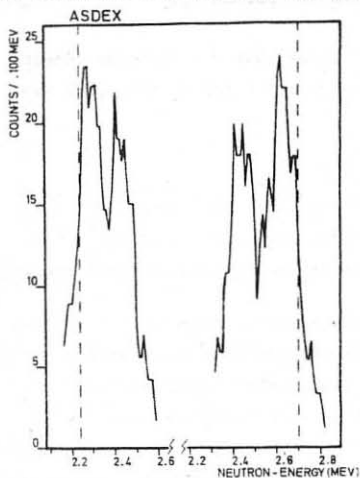


Fig. 1 Neutron energy spectra

space angle of neutron emission the measured spectra cannot show the maximum energy shift.

The 2.45 MeV line is caused by thermonuclear reactions in the bulk plasma. From its FWHM we obtain a deuteron temperature of 1.9 keV integrated over the observation volume and the injection time. Using the measured electron density profile  $n_e(r)/n_e(0)$  for deuteron density and temperature profiles,  $n_D(r)/n_D(0)$  and  $T_1(r)/T_1(0)$ , we calculate  $T_1(0) = 2.0$  keV. This value is about 10 % above the

<sup>†</sup>The members of the ASDEX and NI teams are presented in the paper of F. Wagner et al., this conference.

central electron temperature. The error of the ion temperature determination is at least 20 % in this case.

Neutron rate, classical relaxation model: The energy of an injected fast deuteron with initial energy  $W_0$  decreases in a plasma as

$W(t) = W_0 \exp(-t/\tau_w)$ .  $\tau_w$  is the energy relaxation time. The number of neutrons  $Y_D$  produced by one fast deuteron during its relaxation is given by

$$Y_D = \int_0^{\tau} n_D \sigma \sqrt{\frac{2W}{m}} dt = (n_D/n) \int_{W(\tau)}^{W_0} \sqrt{\frac{2}{mW}} \sigma n \tau_w dW.$$

Here  $m$  is the mass of a deuteron,  $n$  the electron density,  $n_D$  the deuteron density and  $\tau$  the confinement time of the fast deuteron. Due to the strong decrease of cross-section  $\sigma$  with energy the neutron production occurs mainly during the first half of the energy relaxation time, i.e. approximately 15 ms in ASDEX. For our plasma data the energy relaxation parameter  $n\tau_w$  depends only on the electron temperature. Hence the neutron production by injected deuterons depends essentially on the electron temperature of the target plasma. On the other hand the cross-section  $\sigma$  is a function of the ion temperature. For temperatures in the region of some keV it may be approximated by

$$\sigma(W, T_i) = \sigma(W, 0) + 1.16 \times 10^{-28} \sqrt{W} \cdot T_i \text{ barn (} W \text{ and } T_i \text{ in keV).}$$

The neutron rate is calculated by integrating the product of  $Y_D$  times the deposition profile for fast deuterons over the plasma volume.

Due to the lack of a direct measurement we assume that the deuteron density is a constant fraction of the electron density, the total neutron rate may then be expressed by

$$Q(T_i) = (n_D/n) Q_{inj}(0) \cdot (1 + \gamma T_i) + (n_D/n)^2 Q_{therm}(T_i)$$

where  $Q_{therm}$  is the thermonuclear rate of the bulk plasma. Here both,  $Q_{inj}$  and  $Q_{therm}$  are calculated with the electron density profile. Using this relation one can get the bulk ion temperature from the total neutron rate.

Neutron rate, experimental results: The neutron rate measurements have been carried out with BF<sub>3</sub>- and U-238 counters which have been calibrated independently by a Pu-238-B neutron source /1/ and by the nuclear emulsion technique. Figure 2 shows the measured neutron rate for the discharges to which the energy spectra of Fig.1 belong. The discharges are of the H-type /2/ with  $I_p = 420$  kA,  $\bar{n}_e = 4.9 \cdot 10^{13}$  cm<sup>-3</sup> and  $P_{inj} = 1.85$  MW at 45 kV. In Fig. 2  $Q(T_i)$  is calculated using the measured profiles of the electron temperature (ECE) and electron density (HCN interferometer),  $Z_{eff}$  as derived

from the loop voltage (for determination of  $n_D/n$ ) and the deuteron deposition profile as computed by the FREYA code /3/.  $T_i$  was determined by the neutron spectra. The shadowed areas indicate the measurement errors. For  $Q(T_i)$  they are mainly caused by some uncertainties in  $n_D/n$ . Within the limits of the errors all values agree well, thus indicating classical relaxation for the fast deuterons as assumed in the model.

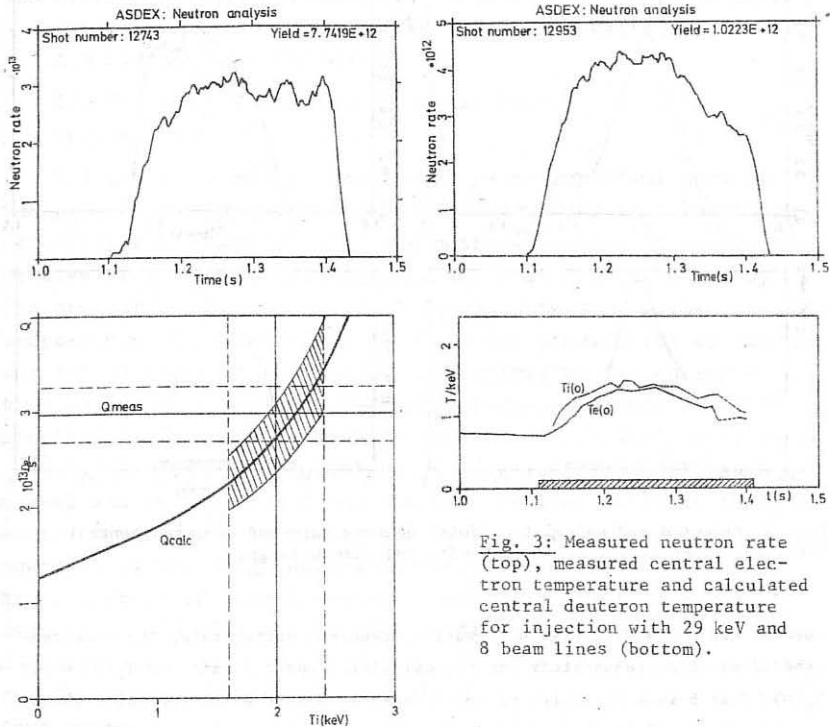


Fig. 3: Measured neutron rate (top), measured central electron temperature and calculated central deuteron temperature for injection with 29 keV and 8 beam lines (bottom).

Fig. 2: Measured neutron rate  $Q(t)$  (top), and calculated  $Q(T_i)$  for  $t = 1.25$  s for injection with  $W_0 = 45$  keV and 4 beam lines (bottom).

Next, results are presented for three other discharge types on ASDEX:

- A)  $P_I=1.3$  MW, 8 beam lines,  $W_0=29$  keV,  $I_p=420$  kA,  $\bar{n}=6.3 \times 10^{13} \text{cm}^{-3}$ , L-type
- B)  $P_I=3.9$  MW, 8 beam lines,  $W_0=45$  keV,  $I_p=275$  kA,  $\bar{n}=4.4 \times 10^{13} \text{cm}^{-3}$ , H-type
- C)  $P_I=2.9$  MW, 6 beam lines,  $W_0=45$  keV,  $I_p=420$  kA,  $\bar{n}=4.9 \times 10^{13} \text{cm}^{-3}$ , H-type

Figure 3 shows the measured neutron rate and central electron temperature

for case A. Further,  $T_i(0)$  is shown as calculated from the neutron rate. It is again about 10 % higher than the electron temperature, demonstrating the validity of the model also for low injection energies and L-type discharges.

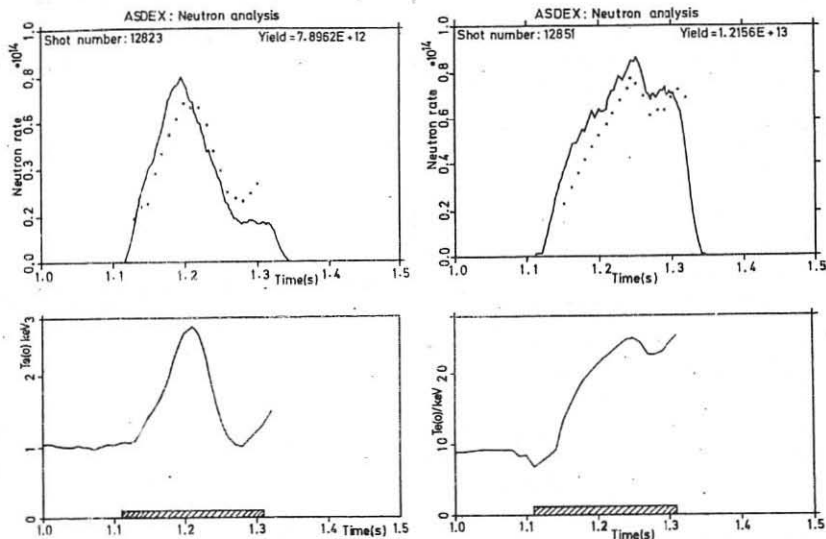


Fig. 4: Measured and calculated (dots) neutron rate and measured central electron temperature for case B (left) and case C (right).

For the cases B and C, Fig. 4 shows the measured neutron rate, the measured central electron temperature and the calculated neutron rate for  $T_i(0) = 1.2 T_e(0)$ . Case B is a discharge of the burst-free H-type which typically shows a strong decrease of neutron production after about 100 ms injection. Our results for  $Q_{calc}$  demonstrate the dependence of neutron production on target electron temperature. In case C the electron temperature, and accordingly the neutron rate, remain nearly constant during injection.

Part of this work was supported by the Deutsche Forschungsgemeinschaft.

#### References

- /1/ G. Assi, H. Rapp, IPP-Report III/70, March 1981
- /2/ F. Wagner, et al., Phys. Rev. Letters 49 (19), November 1982
- /3/ G.G. Lister, D.E. Post, R. Goldston, 3rd Symp. on Plasma Heating in Toroidal Devices, Varenna 1976.