

The Next Generation Global Gravitational Wave Observatory

The Science Book

arXiv:2111.06990v1 [gr-qc] 12 Nov 2021

GWIC
April 2021



*Out of the cosmic rife, I just picked me a star
another came along, from not so far
Thought it would be a real good bet
The best is yet to come*

*The best is yet to come and may be, it'll be fine
You think you've seen the sun
But you ain't seen two rattle and shine*

*A wait till the 3rd-gen's underway
Wait till our feisty stars have met
And wait till you see that everyday
You ain't seen nothing yet.*

— Sanjay Reddy with apologies to Frank Sinatra

SCIENCE BOOK SUBCOMMITTEE

Vicky Kalogera, Northwestern University, USA (Co-chair)
B.S. Sathyaprakash, Penn State USA and Cardiff University, UK (Co-chair)
Matthew Bailes, Swinburne, Australia
Marie-Anne Bizouard, CNRS & Observatoire de la Cote d'Azur, France
Alessandra Buonanno, AEI, Potsdam, Germany and University of Maryland, USA
Adam Burrows, Princeton, USA
Monica Colpi, INFN, Italy
Matt Evans, MIT, USA
Stephen Fairhurst, Cardiff University, UK
Stefan Hild, Maastricht University, Netherlands
Mansi M. Kasliwal, Caltech, USA
Luis Lehner, Perimeter Institute, Canada
Ilya Mandel, University of Birmingham, UK
Vuk Mandic, University of Minnesota, USA
Samaya Nissanke, University of Amsterdam, Netherlands
Maria Alessandra Papa, AEI, Hannover, Germany
Sanjay Reddy, University of Washington, USA
Stephan Rosswog, Oskar Klein Centre, Sweden
Chris Van Den Broeck, Nikhef, Netherlands

STEERING COMMITTEE

Michele Punturo, INFN Perugia, Italy (Co-chair)
David Reitze, Caltech, USA (Co-chair)
Peter Couvares, Caltech, USA
Stavros Katsanevas, European Gravitational Observatory
Takaaki Kajita, University of Tokyo, Japan
Vicky Kalogera, Northwestern University, USA
Harald Lueck, AEI, Hannover, Germany
David McClelland, Australian National University, Australia
Sheila Rowan, University of Glasgow, UK
Gary Sanders, Caltech, USA
B.S. Sathyaprakash, Penn State University, USA and Cardiff University, UK
David Shoemaker, MIT, USA (Secretary)
Jo van den Brand, Nikhef, Netherlands

GRAVITATIONAL WAVE INTERNATIONAL COMMITTEE

This document was produced by the GWIC 3G Committee, the GWIC 3G Science Case Team and the International 3G Science Team Consortium

Final release, April 2021

Additional Authors

P. Ajith, International Centre for Theoretical Sciences, India
Shreya Anand, Caltech, USA
Igor Andreoni, Caltech, USA
K.G. Arun, Chennai Mathematical Institute, India
Enrico Barausse, SISSA, Italy
Masha Baryakhtar, New York University, USA and University of Washington, USA
Enis Belgacem, Utrecht University, The Netherlands
Christopher P.L. Berry, Northwestern University, USA
Daniele Bertacca, Università degli Studi di Padova, Italy and INFN Sezione di Padova, Italy
Richard Brito, Sapienza University of Rome, Italy
Chiara Caprini, Université Paris-Diderot, France
Katerina Chatziioannou, Caltech, USA
Michael Coughlin, University of Minnesota, USA
Giulia Cusin, University of Geneva, Switzerland
Tim Dietrich, University of Potsdam, Germany
Yves Dirian, University of Geneva and University of Zuerich, Switzerland
William E. East, Perimeter Institute, Canada
Xilong Fan, Wuhan University and Hubei University of Education, China
Daniel Figueroa, École Polytechnique Fédérale de Lausanne, Switzerland
Stefano Foffa, University of Geneva, Switzerland
Archisman Ghosh, Ghent University, Belgium
Evan Hall, MIT, USA
Jan Harms, Gran Sasso Science Institute, Italy
Ian Harry, Max Planck Institute for Gravitational Physics, Potsdam, Germany and University of Portsmouth, UK
Tanja Hinderer, Utrecht University, The Netherlands
Thomas Janka, Max Planck Institute for Astrophysics, Germany
Stephen Justham, University of the Chinese Academy of Sciences, China and University of Amsterdam, The Netherlands
Dan Kasen, University of California Berkeley, USA
Kei Kotake, Fukuoka University, Japan
Geoffrey Lovelace, California State University Fullerton, USA
Michele Maggiore, University of Geneva, Switzerland
Alberto Mangiagli, Astroparticule et Cosmologie, France
Michela Mapelli, University of Padova and INFN Padova, Italy
Andrea Maselli, Gran Sasso Science Institute, Italy
Andrew Matas, Max Planck Institute for Gravitational Physics, Potsdam, Germany
Jess McIver, Caltech, USA and the University of British Columbia, Canada
Bronson Messer, ORNL and University of Tennessee, USA
Tony Mezzacappa, ORNL and University of Tennessee, USA
Cameron Mills, Cardiff University, UK
Bernhard Mueller, Monash University, Australia
Ewald Müller, Max Planck Institute for Astrophysics, Germany
Michael Pürrer, Max Planck Institute for Gravitational Physics, Potsdam, Germany
Paolo Pani, Sapienza University of Rome, Italy
Geraint Pratten, University of Balearic Islands, Spain and University of Birmingham, UK
Tania Regimbau, Université Grenoble Alpes, France

Mairi Sakellariadou, Kings College London, United Kingdom
Raffaella Schneider, INFN Roma, Italy
Alberto Sesana, Università di Milano-Bicocca, Italy
Lijing Shao, Kavli Institute for Astronomy and Astrophysics, Peking University, China
P. Thomas Sotiriou, The University of Nottingham, UK
Nicola Tamanini, Laboratoire des 2 infinis - Toulouse, France
Thomas Tauris, Aarhus University, Greece
Eric Thrane, Monash University, Australia
Rosa Valiante, INAF Roma, Italy
Maarten van de Meent, Max Planck Institute for Gravitational Physics, Potsdam, Germany
Vijay Varma, Caltech, USA
Justin Vines, Max Planck Institute for Gravitational Physics, Potsdam, Germany
Salvatore Vitale, MIT, USA
Huan Yang, Perimeter Institute and University of Guelph, Canada
Nicolas Yunes, University of Illinois Urbana-Champaign, USA
Miguel Zumalacarregui, Max Planck Institute for Gravitational Physics, Potsdam, Germany and University of California Berkeley, USA

Contents

Executive Summary	viii
1 Introduction	1
1.1 Prologue	1
1.2 Extreme Matter, Extreme Environments.	2
1.3 Observing Stellar-mass Black Holes Throughout the Universe.	3
1.4 Cosmology and Early History of the Universe.	4
1.5 Extreme Gravity and Fundamental Physics.	4
1.6 Sources at the Frontier of Observations	5
1.7 Summary	5
2 Extreme Matter, Extreme Environments	6
2.1 Nature of Matter at Highest Densities	7
2.2 Nucleosynthesis in Neutron Star Binary Mergers	9
2.3 Formation, Demographics and Merger Sites of Compact Binary Mergers	10
2.4 Jet Physics in Neutron Star Binary Mergers	10
2.5 Outlook for Extreme Matter and Extreme Environments	11
3 Observing Stellar-mass Black Holes Throughout the Universe	12
3.1 A survey of BHs throughout cosmic time	13
3.2 Expanding the BH mass spectrum	14
3.3 High-precision measurements of binary properties	15
3.4 Multiband gravitational-wave observations	16
3.5 Outlook for black hole gravitational-wave astronomy	17

4	Cosmology and the Early History of the Universe	18
4.1	Standard Siren Cosmology	19
4.2	Early History of the Universe	19
4.3	Astrophysical Binary Foregrounds and Large Scale Structure	22
4.4	Outlook for gravitational-wave cosmology and the early Universe	22
5	Extreme Gravity and Fundamental Physics	23
5.1	Nature of Gravity.	24
5.2	Nature of Compact Objects.	26
5.3	Nature of Dark Matter.	27
5.4	Outlook on Exploring Extreme Gravity and Fundamental Physics	28
6	Sources at the Frontier of Observations	29
6.1	Core-Collapse Supernovae	30
6.2	Sources of Continuous GWs	32
6.3	GW Bursts Associated with NS Flares and Glitches	32
6.4	Outlook for GW sources at the frontier of observations	33
	Acronyms and Abbreviations	34
	Bibliography	35

Executive Summary

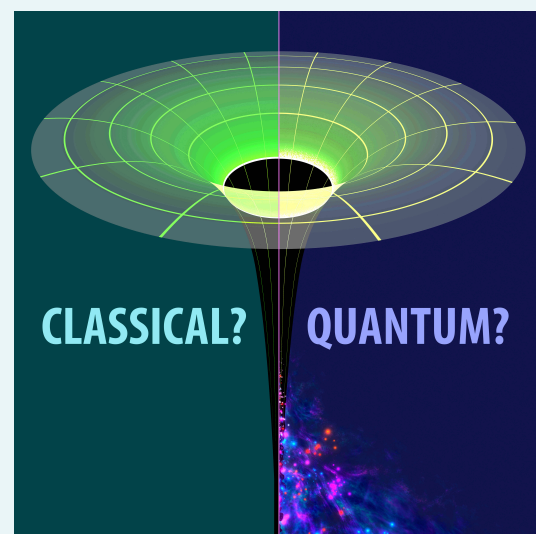
Within the last century, we have discovered incredible new truths about the Universe using telescopes and particle detectors. We now know it is expanding and composed mostly of dark matter and dark energy that we do not yet understand and neutrinos and ultra high-energy particles, whose origin remains a puzzle. We have learned that it harbors astounding entities such as black holes, regions of spacetime so strongly warped that nothing, even light, that falls inside can ever escape. We have learned that it is gravity that shapes the structure of the Universe, from enormous cosmic fibers of matter to galaxies and solar systems, to planets, stars, and black holes. Gravity drives the most extreme phenomena in the Universe, including incredibly violent collisions of black holes that in nearly an instant release millions of times the energy that our Sun will emit in its entire lifetime.

Gravity encodes information about distant cosmic phenomena in a messenger entirely different from light and particles: *gravitational waves*, tiny ripples in the fabric of spacetime emitted by accelerating mass. Unlike light and particles, gravitational waves interact weakly with matter, so they travel vast distances almost entirely unobscured by dust, the Milky Way, or the Earth itself. They offer a crystal clear signature of highly energetic phenomena otherwise hidden from us.

Gravitational waves from a binary black hole merger observed by Advanced LIGO in 2015 gave us our first glimpse of hugely energetic events that are undetectable with light. Detectors on Earth capture gravitational waves emitted by black holes orbiting so quickly that they approach the speed of light before colliding. The first detections also revealed a population of black holes with masses never before observed and charted the existence of black holes formed by collapsed stars farther from Earth than ever before.

BLACK HOLES: WHERE THE COSMOS MEETS THE QUANTUM REALM

A new generation of detectors will allow us to push Einstein's general theory of relativity to the limit and to test alternative theories that aim to resolve the fundamental contradictions between *quantum physics*, which describes the Universe at very small sizes, and *general relativity*, which describes the Universe at very large sizes. Although the two are incompatible with each other, no experiment yet has discovered new physics outside of either theory. Black holes provide a laboratory that smashes these theories together; squeezing enormous masses into infinitesimally small volumes. By sensing black hole collisions with high fidelity, new detectors would allow us to test completely new regimes of highly warped spacetime that could provide critical insight into this current paradox in physics. The precision of next generation detectors would also allow us to search for a differing gravitational-wave signature of new types of exotic compact objects unlike black holes as we currently understand them.



A network of next-generation gravitational-wave detectors will survey the extreme Universe with an unprecedented reach into deep space, beyond the cosmic dawn when the first stars began to shine.

They will sample stellar-mass black hole mergers across the visible Universe, observing these systems over the entire history of the cosmos, and chronicle the evolution of black holes from the earliest mergers and inform if they grow from collapsed stars to supermassive black holes at the center of galaxies, billions of times as massive as our Sun.

The 2017 observation of the merger of two neutron stars by LIGO and Virgo and its aftermath with the full spectrum of electromagnetic (EM) radiation was a spectacular first success of multimessenger astronomy with gravitational waves that reaffirmed theoretical models of the brightest EM events. Next-generation ground-based gravitational-wave detectors will have the unique capability to observe the violent mergers of neutron stars and their aftermath at far greater distances and with the much higher precision required to address outstanding problems in the physics of dense objects. Neutron stars are only the size of a city yet contain more mass than the sun. They are the densest material objects in the Universe, where intense gravity completely obliterates the atomic structure of matter familiar to us. Even quarks—the tiniest constituents that remain confined inside neutrons and protons in ordinary matter—are likely to be liberated in the neutron star core.

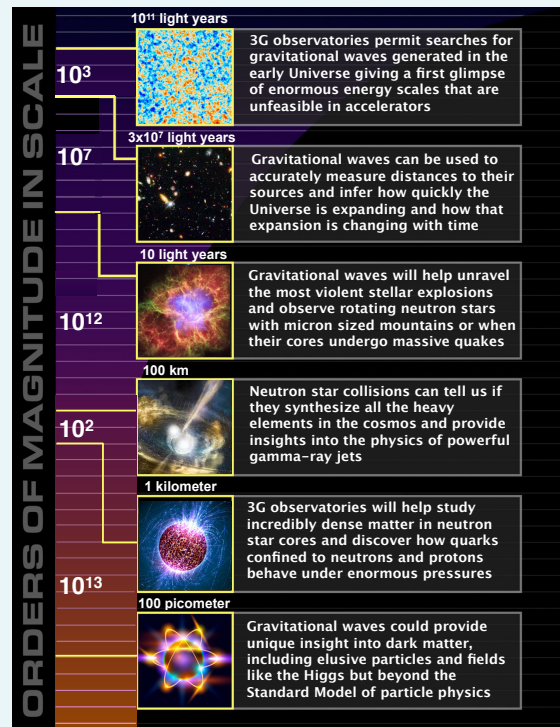
Next generation detectors will allow us to observe neutron star mergers and peek directly into their cores as they tear each other apart by tidal forces before smashing together. These observations are critical to infer new knowledge and understanding about nuclear physics and the states of matter containing quarks. With next-generation multimessenger astrophysics enabled by new detectors, we will learn how much of the Universe's gold and platinum was produced by neutron star collisions.

Furthermore, gravitational waves from binary black holes and neutron stars are *standard sirens*—the distance to the source is encoded in the observed gravitational waves. Thus, merging binaries provide a new precision tool for observational cosmology that will help us gain new insight into how the Universe is expanding and evolving and if dark energy is just a cosmological constant or if there is missing new physics associated with the late-time accelerated expansion of the Universe.

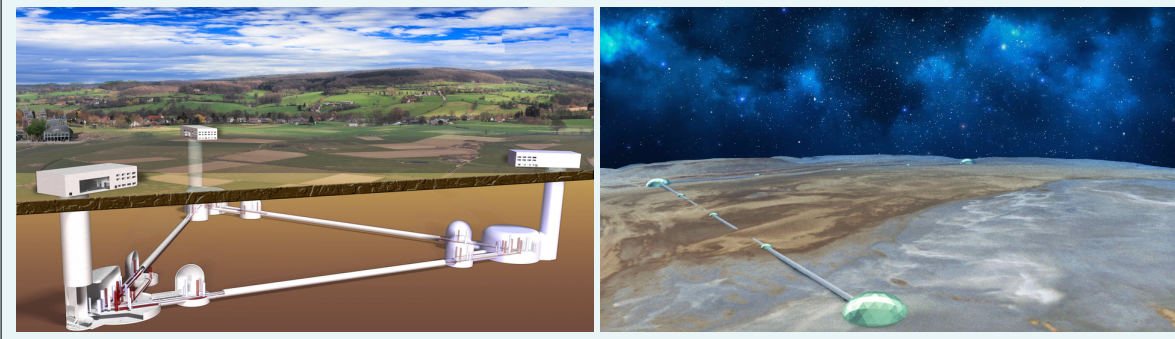
We will also achieve unprecedented insight into cosmic explosion mysteries. Multimessenger astronomy with next generation detectors will allow us to better investigate why core-collapse supernovae explode to

OBSERVING THE UNIVERSE ON ALL SCALES

The diagram shows physical phenomena on different scales explored by gravitational-wave observations. Starting from the scale of the Universe, almost 100 billion light years, represented by the cosmic microwave background at the top, the diagram progressively shows scales that are smaller than the previous ones by the factor shown on the left. On scales about 1000 times smaller are giant galaxy clusters 30 million light years across. Another factor 10 million smaller is the size of a supernova remnant of 10 light years. A trillion times smaller still is the merger environment of binary neutron stars about 100 km across. The core of a neutron star, about 1 km, is a trillion times smaller and contains matter at densities similar to that of atomic nuclei. On scales ten thousand smaller still gravitational waves could probe the nature of dark matter.



ARTISTS CONCEPTION OF EINSTEIN TELESCOPE (LEFT) AND COSMIC EXPLORER (RIGHT)



Artists conception of the Einstein Telescope (left panel) and Cosmic Explorer (right panel) observatories. ET is conceived to be six, V-shaped, underground interferometers, formed out of 10 km sides of an equilateral triangle, while Cosmic Explorer is conceived to be an L-shaped, overground interferometer, with 40 km arms.

seed the formation of new stars and whether starquakes cause mysterious bursts of radio emission. And as with any completely new method of observation, there is also the possibility that next generation detectors will reveal completely new dark phenomena, unseen with light, that we have not yet conceived of.

Today's gravitational-wave detectors are barely sensitive enough to detect the loudest gravitational waves in the Universe, like a simple radio able to pick up only the loudest signals. Next-generation network detector designs leverage cutting-edge technology to surpass current ground-based detectors, making their ability to measure passing gravitational waves more than ten times better than the current instruments.

More powerful detectors will let us listen to the gravitational-wave universe with unprecedented fidelity, fully revealing the rich physics encoded in the waves but currently hidden by observational uncertainty. Einstein Telescope (ET) is a European design featuring six V-shaped interferometers in a triangular topology with 10 km interferometer arms and Cosmic Explorer (CE) is a U.S. design for one or two interferometers with 40 km L-shaped interferometer arms. ET and CE are expected to detect hundreds of thousands of mergers, as well as tens of thousands of multimessenger sources that would also likely emit EM radiation and particles that telescopes and neutrino and cosmic ray detectors can observe. A network of three detectors distributed around the globe will triangulate the gravitational wave signal's location in the sky, critical information that will guide telescopes on Earth and in space in searches for related EM emission.

21st century astronomy will be further revolutionized by the launch of the space-based LISA gravitational-wave observatory, expected in 2034. LISA will sense gravitational waves emitted by more massive systems than ground-based detectors, detecting the signature of orbiting black hole systems up to years before ground-based detectors observe them collide. Combining space-based and ground-based observations will allow us to catalog a much broader expanse of the extreme gravitational Universe than ever before.

Gravitational waves have already given us a first glimpse of the dark, hidden, violent Universe. A global next-generation gravitational wave observatory will propel the field of astrophysics and all foundational science research forward. Observing light, neutrinos and cosmic rays in concert with next-generation gravitational wave detectors will launch enormous advances beyond the current limits of human knowledge, from the quantum realm to the largest cosmological structures in the known Universe.



1. Introduction

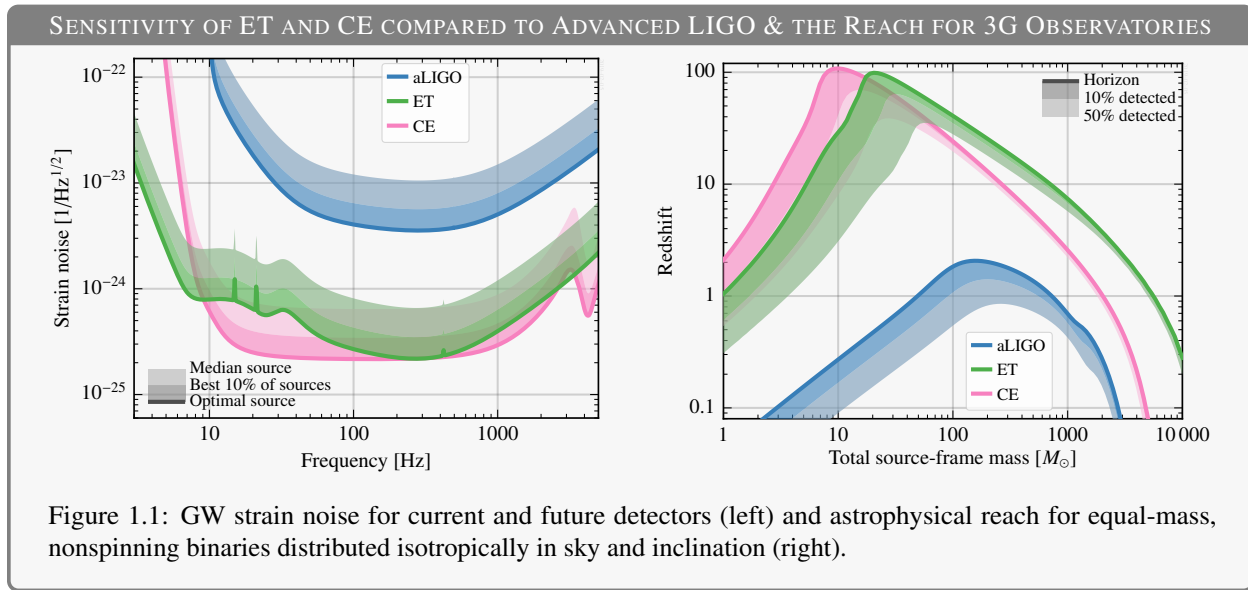
1.1 Prologue

On 14 Sep 2015, the Laser Interferometer Gravitational-Wave Observatory (LIGO) transformed the way we explore the universe. With LIGO, we were able to sense, for the first time, the gravitational-wave (GW) sky. The signal came from the merger of two black holes at a distance of 1.3 billion light years. This was also when binary black holes were discovered and, for the first time, their merger was observed. The masses of the companions, 29 and 36 solar masses, were unexpectedly heavy and the merger converted some 3 solar masses into energy in a mere 200 ms. Since that first discovery, many more black hole mergers have been found by LIGO and the European Virgo, some with masses as small as 2.6 solar masses, heaviest neutron star or lightest black hole, and others as large as 85 solar masses, so large that they could not have formed from the evolution of massive stars. In most cases the component black holes seem to be non-spinning, contrary to what X-ray observations indicate.

Two years later, on 17 Aug 2017, a new era in multimessenger astronomy began with the observation by LIGO and Virgo of the merger of two neutron stars, followed by the detection, 1.7 s later, of a gamma-ray burst from the same source by the Fermi gamma-ray space telescope and INTEGRAL. These observations could together localize the event well enough that its host galaxy was quickly found by optical telescopes. The merger produced spectacular fireworks that were captured by telescopes across the entire electromagnetic (EM) spectrum from radio to infrared and optical to X-rays. This treasure trove of data gave us answers to decades old puzzles in fundamental physics and astronomy: verified that GWs travel essentially at the speed of light, confirmed that binary neutron star mergers are progenitors of short gamma ray bursts and prolific sites for the formation of heavy elements, measured the Hubble constant in a completely new way using GWs for the source's distance and EM observations for its redshift and constrained neutron-star radii to be between 9.5 and 13 km by measuring the tidal deformation of neutron stars.

The LIGO and Virgo detectors are yet to achieve their design sensitivities. They will be augmented with new facilities, the Japanese KAGRA and LIGO-India. Yet, based on the modest glimpses of the sources discovered to date, we know that the full exploration of the GW sky will require a new generation of detectors of a size that demands new facilities. With the aid of such detectors we will be able to observe sources at the edge of the Universe, unveil the properties of matter at the highest densities in the cosmos, provide a new precision tool for observational cosmology and explore the nature of dynamical spacetimes. A detector network with a leap in sensitivity will resolve signals with far greater precision and fidelity that will pave the way for serendipitous discoveries, observing novel phenomena and unearthing new physics.

Beyond Advanced Detectors: During 2008-2011 the design study of a third generation (3G) GW observatory in Europe, Einstein Telescope (ET), developed the concept of a triangular interferometer, 10 km on a side, housing six V-shaped interferometers whose combined sensitivity is a factor ~ 20 better than Advanced Virgo and pushing the low-frequency sensitivity down from 10 Hz to 3 Hz. A similar effort is currently underway in the US to study the science case for and technical design of a 40 km arm length interferometer called Cosmic Explorer (CE), with sensitivity similar to ET (see Figure 1.1, left plot). For the current study we assume that *the 3G network* consists of one ET in Europe and one CE each in the US and Australia. A network of



at least three sites is required to accurately localize sources in the sky and infer their distances. ET alone could measure the wave’s polarization but cannot resolve all the parameter degeneracies to determine the sky position even when the signals last for days.

The science potential of the 3G network is immediately apparent from the dramatic improvement in strain sensitivity that CE and ET are able to deliver (Figure 1.1, right panel). The network makes a leap of 1–2 orders of magnitude in the redshift reach for binary coalescences compared to Advanced LIGO and Virgo. The network will survey a large redshift range for merging binary black holes and provide a massive catalog of detections to constrain their population and origins. The network will explore a wide parameter space of quantum chromodynamics and study high density matter in a region complementary to heavy ion physics experiments. The Box below summarizes the science potential of a 3G observatory, elucidated in the next several paragraphs.

SCIENCE TARGETS FOR THE NEXT GENERATION OF GRAVITATIONAL WAVE DETECTORS

GW astronomy provides a complementary window to EM, neutrino and particle astronomy that could reveal hitherto unseen world. A new generation of detectors will:

- *determine* the properties of dense matter, *discover* phase transitions, and the emergence of quarks
- *reveal* merging black holes across the cosmos and *search* for seeds of supermassive black holes
- *investigate* the particle physics of the primeval Universe and *probe* its dark sectors
- *explore* new physics in gravity and in the fundamental properties of compact objects
- *understand* physical processes that underlie the most powerful astrophysical phenomena

1.2 Extreme Matter, Extreme Environments.

Neutron stars are the densest objects in the cosmos and sites of stupendously strong magnetic fields, up to billions of tesla. Six decades after their discovery, we still lack a clear understanding of the equation of state of their deep cores and the origin of their strong magnetic fields. Neutron stars in binaries are subject to the tidal fields of their companions although the tides raised are extremely small. The extent of tidal deformation depends on the internal structure of neutron stars and the net effect is to accelerate the rate of inspiral allowing to read-off their internal structure from the observed phase evolution of the signal. The merger remnant could be a rapidly rotating, short-lived, hypermassive neutron star that eventually collapses to a black hole. GWs from the merger will lead to tight measurements of NS radii and hence reveal the equation of state of both

cold and hot, supranuclear matter and the deconfinement phase transition of quarks and gluons.

The origin of heavy elements in the Universe has been a long-standing problem. EM observation of GW170817 provided irrefutable evidence that binary neutron star mergers are prolific sites for the production of lanthanides and other heavy elements. The 3G network will facilitate EM follow-up of thousands of mergers, a number that is required to confirm if solar and stellar abundance of heavy elements can be explained by mergers alone or if other production channels, such as supernovae, are necessary.

GW170817 resolved that binary neutron star mergers are progenitors of short gamma-ray bursts. Nevertheless key questions about central engines that produce gamma rays still remain. For example, we do not have a clear picture of the jet properties nor how those properties depend on the progenitor characteristics. EM follow-up facilitated by the 3G network will allow a better understanding of the physics of gamma-ray jets, the opening angle of the jet and its distribution; GWs could tell us the nature of the merger remnant and if the central engine is a transient hypermassive neutron star or a promptly collapsed black hole.

3G observatories will detect binary neutron star mergers from epochs far before the peak of star formation activity. Millions of mergers are expected to be detected by the 3G network. The properties of the detected sources and the environments in which they occur will provide key data to test astrophysical models of the formation and evolution of double neutron star and black hole-neutron star binaries, while also informing the history of star formation activity up to redshifts of 5–8.

1.3 Observing Stellar-mass Black Holes Throughout the Universe.

The 3G network will have nearly all-sky sensitivity, detecting stellar-mass black hole binaries of ~ 10 – 100 solar mass from epochs before the first stars formed at redshifts $z \sim 30$ (Figures 1.1 and 1.2). Consequently, the 3G network could reveal a population of primordial black holes in this mass range formed by quantum processes in the early Universe, in addition to compiling a census of black holes over a range of masses throughout the cosmos.

The merger rate of binary black holes observed so far imply that the 3G network will detect hundreds of thousands of mergers each year. This large population will help us study the merger rate as a function of redshift up to the beginning of the epoch of reionization. It will also help us explore how these rates are correlated with metallicity and galaxy evolution.

Until 2015, it was widely believed that irrespective of how massive a progenitor star was the black hole that resulted from it would be lighter than about 20 solar masses. LIGO and Virgo have detected many black holes with masses in excess of that number. Indeed, GW190521 revealed a black hole of 85 solar masses that, theory says, cannot form from massive stars. Understanding how such black holes form and if they grow through repetitive mergers is an outstanding question in astrophysics. There are a number of competing models. Black holes that form in isolation in globular clusters could sink into the dense cluster cores where they dynamically interact with other holes to form coalescing binaries. Binaries of massive stars formed in active star formation sites could directly evolve into binary black holes that merge within the Hubble time. If primordial black holes significantly contribute to dark matter then they could occasionally form merging binaries in galactic halos. The 3G network will pin down the masses, spins and demographics of black holes, determine principal formation channels and resolve fundamental questions about their origin.

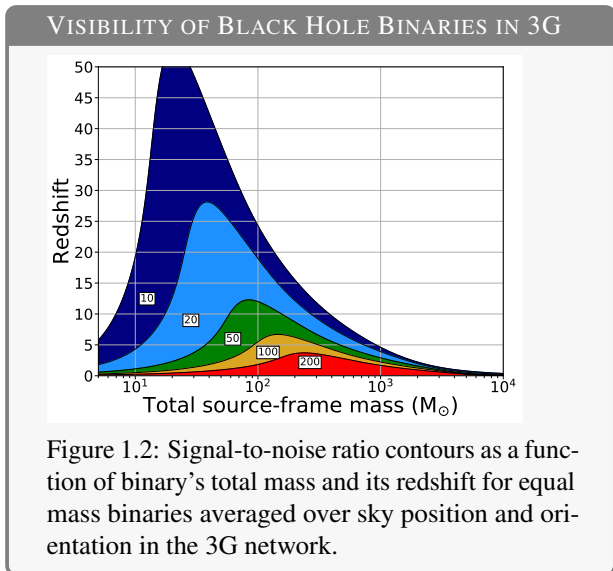


Figure 1.2: Signal-to-noise ratio contours as a function of binary's total mass and its redshift for equal mass binaries averaged over sky position and orientation in the 3G network.

1.4 Cosmology and Early History of the Universe.

GWs from the inspiral and merger of compact binaries can be used to infer the luminosity distance to their sources without the need to calibrate them with standard candles. This is because the orbital dynamics of binary black holes and neutron stars is largely determined by Einstein's theory of gravity. A handful of parameters, e.g. masses and spins of the companion stars, precisely control the pattern of the emitted GWs. The amplitude of that pattern is fixed by the distance to the source, sky position and orientation of the source relative to a detector, which can be inferred with a network of three or more non-collocated detectors. This contrasts with the dynamics of other astrophysical systems, such as supernovae, that require detailed modelling of their composition and environment, making it extremely hard to predict the emitted GW signal with any precision. Consequently, with a population of compact binary mergers observed with 3G detectors, and their redshifts obtained by follow-up EM observations, it will be possible to accurately measure cosmological parameters such as the Hubble parameter, dark matter and dark energy densities and the equation-of-state of dark energy, giving a completely independent and complementary measurement of the dynamics of the Universe.

The cosmological population of point sources create a stochastic GW background. Indeed, with advanced interferometers we could detect the background created by binary black holes and neutron stars throughout the Universe and can do so by cross correlating data from two or more detectors. Such backgrounds would reveal the history of the formation and evolution of these sources and the underlying stellar population. On the contrary, 3G detectors will identify most compact binary mergers in the Universe, giving us a treasure trove of data to study the large-scale distribution of galaxies and their clusters.

Stochastic GWs could also be produced in the early Universe. As the Universe cools from its primeval hot and dense state it undergoes several phase transitions that are expected to generate GW backgrounds. Detection of such backgrounds would dramatically transform our state of knowledge of the underlying particle theory at energy scales that will never be accessible to terrestrial accelerators. Defects, such as cosmic strings, associated with symmetry breaking phase transitions, could also produce stochastic and deterministic signals. The landscape of primordial sources, while uncertain, is a high-risk, high-reward endeavor to pursue in the era of 3G detectors.

1.5 Extreme Gravity and Fundamental Physics.

GWs emanate from regions of strong gravity and large curvature, carrying uncorrupted information from their sources. Imprint in the signal is the nature of the gravitational field, characteristics of the sources and the physical environment in which they reside. Their observation in 3G detectors can put general relativity to the most stringent tests, help explore violations of the theory in strong fields such as the dynamics of black hole horizons, and discover properties of dark matter.

The 3G network offers numerous opportunities to discover failure of general relativity, e.g. in the form of new particles and fields that violate the strong equivalence principle. It is also possible to detect Lorentz invariance violations or variation in Newton's constant, both imprint in the propagation of GWs. One might also see the signature of quantum gravity in the form of parity violation seen in the nature of the GW polarization or in the birefringence of the waves propagating over great distances. Ultra-light Bosonic fields proposed in certain extensions of the Standard Model could be detected via their effect on the orbital dynamics of black hole binaries and spin properties of black hole populations observed by the 3G network.

Black holes are the most compelling explanation for the companions in binary coalescences discovered by LIGO and Virgo detectors. The tell-tale signature of a black hole would be present in the quasi-normal mode spectrum of the merger remnant, whose frequencies and damping times should depend only on the remnant's mass and spin. Signature of additional degrees of freedom would be seen as inconsistency in the remnant's parameters determined by the different modes. Certain alternatives to black holes could mimic the quasi-normal mode spectra, but they could emit additional signals in the form of echoes of the ingoing

radiation reflected from their surface, which could be observable in the 3G network.

Big Bang cosmology is largely consistent with general relativity but the accelerated expansion of the Universe in its recent history cannot be explained by the theory, indicating either its failure or the presence of exotic form of matter-energy density, of which we know very little. Observations on galactic to cosmological scale provide unequivocal indirect evidence for the presence of weakly interacting dark matter, but none has been directly detected in spite of concerted efforts over the past six decades. The 3G network might detect various forms of dark matter including axionic and other dark matter fields around black holes and neutron stars, primordial black holes, etc.

1.6 Sources at the Frontier of Observations

The physics of supernova explosion, glitches in the frequency of pulsars and quakes in highly magnetized neutron stars (or magnetars) are open problems in astrophysics. Many of these systems will generate GWs that could be observed with 3G detectors at distances of several million light years for supernovae and within the galaxy for pulsar glitches and magnetar flares. GW observations of these systems with the 3G network, enhanced by EM and neutrino observatories, will allow us to probe extreme astrophysics and address key questions that have hindered progress in our understanding of the mechanism behind stellar explosions.

From the observed spectrum of GWs it will be possible to determine the physics of core collapse supernova: the different phases of the collapse, the nature of the explosion that dominates the production of GWs and the asymmetry of the collapse and what triggers that asymmetry. Information about the rotation rate of the progenitor star is also encoded in the observed signal and it should be possible to understand how the initial state of the progenitor star determines the final state of the collapse, a black hole or a neutron star. Such observations will be greatly aided by all-sky optical and infrared surveys of the stellar population in nearby galaxies, as well as cosmic rays and neutrinos.

Isolated neutron stars could emit GWs if they are not spherical and don't rotate about their symmetry axis. Indeed, they are persistent sources with the emission lasting for millions of years. Advanced detectors are not likely to detect continuous waves from known pulsars, although all-sky blind searches may reveal hitherto unexpected sources. The 3G network could observe neutron stars whose polar and equatorial radii differ by no more than 10 to 100 microns. This will provide invaluable information about their crustal strengths and the equation-of-state of high density nucleons in their outer cores.

Accreting neutron stars, e.g., in low-mass X-ray binaries, could acquire quadrupole deformations from the in-falling matter that could lead to a perpetual source of GWs. Indeed, the 3G network will help resolve if the accretion torque balanced by the GW back-reaction torque is responsible for the observed limiting spin frequencies of neutron stars in low-mass X-ray binaries. The 3G network will also probe the role of magnetic fields in transient radio emission from magnetars provided the mechanism is caused by crustal quakes that result in the emission of GWs. This could further constrain the equation-of-state of neutron star crusts.

1.7 Summary

LIGO and Virgo discoveries have ushered in a new era in multimessenger physics and astronomy. GW observations can be used to probe the nature of ultra dense matter, reveal quantum chromodynamic phase transitions, study the formation and evolution of stellar mass black holes from the epoch of the formation of first stars, measure cosmological parameters, examine phase transitions in the early Universe, test general relativity in dynamical spacetimes, discover the nature of dark matter and other exotic compact objects, and explore the physics of the most violent processes in the cosmos. The mind boggling reach of the 3G network is difficult to fathom but guarantees serendipitous discoveries, with the potential to unearth new physics. Indeed, 3G observatories will operate in a survey mode wherein signals that do not fit our expectations will be flagged off for further study. The science case for building a new generation of GW detectors that can probe deep into the cosmos and observe a variety of different processes is immensely rich and massively rewarding.

2. Extreme Matter, Extreme Environments

SCIENCE TARGET

Determine the properties of dense matter, discover phase transitions, and the emergence of quarks.

The discovery of the binary neutron star (BNS) merger GW170817 [1] was a watershed moment in astronomy and astrophysics. Multimessenger observations of the source observations provided incontrovertible evidence that BNS mergers are connected to short-hard gamma-ray bursts (GRBs) [2, 3] and its optical localization [4] unveiled that they are prolific sites of heavy element nucleosynthesis [5–7]. It confirmed that GWs and light travel essentially at the same speed, and allowed the first measurement of the Hubble constant using GW standard sirens [8, 9] ushering in a new era in cosmology. Furthermore, the event enabled the first measurement of the neutron star (NS) tidal deformability [2, 10] and provided the most robust and stringent constraints on the NS radius and the equation of state of dense matter under extreme conditions inaccessible to experiments and first-principles theoretical calculations [10–15]. Circumstantial evidence from the EM counterparts for the formation of a black hole (BH) on a timescale of tens of milliseconds provided tighter, albeit model-dependent, constraints on the maximum mass of NSs, and lower bounds on their radii [16–18].

However, a number of important questions were left unanswered (see Sec. 2.1). Accurate observations of a diverse population BNS mergers with detectors of greater sensitivity and bandwidth will be key to shedding light on the nature of extreme matter in extreme environments produced by the mergers.

KEY SCIENCE GOALS

Multimessenger observations of numerous luminous events in the Universe involving dense matter in extreme environments will uncover several key puzzles in fundamental physics:

- **Nature of matter at supranuclear densities.** What are the fundamental properties of the densest matter in the cosmos? How do quarks and gluons manifest in the cores of the most massive neutron stars?
- **Production sites of heavy elements.** What elements are produced in NS mergers and how? Are they able to explain the abundances of elements heavier than iron in the solar system and in stars?
- **Formation and evolution of compact binaries.** How do NS binaries form and evolve? What are their demographics, merger rates, and mass and spin distributions through cosmic time?
- **Central engines of short-hard GRBs.** What is the role of the merger remnant and the physics of central engines powering panchromatic EM counterparts to NS mergers? How do they relate to short GRBs?

Capabilities of Next Generation Detector Networks: The 3G network will compile a survey of the Universe of a large sample of BNS and NS-BH mergers. Table 2.1 shows the detection capability of 3G observatories compared to the network of advanced detectors at their design sensitivity. In computing the event rates the local co-moving merger rate was taken to be $1000 \text{ Gpc}^{-3} \text{ yr}^{-1}$ and redshifts at the epoch of merger were sampled assuming the Madau-Dickinson star formation rate, with an exponential time delay (an e-fold time of 100 Myr) between formation and merger [19]. Tens of thousands of well-localized events

detected by the 3G network will provide ample opportunities for EM follow-up of these mergers as opposed to a handful of them by the current network.

In addition to discovering a great number and diversity of mergers through cosmic time, the wide band sensitivity of 3G detectors will enable tracking the full inspiral, merger and post-merger GW signals. These unique capabilities will help address key science questions on the properties of dense matter. The 3G network will accurately measure the masses and spins NSs and determine their long-sought equation of state, probe the merger dynamics, state of the merger remnant and BH formation, and explore, for the first time, properties of matter at even greater density and temperature if the remnant does not promptly collapse to a BH.

The sub-arcsecond localization of the panchromatic EM counterpart will provide information about environment and geometry of the event, its host galaxy, the physical state and evolution of the ejected material and the nucleosynthesis of heavy elements. Thus, the combination of information derived independently from the GW and EM signals will be immensely powerful to build a complete, self-consistent astrophysical picture, enable cosmological applications, and constrain formation scenarios.

2.1 Nature of Matter at Highest Densities

Neutron stars are precious laboratories for the subatomic physics of matter under unique conditions. The multitude of phenomena connected with multimessenger emissions from BNS mergers is of broad interest to nuclear and particle astrophysics. Our current understanding of the NS interior is captured in Fig. 2.1. The theoretical understanding of matter up to densities present in terrestrial nuclei ($\rho_0 \simeq 2.5 \times 10^{14} \text{ g cm}^{-3}$) are fairly advanced and relevant to the NS crust. However, nuclei dissolve at higher densities $\rho \gtrsim \rho_0/2$ into a uniform liquid of neutrons, with a small admixture of other particles including protons, electrons, and muons [20]. In NSs with large masses, densities in the core may be sufficiently high for exotic states of matter to appear. Furthermore, at densities $\gtrsim 2\rho_0-3\rho_0$ the distance between nucleons becomes comparable to their size, and their quark sub-structure is expected to manifest and phase transitions to new states of matter containing de-confined quarks may occur [21].

The nature and location of the transition from hadronic to quark matter remains unknown, but is a fundamental question with broad implications. The properties of matter inside NSs directly affect their global characteristics, masses and radii, motivating significant observational [22, 23] and theoretical [24–28] efforts to constrain the properties of NS matter and measure NS masses and radii. Radio observations of pulsars have yielded accurate mass measurements of a handful of NSs [23]. The discovery of a massive NSs with $M \simeq 2 M_\odot$ [29–31] has had far-reaching implications for the equation of state of dense matter [32]. However, accurate measurements of the NS radius from X-ray observations have been more challenging

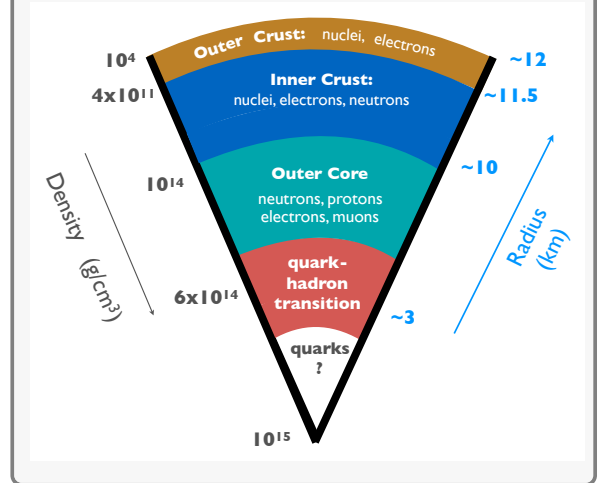
BNS EVENT RATES IN 2G & 3G NETWORKS

Table 2.1: Expected number N of BNS detections per year, the number of events localized to within 1, 10 and 100 deg^2 (N_1 , N_{10} and N_{100} , respectively) and the median localization error M in square degrees, in a network consisting of LIGO and Virgo (HLV), LIGO, Virgo, KAGRA and LIGO-India (HLVKI) and the 3G network.

Network	N	N_1	N_{10}	N_{100}	M
HLV	48	0	16	48	19
HLVKI	48	0	48	48	7
3G	990k	14k	410k	970k	12

INTERNAL STRUCTURE OF A NS

Figure 2.1: Composition of matter in the interior of a NS predicted by theory. Quark degrees of freedom become important at the densities encountered in the inner core. The nature of the transition to matter containing de-confined quarks is unknown.



since they rely on poorly tested models of EM emission from or near the NS surface. Efforts to model and interpret X-ray data from accreting NSs during bursts, and in quiescence, suggest that NS radii are in the range 9–13 km [23, 33], albeit with untested model assumptions. For a few pulsars, NASA’s NICER mission is anticipated to provide reliable radius measurements, using a different method [22]. Results from the first NICER observations are promising, although the errors associated with the extracted NS radius remains large [34, 35].

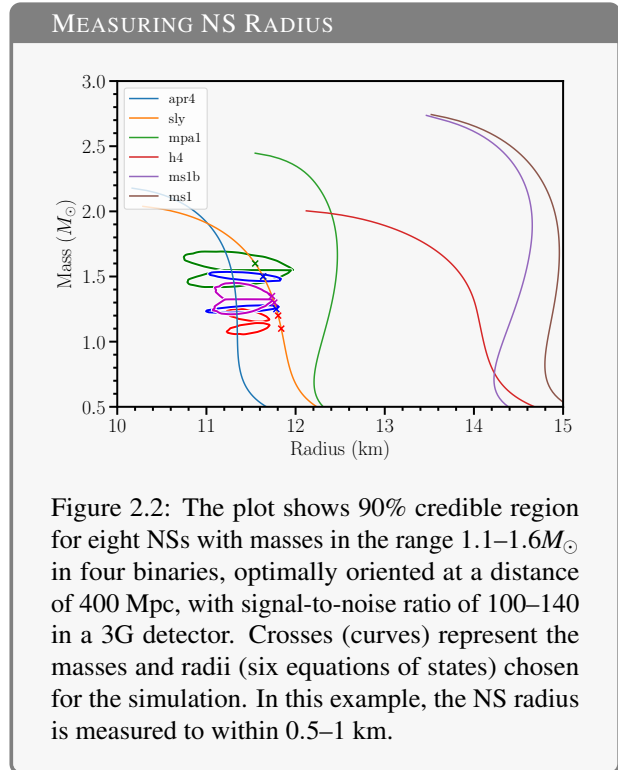
The 3G will probe the properties of dense matter in a *diverse population* of NSs by measuring a variety of matter-dependent GW signatures during the inspiral phase of BNS systems. These arise from tidal effects, including the tidal excitation of a NS’s internal oscillation modes [36–39] that can provide direct information about phase transitions, rotational deformations [40], spin-tidal couplings [41, 42], and the tidal disruption of the NS by a BH companion [43, 44]. The 3G network will measure the radii of several NSs over a wide mass range including both light and heavier NSs to within 0.5 to 1 km, and discern subdominant matter effects on GWs. Both these aspects are critical for measuring the properties of dense matter and discovering phase transitions. Fig. 2.2 shows a projection of the precision with which 3G detectors will allow us to measure the NS mass and radius.

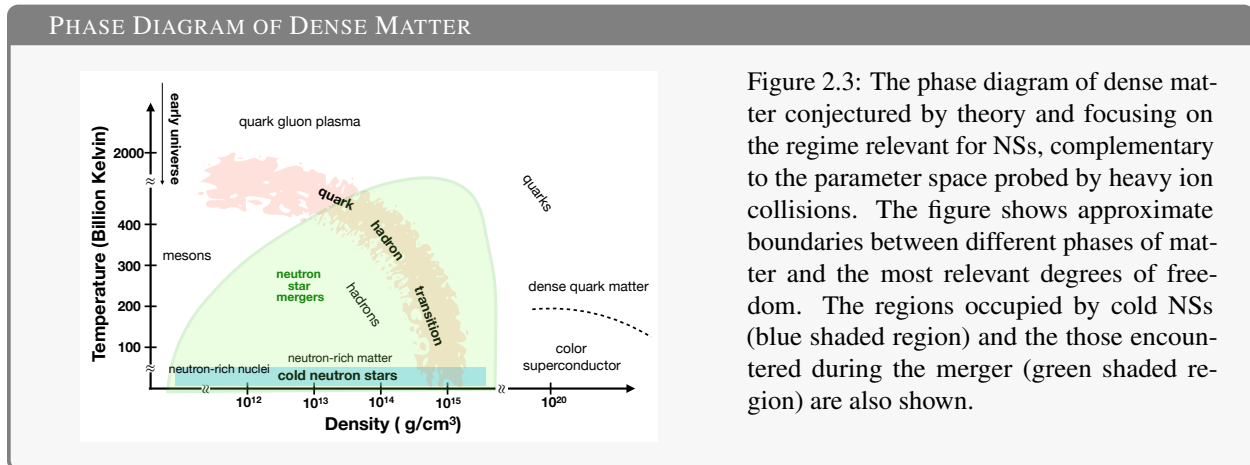
The 3G network will further open an exceptional window onto fundamental properties of *matter in a completely unexplored regime*, at higher temperatures and yet greater densities than encountered in NSs, which is accessible only during the merger and post-merger epochs in BNS collisions.

The merger outcome depends on the companion masses and the equation of state of NS matter [45–48]. Above a critical total mass of the binary the merger will result in prompt collapse to a BH, while for very low-mass progenitors a stable remnant may form. For a wide range of parameters, the merger outcome is a short-lived hypermassive NS that involves complex microphysics and generates a significant amount of GWs [49–58] ultimately collapsing to a BH. The rich GW signals from the merger and post-merger regimes have frequencies in the range 1–5 kHz and are thus difficult to measure with advanced detectors; detailed studies of the complex physics driving the dynamics of NS binary mergers and beyond are a major unique capability of the 3G detector network.

Various phases of strongly interacting matter that will be uniquely accessible with 3G detectors are depicted in Fig. 2.3, which focuses on the regimes relevant to NS binaries that are complementary to those explored by heavy ion collisions. Matter encountered during BNS mergers, shown as light-green shaded region, explores a large swath of the phase diagram of dense matter. The 3G network will enable unprecedented measurements of the new physics encountered during the coalescence and post-merger epochs, with EM and neutrino counterparts providing complementary information to obtain a deeper and more complete understanding of extreme states of matter.

In summary, observations by the 3G network will shed light on many critical questions about the nature of NSs and the fundamental subatomic physics of matter: Are NSs composed solely of similar constituents as nuclei on earth or do they contain condensates of exotic particles or quark matter phases? Do NS mergers





produce novel phases of matter not realized inside nuclei and heavy-ion collisions? What is the nature of the transition from nuclear to quark matter? How do nuclear reactions and neutrinos shape NS merger dynamics?

2.2 Nucleosynthesis in Neutron Star Binary Mergers

A long-standing puzzle in astrophysics is how the elements heavier than iron came into being. About half of these elements are believed to have been created by a nuclear process of rapid neutron capture (the r-process) but it is unclear which astrophysical sites are the main contributors. GW170817 and its associated thermal EM counterpart provided the first direct identification of a NS merger as a prolific site of r-process nucleosynthesis [59]. However, determining the degree to which NS mergers contribute to cosmic chemical abundance and evolution will require a more extensive sample of the rates, locations, timescales, and nucleosynthetic yields of the various types of merger events.

Heavy elements can be synthesized in BNS or NS-BH tidal disruptions when clouds of neutron-rich material are expelled, either dynamically during the merger or through winds blown off the remnant accretion disk. The subsequent radioactive decay of the freshly synthesized elements powers a thermal ultraviolet/optical/infrared EM transient called a kilonova. The brightness and color of the kilonovae are diagnostic of both the total mass of r-process elements and the relative abundance of lighter to heavier elements [60].

Whereas historical studies of chemical evolution have relied on observing fossil traces of r-process elements mixed into old stars, multimessenger observations provide the unique opportunity to study heavy element formation at its production site and to determine how the initial conditions of an astrophysical system map to the final nucleosynthetic outcome. Answering the basic question of the extent to which BNS and NS-BH mergers are the site of r-process production will require multimessenger observations of a large sample of events. GW measurements with the 3G network will pin down the rate of mergers and the binary properties, such as the binary type (BNS or NS-BH), companion masses, spin-orbit alignment, and the merged remnant lifetime, while optical/infrared photometry of the associated kilonovae will determine the average r-process yields and probe the relative abundance distribution of heavy elements. These observations will also illuminate the key physics driving the r-process and kilonova, such as the equation of state of dense matter, the fundamental interactions of neutrinos and the magneto-hydrodynamics of accretion.

Statistical studies of multimessenger observations will reveal how r-process production depends on host galaxy type, location and redshift, allowing us to piece together the history of when and where the heavy elements were formed over cosmic time. Such studies can determine the distribution of delay times between star formation and mergers, thereby addressing whether some of these mergers occurred promptly enough to explain the enrichment of the oldest metal poor stars and the extent to which compact binaries receive strong kicks that may expel them from their host galaxies, a factor that is important for understanding whether mergers can explain the unusually high r-process enhancement seen in some dwarf galaxies [61].

In the era of 3G detectors, optical kilonovae will be detectable by the Vera Rubin Observatory out to 3 Gpc and infrared characterization photometrically by WFIRST/Euclid and spectroscopically by JWST/GMT/TMT/E-ELT would be out to 1 Gpc ($z < 0.2$). The multimessenger information about these event will enable answering questions such as: How nuclear reactions and neutrinos shape NS merger dynamics and nucleosynthesis? How do the properties of nuclei that are far from stability impact the EM emission from material ejected during NS mergers? How do the progenitor properties impact the nucleosynthesis and kilonovae?

2.3 Formation, Demographics and Merger Sites of Compact Binary Mergers

Observations of NS binary merger systems are essential to advance our understanding of how nature assembles these systems—either via standard (isolated) binary star evolution or via dynamical encounters in dense stellar environments. Precise measurements of NS masses and spins in these systems across cosmic time will provide key evidence for their origin in different types of supernova explosions and fossil records of close binary progenitor star interactions and accretion history after birth. Any signs of NS masses being different in NS-BH mergers compared to BNS mergers will yield crucial information on their formation process and the evolution of massive stars.

A key question about compact binary mergers is their demographics, as this could reveal their formation mechanism. Localization of merger events to less than galactic scales (~ 30 kpc) is essential to unambiguously infer associations of mergers with their host galaxies. Without an EM counterpart the vast majority of GW events will have error boxes that greatly exceed the typical radii of potential host galaxies. The census of the binaries, their locations, and environments will provide deep insights into the formation and evolution of NS binaries and their connection with the progenitor stars [62–64].

EM follow-up of NS binary mergers will be critical in pinning down host galaxies. For binaries involving a NS and a BH with a mass ratio that is not too large, depending on the BH's spin, the NS may get tidally disrupted. The debris may result in accretion disk around the BH and lead to EM counterparts that might rival the absolute visual magnitude of the GW170817 kilonova, and, unless they occur in globular cluster cores, will be detectable out to $z = 0.5$ in the reddest filters.

Based on our current understanding, galaxies are assembled by the merger of smaller proto-galaxies and star formation peaks near $z \sim 2$ [65]. Identification of kilonovae beyond $z \sim 0.5$ requires hour-long integrations on 8m class facilities such as LSST or Subaru, rendering the identification of the host galaxies of binary NS mergers near the peak of star formation more challenging in the absence of a gamma-ray burst jet pointing towards the Earth, even with ELTs. Nevertheless, at redshifts $z < 0.5$ 3G detectors will work in concert with astronomy facilities to enable thousands of host galaxy identifications from NS binary mergers where the kilonova counterpart is observable. At larger distances, the identification will be possible only through the detection of an associated gamma-ray burst afterglow, which can be much more luminous than a kilonova if the jet is directed towards the Earth.

2.4 Jet Physics in Neutron Star Binary Mergers

Relativistic explosions and compact-object mergers can generate collimated, energetic jets of fast-moving material and radiation. Prior to GW170817, our understanding of jet physics came from studies of gamma-ray bursts, active galactic nuclei and X-ray binaries. Multimessenger observations provide an entirely new perspective on this topic. For instance, the panchromatic study of GW170817 revealed that there was both a narrow ultra-relativistic jet [9, 66–68] and a wide-angle mildly relativistic cocoon from surrounding material ejected during the merger [7, 69, 70]. This event opened up many questions for future observations to answer. Specifically, what is the connection to the class of cosmological short hard gamma-ray bursts? Does a wide-angle mildly relativistic cocoon always accompany a binary NS merger? Does the jet always successfully escape the cocoon or is it sometimes choked? How do the observed jet properties vary as a function of viewing angle, mass ratio, hypermassive NS lifetime, remnant spin, and ejecta mass? Do mergers

FACILITIES FOR OBSERVING EM COUNTERPARTS TO GWs

Table 2.2: Present (*P*) and future (*F*) EM facilities that are able to observe faint/distant counterparts to GWs. Detection Limit (**DL**, 1 hr exposure time) for UV, optical, and near-IR facilities are expressed in AB magnitudes, for X-rays in 10^{-16} erg s $^{-1}$ cm 2 , and for radio in μ Jy. Distance reach (**D** in Mpc) of facilities for GW170817-like events are also shown.

	Facility	DL	D		Facility	DL	D
Gamma-rays	<i>Fermi P</i>	S/N 5	80	Optical	Keck/VLT	23	500
	AMEGO <i>F</i>	S/N 5	130		Gemini Obs.	23	500
X-rays	<i>Swift P</i>	S/N 5	~80	GMT <i>F</i>	25	1265	
	<i>Chandra P</i>	30	150	Spec. TMT <i>F</i>	25.5	1592	
	ATHENA <i>F</i>	3	480	E-ELT <i>F</i>	26	2005	
	<i>Lynx F</i>	6	450	Keck/VLT	21.5	481	
	STROBE-X <i>F</i>	S/N 5	120	Infrared GMT <i>F</i>	23.5	762	
UV	<i>HST</i> (im) <i>P</i>	26	2000	Spec. TMT <i>F</i>	24	960	
	<i>HST</i> (spec) <i>P</i>	23	400	E-ELT <i>F</i>	24.5	1208	
Optical	Subaru <i>P</i>	27	3200	VLA (S) <i>P</i>	5	91	
Imaging	LSST <i>F</i>	27	3200	Radio ATCA (CX) <i>P</i>	42	51	
IR	<i>WFIRST F</i>	27.5	4800	ngVLA (S) <i>F</i>	1.5	353	
	<i>Euclid F</i>	25.2	1700	SKA-mid (L) <i>F</i>	0.72	634	

produce prompt EM signals or even precursors? What is the distribution of the time delays between the EM and GW signal arrival times? What are the characteristics of a jet from a NS-BH merger? A census of NS binary mergers, and full GW and EM coverage of the signals, joint multimessenger parameter inference will be key in understanding the physical origin of jets, ubiquitous around relativistic sources. For the first time, a direct measurement of the BH spin in a source emitting a collimated jet, will enable to establish the close correlations between the jet power, the spin and the inflow rate from the debris disk, which determines the conditions for launching the jet. The sensitivity of gamma-ray, X-ray and radio telescopes will enable studying jet physics out to 500 Mpc, thus requiring a sample of the order of a thousand events localized to better than few square degrees to map the full parameter space provided by the 3G network.

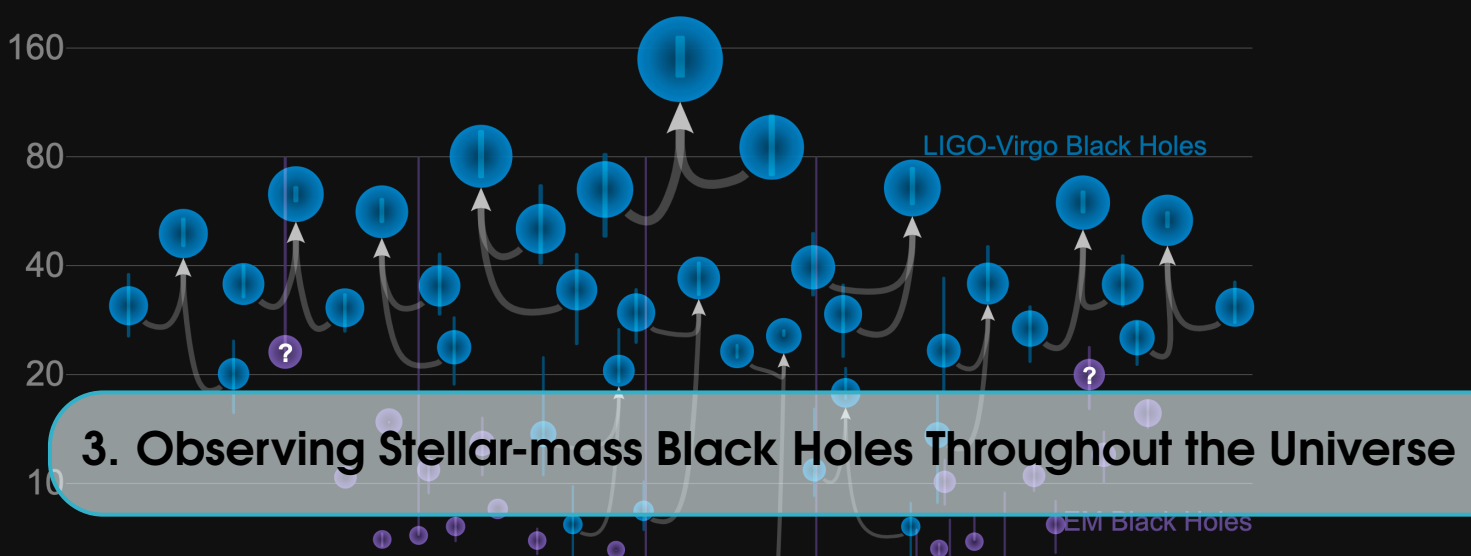
2.5 Outlook for Extreme Matter and Extreme Environments

Observations of BNS and NS-BH mergers with a network of 3G detectors will transform our understanding of the fundamental properties of matter in unexplored regimes of density and temperature and, in conjunction with EM facilities, will address longstanding questions about the formation of heavy elements in the universe, the central engines of highly energetic EM transients, and the formation and evolution of NS binary systems.

SCIENCE REQUIREMENTS

The unique capabilities of 3G detectors required to accomplish these science goals are:

- *an order-of-magnitude greater sensitivity* than 2G detectors enabling observations of the NS binary population in the cosmos and measurements of loud-source properties with exquisite accuracy,
- *a wider frequency range* than 2G detectors that will allow tracking the entire GW signals from the inspiral through the merger, tidal disruption and beyond, and
- synergies with panchromatic EM facilities, that will be critical to fully capitalize on the rich multimessenger science potential of these sources for cosmology, fundamental physics, and astrophysics.



SCIENCE TARGET

Reveal merging black holes across the cosmos and search for seeds of supermassive black holes.

Merging binary black holes are sources unique to GW astronomy—they are the most frequently observed sources to date. We now know that binary black holes (BBHs) form ubiquitously in galaxies and so far appear to be completely dark. These mergers are unrivalled laboratories for testing extreme gravity, and exquisite astronomical sources for gaining insight into the origin and evolution of massive stars in the Universe. With a leap in sensitivity and increased frequency bandwidth, the 3G network will observe them back to the early Universe, chart how the population evolves with time, discover a broader range of masses, and connect stellar-mass black holes (BHs) with the supermassive BHs (SMBHs) found in the centres of galaxies, obtain precision measurements of BH properties, and finally resolve the mysteries of their formation. GW astronomy is perfectly suited to studying BHs, and with 3G detectors we would achieve a complete picture of the family of stellar-mass BBHs.

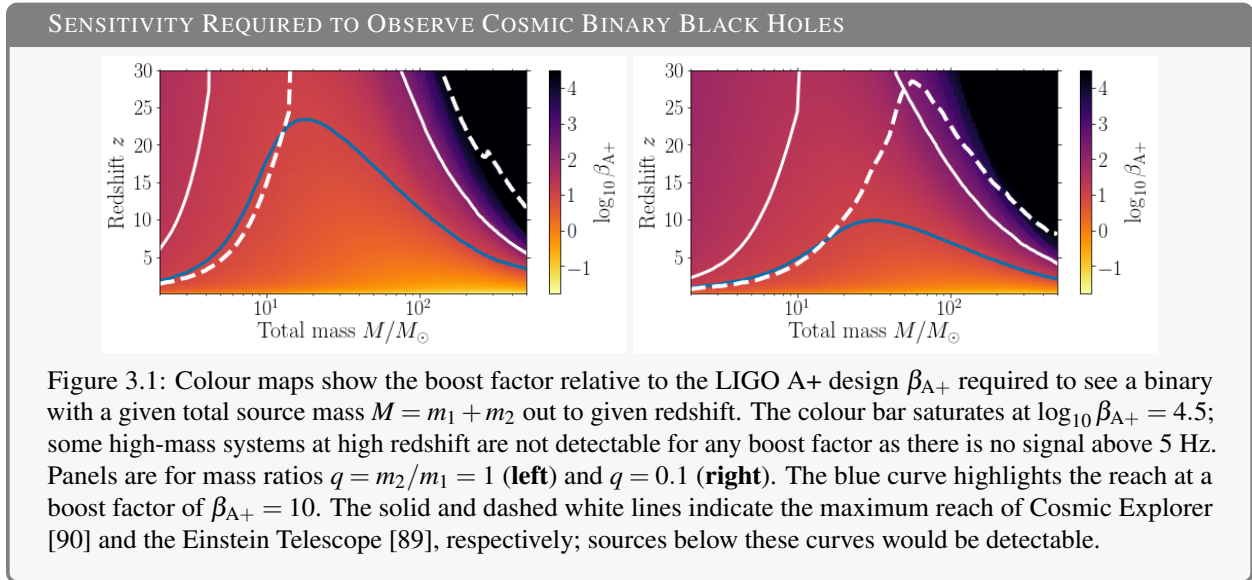
KEY SCIENCE GOALS

The 3G network will uncover BBHs throughout the cosmos, back to the beginning of star formation and detect new sources, if they exist, beyond stellar-mass binaries, such as intermediate-mass BBHs.

- **Discover BBHs throughout the Universe.** What is the merger rate as a function of cosmic time and how does it relate to the star formation rate, metallicity and galaxy formation and evolution?
- **Reveal the fundamental properties of BHs.** What are the mass and spin demographics of BHs throughout the Universe? Are they correlated and do they evolve with redshift? What do they reveal about the formation and evolutionary origin of BHs?
- **Uncover the seeds of SMBHs.** GW observations have proven that intermediate-mass BHs can form at least from BBH mergers. The 3G network promises to explore their population, reveal if intermediate-mass BH mergers occur in nature and serve as the long-sought seeds of SMBHs?

The first three observing runs of Advanced LIGO and Advanced Virgo have yielded the discovery of *more than 50* BBH systems. Already these detections have revolutionized the astrophysics of stellar-mass BHs [71–75] and provided first new tests of general relativity [76–80]. Through the end of the 2020s, the current advanced detector network will continue to be enhanced as sensitivities reach design goals and a new detector in Japan, KAGRA, comes online [81, 82]. In the BBH domain, we will be able to detect a pair of $10M_{\odot}$ BHs out to a cosmological redshift of $z \simeq 1$ when the Universe was 6 Gyr old [82]. The annual BBH detection rates are forecast to be several hundreds of mergers and science benefits will compound through accumulated observing time and growing detected samples [83–88].

Beyond this horizon, step-wise sensitivity improvements with the next generation of ground-based GW observatories will be required if we are to pursue major science questions that cannot be answered by the



current and near-term GW facilities [e.g., 89, 90]. Current-generation GW detectors are able to provide constraints on the merger-rate densities in the local Universe and broad constraints on component masses [72, 91]. However, precise measurements of, for example, spin magnitudes and tilts are of paramount importance to understand the origin and the evolutionary physics of binary systems [83–85, 92–96]. This information is essential to obtain insights on the formation channels of compact binaries. So far we have been surprised by the properties of individual exceptional sources, but the population constraints are too weak to distinguish among formation path possibilities. We highlight here how 3G GW ground-based detectors will enable us to survey deeper, to observe a wider range of frequencies, and to make more precise physical measurements; how observations can be synergistically combined between 3G and space-based GW observatories, and how these results will be transformational in the study of BBH astrophysics.

3.1 A survey of BHs throughout cosmic time

With a 3G detector network, for the first time, we will detect BBH mergers at redshifts beyond $z \sim 1$ and we will measure the evolution of the BBH merger rate out to redshifts of $z \gtrsim 10$ when the Universe was < 500 Myr old [19, 82, 88]. GW astronomy would thereby gain a synoptic view of the evolution of BHs across cosmic time, beyond the peak of the star-formation rate, which took place at $z \sim 2$ when the Universe was 3 Gyr old [65], back to the cosmic dawn around $z \sim 20$ when the Universe was only 200 Myr old and the first stars were forming in pristine dark matter halos.

Measurements of the merger rate as a function of redshift combined with high fidelity measurements of the BH physical parameters will enable conclusive constraints on the BBH formation channels. Stellar-origin BH formation tracks cosmic star formation [97–102], while the density of primordial BHs is expected to be independent of the star formation density [103, 104]; different binary formation channels are predicted to lead to different distributions of delay times between formation and merger [86, 105–114]. Therefore, determining the merger rate as a function of redshift provides a unique insight into the lives of binary BHs. *Only next-generation GW detectors can survey the complete redshift range of merging BBHs and provide a sufficiently large catalog of detections to constrain the full BBH population and their origins.*

To capture BBH mergers across the stellar mass spectrum (up to total masses of $M = m_1 + m_2 \simeq 200M_\odot$) all the way back to the end of the cosmological dark ages ($z \simeq 20$), a major advance in GW detector sensitivity is required. This cannot be delivered by the maximal sensitivity planned for the current ground-based detector facilities. We quantify this sensitivity step by the boost factor β_{A+} relative to the LIGO A+ design [115] between 5 Hz and 5 kHz (and no sensitivity outside this range). In Figure 3.1, we show this boost factor,

required to detect an optimally-oriented, overhead binary at a signal-to-noise ratio (SNR) of 8, as a function of the binary’s total mass and redshift. The boost factors β_{A+} needed to acquire a complete census of BBH mergers throughout the Universe are well within the design aspirations the 3G network; for these specific sensitivity assumptions, BBH mergers of total mass $M \sim 10\text{--}40M_{\odot}$ can be detected out to $z \sim 10^2$.

Observations of the cosmological distribution of coalescing binaries would complement planned EM surveys designed to study stars and stellar remnants back to cosmic dawn [116–120], as well as millihertz GW observations made by the *Laser Interferometer Space Antenna (LISA)* [121], which can observe systems ranging from local stellar-mass binaries (days to years before they enter the frequency range of terrestrial detectors) [122, 123] to SMBH systems in the centres of galaxies [124, 125]. In X-rays SMBHs are detectable due to gas accretion in galactic nuclei. X-ray observatories like *Athena* [126] and the mission concept *Lynx* [127] would detect SMBHs back to high redshift ($z \gtrsim 7$); *Lynx* would observe 10^3M_{\odot} BHs to $z \sim 5$ and 10^2M_{\odot} BHs to $z \sim 2$, while *Athena* will survey these in the nearby Universe. Next-generation GW detectors have the unique potential to observe stellar-mass BH systems all the way back to the early Universe.

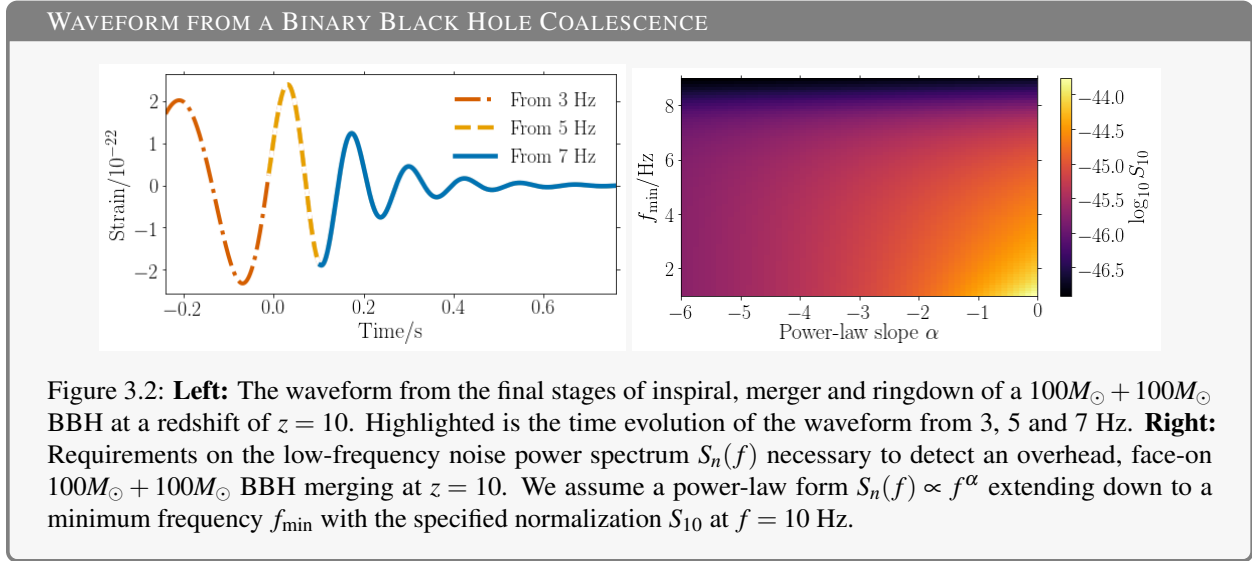
3.2 Expanding the BH mass spectrum

EM astronomy has benefited enormously from advancing observing facilities to cover an expanded range of frequencies. These enable new probes of previously known sources, and allow for the discovery of new types of previously unobserved sources. *3G GW detectors have the unique capability to push the frequency range down to $\simeq 1$ Hz and up to $\simeq 5$ kHz*, while improving performance across the band in between.

The merger frequency for a coalescing binary scales inversely with the mass of the binary, hence observing at lower frequency opens up the potential of detecting more massive BHs. The first intermediate-mass BH (mass in excess of $100M_{\odot}$) has already been observed as the result of a BBH merger [128], but there may be multiple formation paths for these BHs. Reaching down to frequencies of $\simeq 1$ Hz is the most robust means to chart the population of intermediate-mass BHs, and discover any mergers of intermediate-mass BHs—which may be the process through which SMBHs form [129–132]. SMBHs are observed up to redshift $z = 7.54$ [133] as quasars, at lower redshifts as active galactic nuclei [134], and today in massive galaxies in their quiescent state [135], and cover a mass range from $\sim 10^4M_{\odot}$ [136–139] up to $> 10^{10}M_{\odot}$ [140–142]. SMBHs may have light seeds ($\sim 10^2\text{--}10^3M_{\odot}$), formed from massive stars in low metallicity halos which evolve into BHs beyond the pair instability gap [143], or heavy seeds, formed from supermassive (proto)-stars of $\sim 10^4\text{--}10^6M_{\odot}$ growing through continued and fast accretion within their birth clouds, which eventually collapse down to BHs [144–149]. In particular, the observation of high-redshift BHs with mass $\gtrsim 100M_{\odot}$, beyond the (pulsational) pair-instability supernova mass gap (where supernova explosions are hypothesised to completely disrupt the star leaving behind no remnant) [150–155], would be key to understand not only the properties of very massive ($\gtrsim 250M_{\odot}$) metal-poor stars [156, 157], but also the assembly of the first massive BHs in the Universe [158–162].

In Figure 3.2, we illustrate the importance of sensitivity in the 1–10 Hz regime. Even with detectors sensitive to 3 Hz, we see only one cycle of a $100M_{\odot} + 100M_{\odot}$ circular binary with non-spinning components at $z = 10$ before merger. This system is not observable above 10 Hz. Therefore, the objective to observe the most massive BBHs of stellar origin and the potential seeds of SMBHs in the Universe’s early history requires new detectors sensitive to currently inaccessible frequencies below ~ 10 Hz.

The detectability of intermediate-mass BHs places requirements on low-frequency sensitivity. We can model the low-frequency noise power spectral density of the detector as a power-law $S_n(f) = S_{10}(f/10 \text{ Hz})^{\alpha}$ and assume that the power law extends to some minimal frequency f_{\min} , below which the detectors have no sensitivity. In Figure 3.2, we show the combination of power law α , minimum frequency f_{\min} and the normalisation S_{10} necessary to detect an optimally located and oriented merger of two $100M_{\odot}$ intermediate-mass BHs at $z = 10$. There is a trade-off between the power-law slope, minimal frequency and overall normalization, such that a range of specifications can fulfill the science requirements.

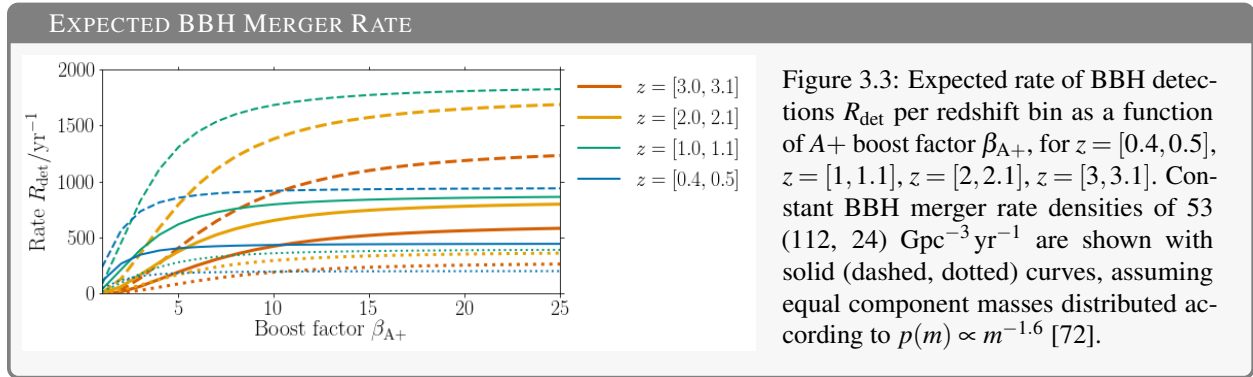


3.3 High-precision measurements of binary properties

A leap in sensitivity combined with the increased frequency bandwidth of the 3G detectors will enable high-precision measurements of the properties of individual binaries [163–165]. Parameter uncertainties are inversely proportional to the SNR [166]. The increase in SNR made possible by the greater sensitivity will lead to exquisite measurements of the loudest events. Increased bandwidth enables the coalescence to be tracked for a longer time, improving estimates of quantities like the spins. Masses, spins, merger redshifts, orbital eccentricities and (where possible) associations with host galaxies all give complementary insights into binary physics. High-precision measurements of individual systems allow us to make detailed studies of their origins and fundamental physics [167–172]. Combining many events together lets us study the properties of the population. *The unique and critical advantage of BBH observations with the 3G network is the combination of high-precision measurements for a very large number of detected sources*, something that cannot be delivered by the current detectors.

As an example, consider a highly precise reconstruction of the BH mass spectrum. At high masses, a gap is predicted to exist between $\sim 50M_{\odot}$ and $\sim 130M_{\odot}$ due to (pulsational) pair-instability supernovae [152, 154, 155, 173], although nature has some ways to form BHs in it [128]. At lower masses, there is potentially a gap between the maximum neutron star mass and the minimum stellar BH mass [75, 174–177]. Determining the precise bounds for these gaps would provide insight into the mechanics of supernova explosions [178–181], insights into the neutron star equation of state [10, 11, 16, 182–186], and even details of nuclear reaction rates [155, 187]. It can be shown that: (i) for the high-mass gap, if the desired accuracy on the mass gap boundary measurement is $\sigma_g \sim 1M_{\odot}$, with a conservative individual mass uncertainty for near-threshold detections of order $\sigma_m \sim 10M_{\odot}$, $N \gtrsim 500$ detections are required; (ii) for the low-mass gap, $\sigma_g \sim 0.3M_{\odot}$ and $\sigma_m \sim 3M_{\odot}$, would require $N \gtrsim 1500$ BBH detections. To provide robust answers to questions regarding massive star evolution and BBH formation, we need to trace the dependence of the boundaries of the mass gaps on metallicity and hence redshift. Therefore, it is desirable to observe ~ 1000 sources in each redshift bin of width $\Delta z = 0.1$, since we may expect knowledge of the star formation rate and metallicity distribution to be available at this resolution on the timescale of the 3G network [65]. Such observations would provide $\sim 3\%$ fractional accuracy on the merger rate per redshift bin, sufficient to determine the redshift evolution of the rate, and constrain details of the binary evolution at that redshift [86, 87].

With this in mind, we plot the number of expected BBH detections for a next-generation detector as a function of its boost factor relative to A+ in Figure 3.3. This assumes a BBH merger rate that does not evolve in redshift and is roughly consistent with current GW observations [72]. From this, the target of ~ 1000



detections per redshift bin is achievable with boost factors of $\beta_{A+} \sim 10$ after only 2 years of observing time. These factors are possible only with next-generation GW detectors.

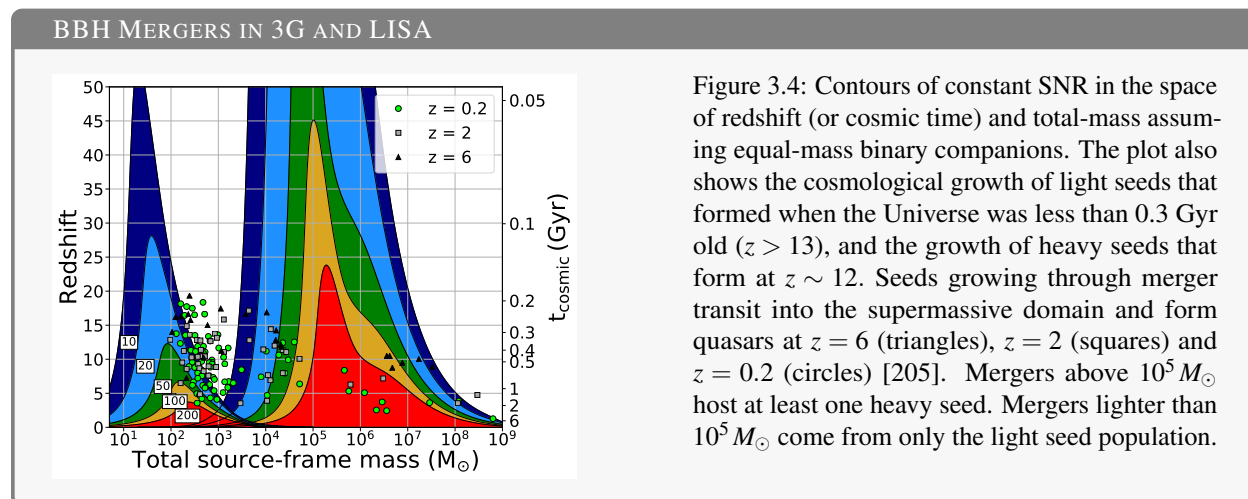
We note that observing across a broader range of frequencies gives a more complete picture of BBH properties. The precession of component spins misaligned with the orbital angular momentum occurs over many orbits [188, 189]. Its imprint is easier to discern over longer inspirals, and hence becomes more apparent with low-frequency data. Orbital eccentricity is rapidly damped through GW emission [190]. This means that it is near immeasurably small for current GW detectors [191–193], but reachable for 3G detectors. Both the spins and the orbital eccentricity are indicative of the formation channel; enabling their measurement for large samples will have a transformative effect on our ability to answer questions about BBH origins.

3.4 Multiband gravitational-wave observations

Joint observations of GW events by *LISA* at millihertz frequencies and 3G detectors at higher frequencies maximises their science potential. If *LISA* had been observing in 2010, it would have detected GW150914 years before it was observed by LIGO [122]. *LISA* will potentially see up to hundreds of stellar-mass BBH mergers of $M > 20\text{--}30M_{\odot}$, up to $z \approx 0.3$ [122, 194]. A small fraction of these will sweep across the detector band within few years that will eventually be detected by ground-based detectors. Multiband observations enables unrivalled measurements of BBH properties.

LISA would provide a precise measurement of the system’s eccentricity to a precision of $\Delta e < 0.001$ [195], sky localization to 0.1 deg^2 , and time to coalescence within few seconds, several weeks prior to coalescence [122]. This enables EM telescopes to be pointed in the right direction before the merger, permitting a much deeper coverage from radio to gamma-ray than what is possible without any early warning. Alternatively, one can use the information extracted by the 3G network to dig out sub-threshold *LISA* events [196]. From an astrophysical standpoint, the eccentricity information from *LISA* can be combined with the spin measurement from 3G detectors to better constrain different formation channels [197–199]. Multiband observations will also facilitate tests of general relativity [200–203] by enhancing the sensitivity to specific deviations arising in the long inspiral as predicted, for example, from dipole radiation not predicted in general relativity [204].

Figure 3.4 shows the cosmological growth of the earliest light and heavy seeds that transit into the supermassive domain, inferred using a semi-analytical model for the formation of quasars at $z = 6$, $z = 2$ and $z = 0.2$ [205]. Binaries of light seeds ($\sim 10^2 M_{\odot}$) are accessible to the 3G network with an SNR of 10–20 at $6 < z < 15$. They then enter the *LISA* domain with larger SNRs as they grow to a few $10^4 M_{\odot}$. Mergers above $10^5 M_{\odot}$ come from at least one heavy growing seed in the binary, lighter mergers arise from the light-seed population. Combining the observations in the two different frequency domains will provide the first ever census of coalescing BBHs forming in the Universe. The comparison between the detection rate between light and heavy seed BHs in the two GW bands will be instrumental in determining, at statistical level, the relative contribution of light and heavy seeds in building up the population of SMBHs, and the role of mergers versus accretion in determining their growth. Figure 3.4 highlights that intermediate-mass BHs are a prime



multiband GW astronomy target [206, 207]. As well as both 3G and space-based detectors being able to see the same intermediate-mass BH binary signal at different phases in the inspiral, two GW bands allow the intermediate-mass BH population to be explored in different regimes. While *LISA* will be sensitive to mergers of $M \geq 10^3 M_\odot$ binaries out a redshift of $z > 20$, 3G detectors will be able to access $M \sim 100 M_\odot$ populations at comparable redshifts. *Multiband GW observations will quantify the continuity between the stellar-mass, intermediate-mass, and SMBH populations.*

3.5 Outlook for black hole gravitational-wave astronomy

The 3G network will enable the measurement of the cosmological evolution of the mass and spin distributions of BBHs and probing their dependence on star formation history and metallicity evolution with redshift. They will make robust discovery of intermediate-mass black holes, if they exist, and reveal the boundaries of any mass gaps. The precise measurements of physical properties for large numbers of BH systems, back to the cosmic dawn, would lead to constraints on the physics of massive star evolution in single and binary systems, and to place bounds on the different formation channels of merging BBHs. Additionally, 3G observations could solve the long-standing mystery of the nature of SMBHs' seeds. 3G data would complement those from future EM and space-based GW observatories, enabling the maximum scientific return from these facilities. We have an unparalleled opportunity to advance the frontiers of stellar astrophysics, the fundamental physics of compact objects, and the formation mechanisms for the entire spectrum of BHs.

SCIENCE REQUIREMENTS

The 3G network will transform our BH studies by enabling cosmological probes of their formation across cosmic time:

- an improvement of a factor of 10-20 in strain sensitivity will allow us to probe the complete BH population to the edge of the universe,
- reaching down to frequencies of about 1 Hz is needed to detect a population of intermediate-mass BHs, and quantitatively probe mergers of such BHs and uncover their link to *LISA* sources and SMBH growth, and
- both the leap in sensitivity and expansion to low frequencies are needed for precise BBH properties, masses, spins, and possibly eccentricities, to firmly uncover their formation origins.



4. Cosmology and the Early History of the Universe

SCIENCE TARGET

Investigate the particle physics of the primeval Universe and probe its dark sectors.

Gravitational waves (GWs) offer unique new probes of the early universe and its composition—the weakness of gravity relative to other known forces implies that GWs decoupled very early from the primordial plasma in the Universe, a fraction of a second after the Big Bang. Detection of such waves would provide unique opportunities to study the evolution of the very early Universe and of the physical laws that apply at very high energy scales, inaccessible in traditional laboratories. Particle physics processes that drove inflation and phase transitions in the early Universe would leave imprints in the form of spectral, spatial, and polarization properties of the stochastic GW background. In the more recent history, merging neutron star (NS) and black hole (BH) binaries generate signals that precisely determine the luminosity distances of the binaries, independently of the cosmic distance ladder. Such coalescing binaries can be used as *standard sirens* to measure the expansion and acceleration rates of the Universe as a function of redshift, thereby inferring fundamental properties of the dark sectors and of gravity itself.

KEY SCIENCE GOALS

Future GW observations will enable exploration of *particle physics, early Universe, and cosmology*:

- **Standard Siren Cosmology.** What is the precise value of the Hubble constant? Is dark energy fully described by a cosmological constant, or does its equation-of-state vary with redshift?
- **Early Universe.** What particle physics laws and energy scales drove the Universe's early evolution? How did it transition from one evolutionary phase into another and to the present Universe?
- **Modified Theories of Gravity.** Do GWs propagate from their sources in the same way as EM waves do? How do modified theories of gravity affect the propagation of GWs from their sources?

Gravity assembles structures in the Universe from the smallest scale of planets to the largest scale of galaxy clusters and the Universe itself. GW observations can, therefore, elucidate the Universe's evolution and its constituents, complementary to the EM, neutrino, and particle observations. Indeed, GW observations forever changed the role of gravity in our exploration of the Universe. In particular, the multi-messenger nature of GW170817 [2, 59] generated a treasure trove of data ushering in a new era in cosmology [8].

General relativity (GR) completely determines the time evolution of the amplitude and frequency of GWs generated by binaries of NSs and BHs. Matched filtering the GW data with the predicted GR waveform readily infers the luminosity distance to the binary's host galaxy [208]. Thus, coalescing binaries have been hailed as *standard sirens*. Just like the traditional standard candles (e.g. type Ia supernovae) they provide a tool for measuring the dynamics of the Universe, but are not susceptible to the systematic biases of the cosmic distance ladder. Standard sirens can be used to make completely independent measurement of the Universe's expansion rate as a function of redshift and to infer cosmological parameters describing the dark sector.

The measured compact binary merger rates [1, 72] imply that future detectors will observe a stochastic background formed from the astrophysical population of binaries at cosmological distances [71, 209]. Such observations will reveal the history of star formation from a time when the Universe was still assembling its first stars and galaxies. Buried under this background could be stochastic signals of primordial origin (see, e.g., [210–220]) that provide insights into the physics and energy scales of the earliest evolutionary phases in the history of the Universe—scales not accessible to current or planned particle physics experiments. The 3G network will probe those energy scales with its excellent sensitivity to stochastic backgrounds.

4.1 Standard Siren Cosmology

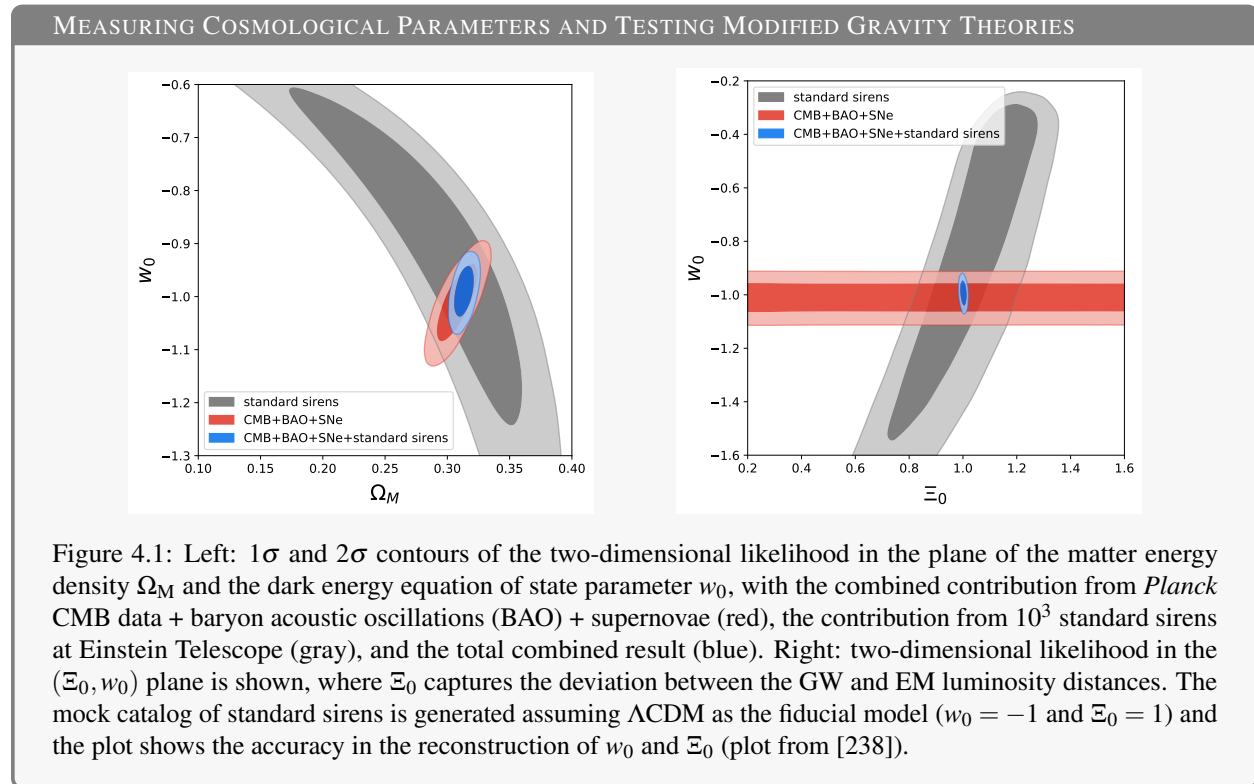
Hubble Constant: GWs from a compact binary merger support a direct measurement of the source’s luminosity distance. Combined with the redshift obtained from the EM counterpart, or with galaxy catalogs, this gives an estimate the present value of the Hubble parameter H_0 [8, 208]. It is estimated that about ten mergers with EM counterparts would be required to reach an accuracy of 5% and 200 to reach 1% [221–223]. While binary NS events are promising based on GW170817, NS-BH and BH-BH mergers due to precession of the orbital plane induced by spin-orbit coupling [224] or the presence of higher multipole modes in the observed waveform [225, 226], can break the degeneracy between the orbital inclination and luminosity distance, and provide more accurate distance measurements. EM observations could also break this degeneracy [227]. There is significant potential in statistical methods as well, where binary mergers without EM counterparts are combined with galaxy catalogs to make inferences [228]. For example, 3G detectors could localize some sources within a volume where on average only one galaxy is present [164, 226, 229], although the method is limited by the peculiar velocity at the redshift of interest and the distance uncertainty $\sim 1\%$ from GW observations.

Precision Cosmology: GWs offer a new approach to probing the dark energy properties either by using redshift measurements from EM counterparts, or by using statistical methods [230–238]. The 3G network will measure standard sirens out to large redshifts, $z \sim 10$, significantly farther than what is possible with the standard candles. Being susceptible to a completely different set of systematic errors from those due to type Ia supernovae this approach will provide a complementary probe precision cosmology and of comparable sensitivity to those due to the cosmic microwave background, baryon acoustic oscillation, and supernovae. An example is shown in Figure 4.1 (left) [237, 238], where addition of standard sirens improves the accuracy of the measured dark energy equation of state parameter w_0 .

Modified Gravity Theories: Standard sirens are sensitive to another powerful signature of the dark energy sector that is not accessible to EM observations. A generic modified gravity theory induces modifications, with respect to the standard model of cosmology, in the cosmological background evolution and perturbations. Indeed, theories with extra dimensions [239], some scalar-tensor theories [240–244], as well as a nonlocal modification of gravity [237, 238, 245–247], are characterized by GWs propagating at the speed of light but with their amplitude decreasing differently with the scale factor than in GR. Consequently, the standard sirens would measure a different luminosity distance compared to their EM counterparts. The 3G network will be sufficiently sensitive to search for this deviation, and probe multiple classes of modified theories of gravity in the context of their dark energy content (Figure 4.1, right) [237, 238].

4.2 Early History of the Universe

Stochastic GW backgrounds could either be astrophysical in origin, generated by a myriad of individual sources, or it could be generated by quantum processes associated with inflation and spontaneous symmetry phase transitions breaking in the early Universe. Figure 4.2 shows examples of GW energy density spectra for some of the cosmological background models in comparison with the best current upper limits and future expected detector sensitivities. Cosmological background is arguably the most fundamentally impactful observation that GW observatories could make. The astrophysical background, however, may mask the

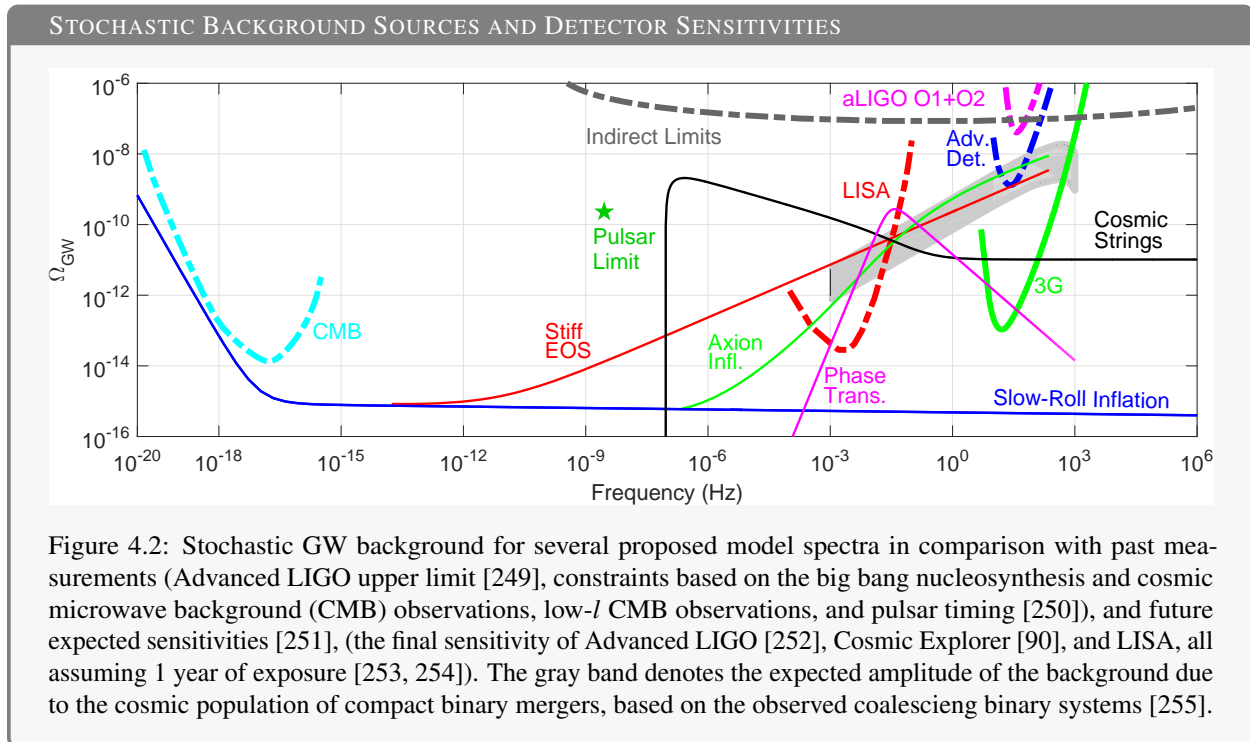


primordial background over much of the accessible spectrum, while still carrying important information about the evolution of structure in the Universe. Techniques are being developed to identify and estimate these various contributions to the stochastic GW background [248].

Irreducible GW background from inflation: Inflation represents the leading framework to explain the properties and initial conditions of the observed Universe. During inflation massless fields experience quantum fluctuations, and due to the accelerated expansion small fluctuations with wavelengths initially smaller than the Hubble radius are amplified and stretched to super-Hubble scales. This applies, in particular, to tensor perturbations [210–212] that re-enter the Hubble radius after inflation and turn into a stochastic GW background. In the standard slow-roll inflationary model (shown in Figure 4.2), the energy density spectrum of this background is likely below the proposed sensitivity of the 3G network, although upgrades to 3G detectors could be sufficiently sensitive to observe this background.

Beyond the irreducible background from inflation: Additional processes during or immediately after inflation could lead to significant amplification of the stochastic background in the frequency band of 3G detectors. For example, coupling of axions to the scalar field of inflation could extend GW production with significant enhancement at higher frequencies [256, 257] as shown in Fig. 4.2. Furthermore, the resulting background is distinctly chiral (i.e. only one polarization is excited) and non-Gaussian [258], unique predictions which can be used to differentiate it unambiguously from other GW backgrounds. Similar trends in the background spectrum are also possible if inflation is followed by a phase with a stiff equation of state ($w > 1/3$), that could be detectable with 3G detectors [213–215, 259].

First order phase transitions: Following the end of inflation, the Universe has undergone quantum chromodynamic, electroweak, and, possibly, other phase transitions. Currently, an experimentally verified physical model lacks energy scales higher than the electroweak scale. Several proposed extensions of the Standard Model predict the occurrence of phase transitions. Any experimental confirmation that such phase transitions took place in the early Universe would constitute a step change in our understanding of particle theory at energy scales inaccessible to terrestrial experiments.



To be an efficient direct source of GWs, a phase transition must be of first order. First-order phase transitions proceed through the nucleation of bubbles of the, energetically more favourable, true vacuum in the space-filling false vacuum. The dynamics of the bubble expansion and collision is phenomenologically rich, and the sources of GWs are the tensor anisotropic stresses generated by these multiple phenomena: the bubble wall's expansion [217, 260], the sound waves in the plasma [219], and the subsequent magnetohydrodynamic turbulence [220, 261]. The nature of the phase transition and its energy scale determine the amplitude and the spectral shape of the GW background. An example of such a background is shown in Figure 4.2 which is potentially within reach of the 3G network [254].

Cosmic Strings: Topological defects such as cosmic strings may arise in the aftermath of a phase transition [262]. Often, the string tension is the only free parameter and it defines the energy scale of the phase transition and the accompanying spontaneous symmetry breaking scale that leads to the formation of cosmic strings. It is also possible to form a network of fundamental cosmic (super)strings. Cosmic strings predominantly decay by the formation of loops and the subsequent GW emission by cosmic string cusps and kinks [263, 264]. Searches for individual bursts of GWs from cosmic strings and for the stochastic background from a string network have placed a strong constraint on the string tension for the three well-known models [265–268]. The 3G network will either detect cosmic strings or improve on these bounds by eight orders of magnitude, depending on the model (see Fig. 4.2).

Dark Photons: A dark photon is proposed to be a light but massive gauge boson in an extension of the Standard Model. If sufficiently light, the local occupation number of the dark photon could be much larger than one, so it can then be treated as a coherently oscillating background field that imposes an oscillating force on objects that carry dark charge. The oscillation frequency is determined by the mass of the dark photon. Such effects could result in a stochastic background that could be measured by 3G detectors, potentially exploring large parts of the parameter space of such models [269].

4.3 Astrophysical Binary Foregrounds and Large Scale Structure

The cosmological population of compact binary mergers will give rise to a stochastic foreground of GWs [270–275]. The amplitude of this foreground is estimated to be $\Omega_{\text{GW}} \sim 10^{-9}$ at 25 Hz, and is likely to be detected by Advanced LIGO and Advanced Virgo [209]. The 3G network, thanks to better low-frequency sensitivity, will probe $\Omega_{\text{GW}} \sim 10^{-13}$ – 10^{-12} [270, 276, 277]. The 3G network could, therefore, detect the predicted spatial anisotropy [278, 279] and non-gaussianity [279] in the energy density of this foreground. In particular, measurements of higher order correlation functions would be extremely useful in understanding the large scale structure of the Universe.

Figure 4.2 shows that the amplitude of this foreground is several orders of magnitude stronger than most cosmological GW backgrounds, hence masking the cosmological signals discussed above. However, a large fraction of the binary merger signals will be individually detected by the 3G network, allowing for the possibility to subtract this foreground to probe a cosmic background of $\Omega_{\text{GW}} \sim 10^{-13}$ [248, 270, 277]. Small errors in subtraction could lead to a substantial residual foreground—novel methods are being explored to enable more effective subtraction [280].

Astrophysical sources other than compact binary mergers could also contribute to the astrophysical GW foreground, including isolated NSs [281–283], core collapse supernovae [284, 285] and population III binaries [286]. Distinguishing these unresolved foreground sources from the cosmological background will require using spatial correlations (both GW-GW and GW-EM correlations) and careful measurements of the foreground spectra. Studying the properties of the astrophysical foreground will provide unique new information on galactic and stellar physics and allow for new types of constraints on astrophysical models.

4.4 Outlook for gravitational-wave cosmology and the early Universe

3G observatories will revolutionize our understanding of the evolution of the Universe, from its earliest moments to the recent past. Measurements of the frequency spectrum, spatial anisotropy, and polarization content of the primordial stochastic GW background are likely to reveal imprints of the physical laws and processes that drove the earliest phases of the Universe’s evolution and take place at energy scales inaccessible to current or planned particle physics experiments. Observation of the astrophysical, compact-binary stochastic GW background and its properties would provide information on the formation and evolution of matter in the Universe, going back to the time when the first stars and galaxies were formed. More recent compact binary mergers will be used as standard sirens to provide novel and independent measurements of the expansion and acceleration of the late Universe, constraining the fundamental nature of gravity and cosmological parameters. While the irreducible inflationary stochastic GW background is likely not within reach of the 3G network, it may be within reach of the follow-up upgrades, which should be factored into the site selection and facility design of the 3G detectors.

SCIENCE REQUIREMENTS

With 3G observatories we can investigate the fundamental physics of the primeval Universe and probe its dark sectors:

- a factor of 10-20 improvement beyond the sensitivity of advanced detectors is critical to infer cosmological parameters at a precision that is competitive with current measurements,
- lowering the frequency response down to 5 Hz will allow the observation of astrophysical foregrounds and their subtraction via better estimation of source parameters, and
- a factor of 10 improvement in strain sensitivity (100 in energy density) over advanced detectors is required to observe GW background from early Universe phase transitions.



5. Extreme Gravity and Fundamental Physics

SCIENCE TARGET

Explore new physics in gravity and in the fundamental properties of compact objects

General Relativity (GR) is a mathematically elegant and physically appealing theory that has been tremendously successful in explaining all relevant astronomical observations and laboratory experiments [287–289]. Nonetheless, it is widely believed that GR is at best incomplete [289, 290], representing an approximation to a more complete theory that cures some or all of its deep conceptual problems. Black hole (BH) information loss, spacetime singularities, cosmological constant and the lack of a viable formulation of quantum gravity, have all added to the suspicion that GR violations will eventually show up in observations [290–293]. Over the past decade new insights into the relationship between entanglement entropy and the architecture of spacetime [294] on the one hand and the connection between asymptotic symmetries, the BH entropy and infrared behaviour of quantum gravity [295] on the other are all hinting towards a modified theory of gravity. At the same time, the discovery of gravitational waves (GWs) has provided a powerful new tool to test GR in a realm of the theory that is inaccessible to other experiments and observations.

GW150914 was not only the first direct detection of GWs but also the first ever observation of a binary black hole (BBH) [296]. Since then tens of BBH mergers have been observed [297], GW190521 being the most massive system discovered so far [128] that converted 9 solar masses to pure energy in a mere 100 ms. BBH mergers are arguably the most powerful phenomena in nature, save for the Big Bang, the signals coming from a region where both the spacetime curvature and gravitational field are extremely large. They have helped test Einstein's gravity in regimes where the theory has never been tested before [80]. 3G observatories will make a step change in studying gravity and the nature of ultra-compact objects.

KEY SCIENCE GOALS

The 3G GW observatories will enable unprecedented and unique science in *extreme gravity* and *fundamental physics*:

- **The nature of gravity:** Are the building-block principles and symmetries in nature, e.g. Lorentz invariance and equivalence principle, invoked in the description of gravity valid at all scales?
- **The nature of compact objects:** Are black holes and neutron stars the only ultra-compact objects in the Universe? If other compact objects exist what are their signatures in gravitational waves?
- **The nature of dark matter:** Is dark matter composed of particles or dark objects or is it a manifestation of failure of general relativity?

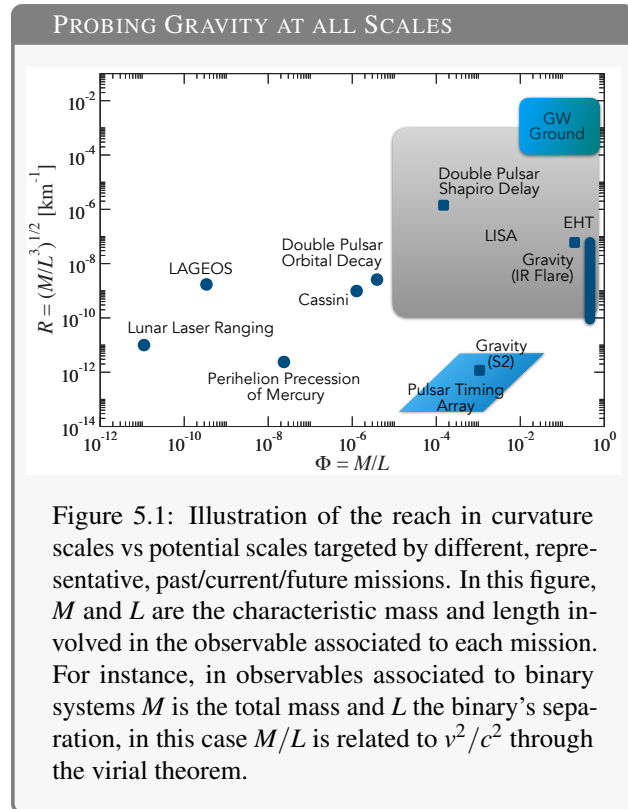
GWs are copiously produced in regions of strong gravity and relativistic motion. Yet the waves carry pristine information about their sources because they interact feebly with matter and remain unscathed as they propagate over billions of light years to Earth. This makes them ideal for testing GR in new ways that could reveal subtle departures from the theory [77, 170, 298, 299]. We can now directly probe the

two-body dynamics (specifically, the complex orbital motion predicted in GR such as precessional effects), the spacetime structure near BHs and the dynamics of their horizons, and if the Universe depicted by electromagnetic (EM) and particle messengers is the same as the one revealed by GWs. In addition, imprint in the observed GWs are the signatures of compact objects that can be deformed in a way that BHs cannot be. This should help us detect a new class of ultra-compact compact objects consisting of boson condensates or other exotic particles and fields should they exist. Finally, the presence of dark matter around compact binaries or their accumulation in neutron star (NS) cores, will also modify the observed signals. Hence, GWs could be used to measure the properties of dark matter particles—their masses and interaction cross sections with hadronic matter.

5.1 Nature of Gravity.

Probing the nature of gravity and its possible implications on fundamental physics is a high-reward, even if uncertain, prospect of GW observations. To our knowledge, astrophysical BHs and relativistic stars exhibit the largest curvature of spacetime accessible to us. They are, therefore, ideal systems to observe the behaviour of spacetimes under the most extreme gravitational conditions. New physics indicative of departures from the basic tenants of GR could reveal itself in high fidelity waveforms expected to be observed in the 3G network. Such signals would provide a unique access to extremely warped spacetimes and gain invaluable insights into GR or what might replace it as the theory of gravity.

Figure 5.1 provides a perspective of the reach of different missions/facilities and their target regime with respect to characteristic spacetime curvature (R) and gravitational potential Φ (which for binary systems can be equated with v^2/c^2 , where v is the characteristic velocity of companion stars and c the speed of light).



To this end, beyond having access to the sensitivity and frequency windows of 3G detectors, guidance from theory, together with further refinements in data analysis, will be of utmost importance to harness this potentially revolutionary opportunity. On the theoretical front a major challenge in extracting new physics with GWs is that, in principle, one needs to model the characteristics of the emitted signal for the desired physical scenarios beyond the framework of GR and then confront it with the data [300, 301]. The powerful perspective of effective field theory [302–305]) allows one to build extensions to GR with higher-order corrections and search for new physics, even before a new fundamental theory and its low-energy phenomenology is fully developed (see, e.g., Refs. [306–308]). Among possible departures under scrutiny are the following:

New fields, particles and polarizations: GR has not only passed every experimental and observational test that it has been subject to but it is also a very robust theory. In fact, any physically meaningful departures from GR necessarily require the presence of extra degrees of freedom, e.g. scalar and vector fields, in addition to the metric tensor [309]. Such additional degrees of freedom also generically arise in the low-energy limit of quantum gravity theories, which often lead to violations of the strong equivalence principle. Among possible theories, those with an additional scalar field, e.g. the Brans-Dicke theory, are relatively simple [310, 311].

Yet they could give rise to exciting new phenomenology in strong gravitational fields indicative of failure of GR [312, 313]. Such theories, therefore, serve as excellent proxies of the type of new physics we can hope to discover. Additionally, if a binary’s companions can become dressed with a scalar configuration [314–317], we can expect the emission of scalar GWs, in addition to tensorial ones, with the dominant component being dipolar emission [300]. Such emissions lead to additional polarizations, beyond the two in GR, that could be detected directly with a network of detectors [76, 318] or indirectly inferred from their effects on the source’s dynamics and consequent impact on the observed GWs [300].

Graviton mass and speed of GWs: Recently, the possibility that gravitons could have mass has resurfaced in theoretical physics within extensions of GR [319, 320]. In a massive graviton theory GWs would be dispersed as they propagate from their sources to Earth, causing a change in the phase evolution of the observed signal relative to GR. The current best bound on the graviton mass m_g comes from modifying the GR dispersion relation for GWs and this sets the bound $m_g < 5.0 \times 10^{-23} \text{ eV}/c^2$ [80]. The 3G network could improve this bound by two orders-of-magnitude by detecting sources from as far as $z \sim 50$ and further constrain massive graviton theories.

Lorentz invariance and parity violations: Lorentz invariance is regarded as a fundamental property of the Standard Model of particle physics, tested to a spectacular accuracy in particle experiments [321]. In the gravitational sector, constraints are far less refined. Theories with Lorentz invariance violation (e.g., Hořava–Lifschitz [322] and Einstein–æther [323]) give rise to significant departures from GR on the properties of BHs [324, 325], existence of additional polarizations [326], and the propagation of GWs through dispersion and birefringence [78, 327]. Furthermore, parity violations in gravity arise naturally within some flavors of string theory [328], loop quantum gravity [329] and inflationary models [304]. The associated phenomenologies are, to some degree, understood from effective theories [330]. For instance, they give rise to BHs with nontrivial pseudo-scalar configurations that violate spatial parity [331]. The resulting scalar dipole leads to a correction to the GWs produced in a binary inspiral and merger signal [332–334]. Additionally, parity violating theories can exhibit birefringence, thus impacting the characteristics of GWs tied to their handedness [335]. All of these effects will be greatly constrained by the 3G network as it will observe sources at redshifts of $z \sim 50$ and higher [335] and will have the ability to measure additional polarizations.

Ultra-light Bosonic Clouds: Ultralight bosons have been proposed in various extensions of the Standard Model [336]. When the Compton wavelength of ultralight bosons (masses in the range 10^{-21} eV – 10^{-11} eV) is comparable to the horizon size of a spinning stellar-mass or supermassive BH, superradiance can cause BH spin to decay, populating bound Bohr orbits around the BH, with an exponentially large number of particles [337–339]. Such bound states, in effect gravitational atoms, have bosonic clouds with masses up to $\sim 10\%$ of the BH mass [340–342]. Once formed, the clouds annihilate over a longer timescale through the emission of coherent, nearly-monochromatic, GWs [340, 343].

Presence of such clouds could be detected via blind searches in the Milky Way [343–347] or directed searches aimed at a candidate BH, such as that formed in a merger event [343–345, 348, 349], or observations of a stochastic background from an unresolved population [346, 347]. Annihilation of such clouds can also impact a binary’s dynamics and thus the GWs produced during the inspiral [350]. Such a modified GW signal is a promising target for 3G detectors with a few to hundreds of events per year expected for bosons in the $\sim 10^{-13}$ – 10^{-12} eV range [344]. Measuring the spin and mass distribution of merging BBHs can provide evidence for characteristic BH spindown from superradiance [344, 345], which would allow exploration of a new parameter space for ultralight bosons [163]. In addition, the presence of such clouds can be probed through the imprint of finite-size effects on the compact objects in a binary system [350]. Some dark-matter candidates alternative to weakly interacting massive particles (e.g., fuzzy dark matter [351], axion-like particles, and other ultralight bosons [336]) predict exotic compact objects (or ECOs) either in the form of boson stars or in the form of condensates that form spontaneously due to BH superradiant instabilities.

Large, non-local, quantum effects: Semi-classical arguments have been put forward to support the possibility of exotic states of matter or dressed compact objects with further structure stemming from quan-

tum gravitational origin. Examples of electromagnetically dark but horizonless compact objects include fuzzballs [352, 353], gravastars [354], dark stars [355, 356], and others [357, 358]. Additionally, new non-local physics at the horizon scale has been suggested by firewall arguments [359] as well as other quantum effects [360–363]. These scenarios can generically give rise to signatures that can potentially be detected by the 3G network. Thus, GWs facilitate a unique window to these arguably speculative ideas, with far reaching consequences if observed.

5.2 Nature of Compact Objects.

Observational evidence so far suggests that compact massive objects in the Universe exist in the form of BHs and NSs. Binary systems composed of such objects provide ideal scenarios to unravel both astrophysical and fundamental physics puzzles such as elucidating the connections of strong gravity with the most energetic phenomena in our Universe, exploring the final state conjecture [364] (namely, the end point of gravitational collapse is a Kerr BH), and probing the existence of horizons.

Nature of black holes: Kerr BHs are characterized by just two parameters, their mass and spin angular momentum. This remarkable property implies that the oscillations of a perturbed BH are rather unique. Indeed, a perturbed BH returns to its quiescent state by losing the energy in its deformation into GWs. The emitted waves consist of a spectrum of damped sinusoids called *quasi-normal modes* whose frequencies and decay times are determined by the BH's mass and spin.

These parameters of the BH can be inferred from the measurement of a *single* mode frequency and its decay time. Detection of several modes would facilitate multiple null tests of the Kerr nature of BHs [168, 365–371]. Such tests look for consistency in the masses and spins inferred from the different modes (see Fig. 5.2). The sensitivity of a 3G network is necessary for the precision with which such consistency tests can be carried out. Multiple loud events expected to be detected by the 3G network will provide exquisite tests of the Kerr nature of compact objects.

Beyond black holes: From a phenomenological standpoint, BHs and NSs are just two species of a larger family of compact objects. More exotic species are theoretically predicted in extensions to GR, but also in particular scenarios within GR [363, 372]. For instance, extremely compact objects (ECOs) arise from beyond-standard model fundamental fields minimally coupled to gravity (e.g., boson stars [373]), in Grand Unified Theories in the early Universe (e.g., cosmic strings [262]), from exotic states of matter, as dressed compact objects with further structure stemming from quantum gravitational origin [360, 362] or new physics at the horizon scale (e.g., firewalls [359]), or as horizonless compact objects in a variety of scenarios, for example, fuzzballs, gravastars, and dark stars [352, 354–358]. GW observations provide a unique discovery opportunity in this context, since exotic matter might not interact electromagnetically or any EM signal from the surface of an ECO might be highly redshifted [363]. Example GW signatures from the inspiral epoch include dipole radiation as well as a variety of matter effects as in the case of NSs [372].

An ECO could be parameterized by the gravitational redshift z_g near its surface. This parameter can

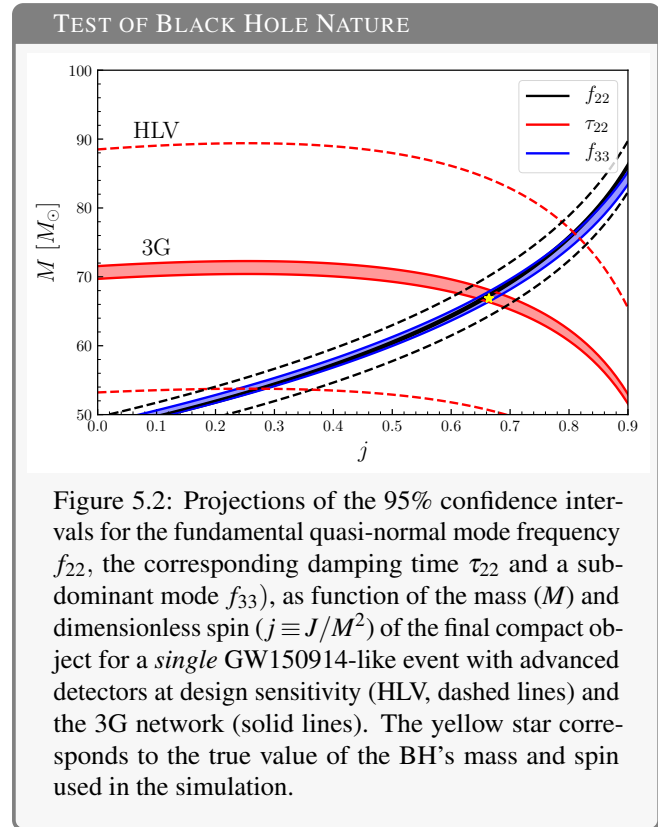


Figure 5.2: Projections of the 95% confidence intervals for the fundamental quasi-normal mode frequency f_{22} , the corresponding damping time τ_{22} and a subdominant mode f_{33} , as function of the mass (M) and dimensionless spin ($j \equiv J/M^2$) of the final compact object for a *single* GW150914-like event with advanced detectors at design sensitivity (HLX, dashed lines) and the 3G network (solid lines). The yellow star corresponds to the true value of the BH's mass and spin used in the simulation.

change by several orders of magnitude depending on the model. BHs have $z_g \rightarrow \infty$ while NSs and the most compact theoretically constructed boson stars have $z_g \sim \mathcal{O}(1)$. For sufficiently large values of z_g compact objects could behave like BHs with increasing precision. Studies of the geodesic motion and quasi-normal modes indicate that ECOs with $z_g \lesssim 1.4$ display internal structure effects that can be discerned in future GW observations. For larger values of z_g , ECOs mimic BHs [363, 374, 375] as departures are redshifted to ever smaller values. Interestingly, models of near-horizon quantum structures—motivated by various scenarios [352–354, 359]—can reach redshifts as high as $z_g \sim \mathcal{O}(10^{20})$ for ECOs in the frequency band of ground-based detectors. GWs could be our only hope to detect or rule them out.

Additionally, while the ringdown signal that follows from the merger of a compact binary can be qualitatively similar to that of a BH, quasi-normal modes of, e.g., gravastars, axion stars and boson stars, are different from Kerr BHs [170]. In addition to gravitational modes, matter modes might be excited in the ringdown of an ECO, akin to the fluid modes excited in a remnant NS [372]. In the case of certain BH mimickers the prompt ringdown signal is identical to that of a BH. However, these objects generically support quasi-bound trapped modes which produce a modulated train of pulses at late time. These modes appear after a delay time whose characteristics are key to test Planckian corrections at the horizon scale. The 3G network will have unprecedented ability extract all these modes and explore the existence of ECOs [363].

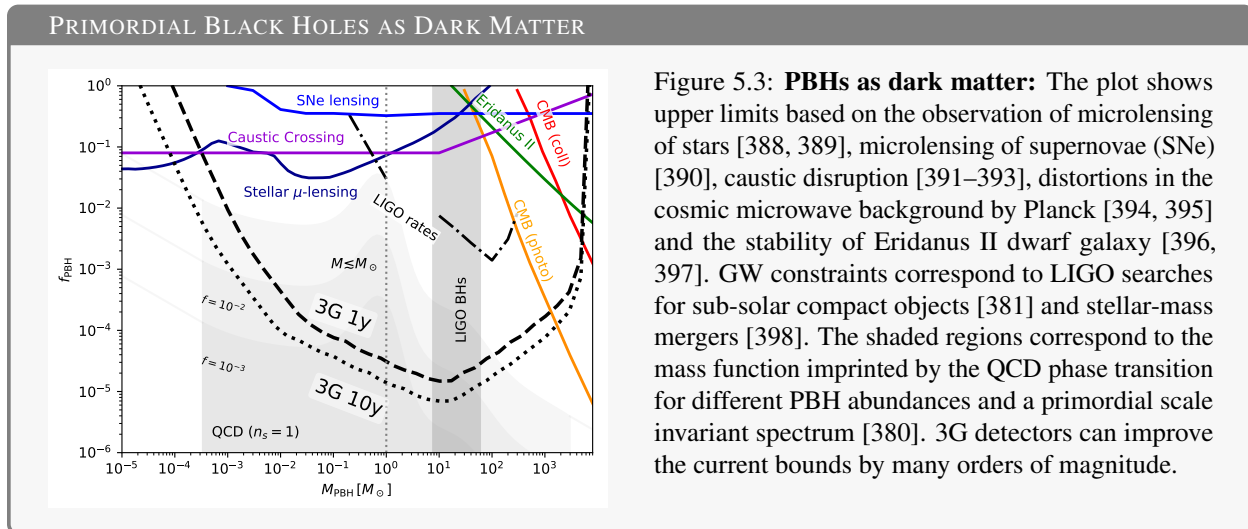
5.3 Nature of Dark Matter.

The exquisite ability of 3G detectors to probe the population and dynamics of electromagnetically dark objects throughout the Universe and harness deep insights on gravity can help reveal the nature of dark matter and answer key questions about its origin.

Black holes as dark matter candidates: LIGO and Virgo discoveries have revived interest in the possibility that dark matter could be composed, in part, of primordial black holes (PBHs) of masses $\sim 0.1\text{--}100M_\odot$ [376–378]. Such BHs might have formed from the collapse of large primordial density fluctuations in the very early Universe or during inflation [103, 379]. The exact distribution of their masses and spins depends on the model of inflation, which might be further affected by processes in the early Universe such as the quantum chromodynamics phase transition (QCD) [294, 380]. Detection of binaries composed of objects much lighter than stellar-mass BHs or mass and spin distributions showing an excess in a certain range, could point towards the existence of PBHs [381]. Identifying mergers at redshifts $z > 30$, when first stars were yet to form, would be another hint towards this formation channel [382]. With the sensitivity to observe mergers at redshifts as large as $z \sim 50$, the 3G network will be uniquely positioned to determine the mass, spin and redshift distributions of BHs, which will be crucial to test the hypothesis that dark matter consists of PBHs [383]. Figure 5.3 shows the current bounds and discovery potential of 3G detectors.

Detection of dark matter with compact objects: Apart from probing whether dark matter can be partially made up of PBHs, GWs can also scrutinize models where dark matter consists of particles beyond the Standard Model, e.g., weakly interacting massive particles [384], fuzzy dark matter [351] or axion-like particles [336]. Indeed, BBHs evolving in a dark-matter rich environment will not only accrete the surrounding material, but also exert a gravitational drag on the dark matter medium, which affects their orbital dynamics [385–387]. Even though their magnitude is small, drag and accretion could have a cumulative effect over a large number of orbits that could be detected by a combination of observatories in space and 3G detectors [372].

Additionally, dark matter that interacts with standard model particles can scatter, lose energy, and be captured in astrophysical objects [399–402]. The dark matter eventually thermalizes with the star, and accumulates inside a finite-size core. The presence of dark matter would change the core’s equation of state imprinting its signature into GWs emitted during the inspiral and merger of such objects in a binary system [403]. In certain models, asymmetric dark matter can accumulate and collapse to a BH in the dense interiors of NSs. The core can grow by accumulating the remaining NS material, in effect turning NSs into light BHs in regions of high dark-matter density such as galactic centers [404, 405]. This provides a



mechanism for creating light BHs that could be observed by 3G detectors. However, BHs that result from implosion of dark matter accreting NSs will always be heavier than $\sim 1 M_{\odot}$, any BH candidates of mass $\lesssim 1 M_{\odot}$ could only be primordial in origin.

5.4 Outlook on Exploring Extreme Gravity and Fundamental Physics

Einstein’s description of gravity led to a revolution in our thinking of the very nature of spacetime itself. Gravity is the manifestation of the curvature of spacetime caused by matter and energy density. GR has so far passed every test to which it has been subject to, yet some of its predictions are deeply troubling. The physical singularity at the Big Bang, loss of information when matter and energy fall into a BH are but examples of predicaments faced by the theory for which no satisfactory resolutions exist. Moreover, observations are hinting that our knowledge of the constituents of the Universe is underwhelmingly poor and breakthroughs in the detection of new particles and fields is keenly awaited.

SCIENCE REQUIREMENTS

GW observations by the 3G network could provide answers to some of the most fundamental questions about spacetime and matter. The capabilities required to answer them are:

- high fidelity observations (SNR>1000) of a large number of sources to discover rare events that carry unique signatures of new particles and fields, dark matter and violations of GR,
- a network of at least one ET and two CEs to constrain or detect additional polarizations,
- access to frequencies below the sensitivity of current detectors to track GWs over a much longer period to test GR predictions and to detect dipole radiation, and
- ability to identify black holes at large redshifts $z \geq 50$ to detect birefringence, improve bounds on graviton mass and discover primordial black holes.



6. Sources at the Frontier of Observations

SCIENCE TARGET

Understand physical processes that underlie the most powerful astrophysical phenomena.

In addition to revealing dark astrophysical processes, gravitational waves (GWs) also offer a new window into extreme phenomena inside compact astronomical sources. While electromagnetic (EM) observations reveal the conditions in the surface regions, GWs encode information about the internal dynamics of these systems, particularly the way the distribution of matter is changing. This will allow us to address longstanding mysteries about the underlying physical mechanisms that power nature's most extreme phenomena. For example, GWs emitted along with light and neutrinos during the aftermath of the collapse of a massive star's core will yield unique insight into why the exterior of the star explodes as a supernova and the nature of the remnant that it leaves behind. GWs from magnetar flares or glitching radio pulsars will elucidate the physical mechanism that causes these bursts of light and provide an unprecedented glimpse deep into the interior of neutron stars (NSs). GWs produced by spinning NSs will allow us to study the structure of ultra-dense NS matter and test particle physics theories in extreme conditions of strong gravity. Sensing these sources of GWs with a network of 3G detectors in concert with EM observatories and neutrino detectors will allow us to learn new physics that is otherwise inaccessible.

KEY SCIENCE GOALS

Observations with the 3G network, further enhanced with EM and neutrino observatories, will allow us to probe new extreme environments and answer key questions about exotic astrophysical transients:

- **GWs from core-collapse supernovae.** How do massive stars explode to form NSs and black holes? What are the physical processes involved and how can they be harnessed to advance our understanding of dense matter, neutrinos, and dark matter?
- **Continuous GW emission from isolated or accreting NSs.** How does dense matter support elastic and magnetic stresses? Does GW emission limit the spin frequencies of NSs?
- **Bursts of GWs from magnetars and other pulsars.** What is the role of magnetic fields in bursts of EM radiation emitted by NSs? How stable is the ultra-dense matter of NSs?

At least one binary NS merger has been detected as a multi-messenger source with a GW signature and emission across the EM spectrum [59]. Although no other multimessenger GW sources have yet been detected, other astrophysical phenomena are expected to produce detectable signatures in multiple messengers. Core-collapse supernovae in our galaxy, rare events with broad implications for astrophysics, nuclear physics and particle physics, remain the most promising multimessenger sources in the Universe. Simulations indicate that the intense GW, neutrino and EM emissions coincident with or following the collapse and explosion of a massive star would produce a large signal in terrestrial detectors, and would allow us to discern fine spectral and temporal features. Spinning asymmetric pulsars, NSs that emit a lighthouse-like beam of EM radiation as they spin, are also expected to be continuous GW sources. A sudden speed up in these otherwise regular

pulses, called *glitches*, and episodic energetic outbursts in highly magnetized NSs, called *giant flares*, may also produce a coincident signal in both EM and GWs. Other fast and energetic phenomena, including fast radio bursts (FRBs) whose origin and prevalence remain most mysterious, may also emit GWs if they are associated with NSs or magnetars. With a network of 3G detectors, each of these potential multimessenger GW sources would yield new insight into key problems in modern physics and astronomy.

6.1 Core-Collapse Supernovae

The physical processes that drive core-collapse supernovae to violently explode and form vast nebula that can seed new stars and produce heavy elements remain mysterious [406]. Simulations that take into account general relativity and the extreme nuclear and neutrino physics needed to model core-collapse supernovae predict that both neutrinos and large scale asymmetric matter flows that are necessary for explosion, also produce strong GW emission [407]. These studies indicate that the formation of the proto-neutron star (PNS) halts collapse, and the shock wave generated by its formation stalls prematurely unless revived by dynamics that breaks spatial symmetries, neutrino heating, and or magnetic field interactions. A GW signal from a galactic core-collapse supernovae, would encode this inner dynamics, and offer vital new information. Furthermore, since neutrinos emitted from the hot and dense PNS are also detectable, GWs provide the complementarity necessary to unlocking longstanding mysteries about the explosion mechanism. Together, GWs and neutrinos would also provide detailed information about fundamental processes that can reveal nuclear and particle physics inaccessible in the laboratory.

Simulations in two (2D) and three (3D) dimensions have shown that GWs are generated by rotational flattening, pulsations of the newly formed PNS, convection, non-radial accretion flows and instabilities, and other asymmetries associated with the effects of strong magnetic fields. The dominant GW emission occurs during a phase of neutrino-driven convection and an instability in the shock wave of matter that has bounced off of the PNS's surface and stalled in place, called *Standing Accretion Shock Instability* (SASI) [407–409]. Oscillations in the star's matter in the near-surface layers of the PNS (the $\ell = 2$, f- and g-modes) [408–410] also contribute strongly to GW emission. Indeed, GW emission from each phase of a core-collapse supernova provides diagnostic constraints on the explosion mechanism and the dynamics of the nascent PNS.

Core collapse and bounce: General-relativistic studies [411–414] show that the GW burst signal from core bounce has generic characteristics for a wide range of progenitor rotation rates and profiles [415] and is a good probe of the bulk parameters of the collapsing iron core [416, 417].

Neutrino-driven turbulent convection outside and inside the PNS: Milliseconds after core bounce, prompt convection in the region between the PNS and standing shock produces a short-period (of the order of tens of ms) burst of GWs peaking at ~ 100 Hz. Subsequent stochastic mass motions that persist for tens of ms post bounce, can lead to significant broadband emission (10–500 Hz with a peak at about 100–200 Hz) [408–410, 418–423]. GWs from the inner PNS convection zone is also broadband and range from 500 Hz to a few kHz [409, 413, 420, 422–424].

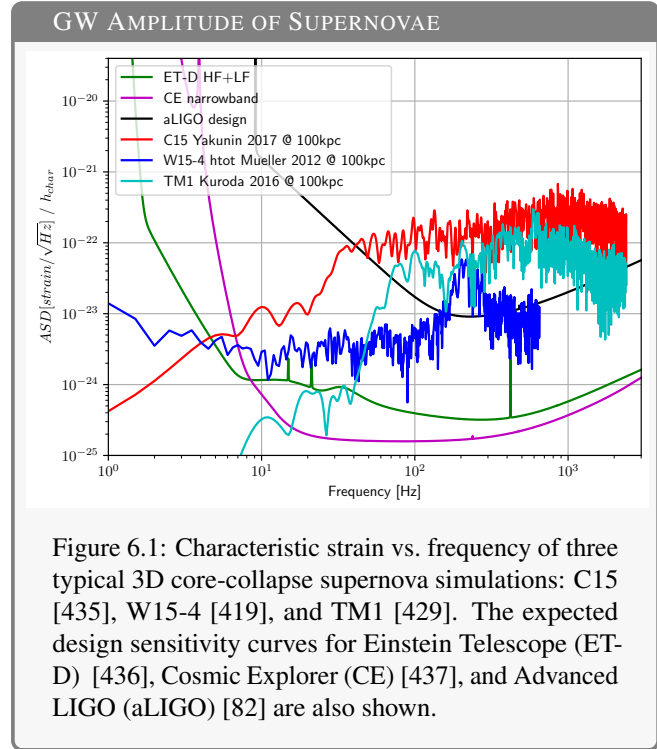
PNS oscillations: The dominant fundamental oscillation modes, the quadrupolar g-modes [425] and the fundamental f-mode, of the nascent PNS excited by accretion or convection and driven by gravity and pressure forces, generate GW emission [409, 425–427]. The frequency and amplitude evolution of the dominant f-mode is largely determined by the PNS mass, radius, and temperature and thus a sensitive probe of the equation of state and of neutrino transport. Simulations indicate GW emission in the frequency band ~ 200 –500 Hz at early times ($< \text{few } 100$ ms) when the core is more extended and ~ 500 –2000 Hz at later stages [419, 425].

SASI: This is an instability associated with shock waves that exists in both 2D and 3D simulations. It is characterized by a nonlinear sloshing mode in 2D, and by both sloshing and spiral modes in 3D [428]. The SASI produces characteristic time modulations both in neutrino and GW signals and provides information about the propagation of the shock wave and explosion dynamics [408, 410, 420–423, 429, 430].

Black hole formation: When a massive rapidly spinning star collapses to form a black hole it will be accompanied by an intense burst of GW emission, followed by a fast ringdown as the newly formed black hole settles down to a Kerr spacetime [431]. By contrast, black hole formation a few seconds after collapse in non-rotating or slowly rotating progenitors is likely to manifest as an abrupt cutoff of the GW emission after their characteristic frequencies (set by oscillation modes of the PNS) increase to several kHz [427, 432].

A ten-fold increase in sensitivity relative to current ground-based detectors would boost the distance at which core-collapse supernova are detectable to 100 kpc, which would include the Large Magellanic Cloud and a number of smaller dwarf galaxies, increasing the chance of detection. Fig. 6.1 plots the spectrum of 3D core-collapse supernova GW signals for a source placed at 100 kpc. They have signal-to-noise ratios for Advanced LIGO at design sensitivity in the range 0.5–6, which is below reliable detectability levels. In contrast, they reach values in the range of 12–130 for the 3G network, levels that would not only allow us to detect the expected signals, but also to determine detailed properties of the progenitor star and the physics of the explosion [433, 434].

In addition to providing critical clues needed to unravel the mystery of how core-collapse supernovae explode, GW observations of these systems will yield a unique insight into the state of hot and dense matter.



GWs emitted by PNS modes and the SASI signals at ~ 100 Hz to 250 Hz will allow us to constrain the physics of hot, ultra-dense nuclear matter in the newborn NS [409, 429]. Recent work has also demonstrated that first-order phase transitions from nuclear to quark matter in the PNS core would imprint unique signatures on the detectable GW emission [438]. The time evolution of the GW frequency will also allow us to chronicle the mass accretion history before and after the shock wave is reignited. The duration and the strength of the GW emission is expected to be correlated with the mass of the progenitor star [409, 422, 423]. With multiple detections the correlations between GWs, neutrinos and EM radiation will help unravel the connection between stellar mass, explosion dynamics, and the nature of the remnant.

Importantly, the onset of the neutrino emissions in core-collapse supernovae coincides with the onset of GW emission to within a few ms [408, 410, 420, 429, 439]. The detection of neutrinos by Super-K/Hyper-K [440], Deep Underground Neutrino Experiment [441], Jiangmen Underground Neutrino Observatory [442], IceCube [443], Large Volume Detector [444], Borexino [445], KamLAND [446], and yet more sensitive neutrino detectors anticipated for the 2030's, will allow us to optimally extract the GW signal [430]. Both signals are produced at the same interior locations, resulting not only in time-coincidence, but have correlated modulations and polarizations, which aid with signal extraction and interpretation. If the progenitor core is rotating, there are additional, distinctive modulation signatures [413]. Thus, joint multimessenger analyses can not only enhance detectability but also more reliably probe physical processes, especially at the highest densities and temperatures. This aspect is critical to studies that aim to harness the extreme conditions encountered in supernovae to either discover or constrain new physics pertaining to dense matter, neutrinos, and dark matter. The core spin could be measured with GW and neutrino multimessenger detections, as the GW frequency is twice the modulation frequency of the neutrino signal [413, 430, 447, 448]. Bounce and

explosion times are very difficult to localize with GWs alone, they require corresponding neutrino detections to pin point the source with accuracy. The 3G network will be critical to extracting all the physics from observable core-collapse supernovae and will synergistically leverage planned telescopes, satellite missions and neutrino detectors.

6.2 Sources of Continuous GWs

The detection of continuous GWs from NSs in the 3G network would provide evidence for deformations and clues about NS structure and their thermal, spin, and magnetic field evolution [449]. The emission of continuous GWs at detectable amplitudes requires a large mass quadrupole in a rapidly rotating compact object. A variety of mechanisms have been proposed to sustain non-axisymmetric distribution of matter and energy, also called ellipticity, in a rotating compact object [450]. Most prominent examples include elastic stresses in the crust, deformations due to magnetic fields, and the growth of r-modes in accreting NSs (a fluid mode of oscillation for which the restoring force is the Coriolis force) [451, 452]. Deformed NSs are nearly monochromatic sources of continuous GWs because the change in frequency is often smaller than $10^{-9} \text{ Hz s}^{-1}$ over long periods of at least a few weeks and typically years. Rapidly spinning NSs are powerful sources since the intensity of the GW radiation is proportional to the sixth power of the spin frequency. Such high spin can be imparted at birth due during a core-collapse supernova [453], or from accretion of matter and angular momentum from a companion star [454].

Isolated NSs: A solid NS crust can sustain (nonaxisymmetric) deformations or ellipticity, whose size depends on the composition, material properties and its evolution [455]. Although maximum fiducial ellipticities $\sim 2 \times 10^{-6}$ can be realized in the crust [456], fiducial ellipticities of $\sim 10^{-9}$ that provide a floor on the spin-down of millisecond pulsars seem more likely [457].

Accreting NSs in Binaries: NSs in binary systems can also emit continuous GWs. They are more likely to present larger deformations than their isolated siblings, due to their internal magnetic fields [458]. Surface magnetic field compressed by infalling material can produce large quadrupolar ellipticity [458] and asymmetric heating in the crust due to accretion could lead to thermal deformations [459]. The excitation of GW-emitting unstable r-mode is also more likely in accreting NSs [460, 461] and may explain why we do not observe NSs spinning at their theoretical upper limit [462–464].

Current observations limit the ellipticity of a canonical NS at 10 kpc emitting GWs above 500 Hz (150 Hz) to be $< 10^{-5}$ (10^{-4}) [465] and targeted searches of known nearby pulsars provide more stringent bounds. For example, observations of the pulsar J0711–6830 require that the ellipticity be less than 1.2×10^{-8} [466]. However, to detect fiducial ellipticities of $\lesssim 10^{-9}$, new detectors with a substantially lower noise floor compared to Advanced LIGO and Virgo are necessary. Multimessenger observations of the source in GW, radio, X-rays, and/or gamma-rays can closely track the spin torques and orbital evolution of Galactic NSs to unravel the physical mechanism responsible for deformations and GW emission.

6.3 GW Bursts Associated with NS Flares and Glitches

NSs can produce GW bursts due to the dynamics associated with giant flares of X-rays produced by highly magnetized NSs or sudden spin-up called glitches in otherwise stably rotating NSs. The coincident detection of GWs from these sources would be transformative as it will allow us to study mechanisms that operate in the deep interior and with unprecedented detail.

Magnetar flares: Magnetars, highly magnetised NSs with magnetic fields exceeding 10^{14} G, are observed as anomalous X-Ray pulsars or soft gamma-ray repeaters [467]. Soft gamma-ray repeaters show recurrent X-ray activity that include frequent short-duration bursts (10^{36} – 10^{43} erg s^{-1} with durations of ~ 0.1 s) and, in some cases, energetic giant flares [468] (10^{44} – 10^{47} erg s^{-1} within 0.1 s with X-ray tails that can extend to several 100 s). To date, three giant flares [469–471] have been detected, and several bursts [472, 473] have been observed that showed quasi-periodic oscillations. Since these events are thought to involve

substantial structural changes within the NSs and due to the large involved energy, magnetars are potential GW sources [474, 475]. If a significant fraction of the X-ray energy is channelled into GWs a magnetar at 10 kpc with a magnetic field at the pole of $B_{\text{pole}} \sim 10^{15}$ G, would produce a strain of $h \sim 10^{-27}$ in the detector. The signal would consist of a high frequency component corresponding to the f-mode around 1–2 kHz, and a low frequency component associated to Alfvén oscillations with frequency ~ 100 Hz, which depends on the magnetic field strength.

Pulsar glitches: Radio pulsars known for their very stable spin periods can occasionally exhibit a sudden increase in their rotation frequency. These are called *glitches* and several hundred glitches have been observed in over 100 pulsars [476]. Physical models for the explanation of glitches involve a substantial rearrangement of the NS structure on a short time scale, and are therefore expected to produce bursts of gravitational radiation. This dynamics is, however, not well understood and the predictions of the emission of GWs and their detectability vary widely. The most optimistic scenarios suggest that a signal should be marginally detectable even by Advanced LIGO and Virgo [477–479], while in pessimistic scenarios even 3G instruments cannot detect the signal [480]. Moderately optimistic scenarios predict the signals to be detectable by the 3G network [481]. As with magnetar flares, the coincident detection of GWs with glitches would be a breakthrough, and even non-detection by 3G instruments would be able to distinguish between different scenarios.

6.4 Outlook for GW sources at the frontier of observations

A network of 3G detectors is essential to correlate multimessenger signals from a host of extreme phenomena from compact sources. It would offer unprecedented opportunities to learn about the birth and extreme behavior of stellar remnants, *and obtain new insight into physical processes in dense matter and those that underlie explosive phenomena such as core-collapse supernovae*. A galactic core collapse stands out as a singular source for multimessenger astrophysics, the coincident detection of GWs, neutrinos and EM radiation will be key to understanding the mechanisms that power core-collapse supernovae explosions. Temporal and spectral correlations between these three messengers will provide definitive insights into the formation of NSs and stellar mass black holes, properties of matter at extreme density and temperature. Continuous GWs from spinning NSs, and bursts of GWs associated with magnetar flares and pulsar glitches also has the potential to unravel the mechanisms at play and the properties of dense matter and extreme magnetic fields. Moreover, with its astounding reach, the potential of 3G detectors is difficult to overstate and serendipitous discoveries are guaranteed to take place.

SCIENCE REQUIREMENTS


Decades after their discoveries much of the physics of most explosive processes in the Universe remains hidden. Understanding these multimessenger phenomena would require:

- GW observatories that are *10–100 times* more sensitive than current ground-based facilities in the 50-1000 Hz range are key to unravel the inner dynamics of supernova explosions.
- Current understanding of compact multimessenger sources such as magnetar quakes and pulsar glitches dictate a factor of ten greater sensitivity than advanced detectors at high frequencies.
- To observe continuous waves from NSs with ellipticities of one part in a billion requires sensitivity improvements by a factor of ten over the entire frequency range.

Acronyms & abbreviations

ACRONYMS AND ABBREVIATIONS USED IN THE TEXT

3G	<i>Third-generation</i> , the next generation of ground-based gravitational-wave observatories consisting of 1 ET in Europe, 1 CE each in the US and Australia
BAO	<i>Baryon Acoustic Oscillations</i>
BBH	<i>Binary Black Hole</i> , a binary system of two black holes
BH	<i>Black Hole</i>
BNS	<i>Binary Neutron Star</i> , a binary system of two neutron stars
CE	<i>Cosmic Explorer</i> , concept for a US third generation interferometer with 40 km arms
CMB	<i>Cosmic Microwave Background</i>
ECO	<i>Exotic Compact Object</i> , an alternative to a neutron star or a black hole
EM	<i>Electromagnetic</i>
ET	<i>Einstein Telescope</i> , concept for a European triangular shaped interferometer with 10 km arms
GR	<i>General relativity</i>
GSF	<i>Gravitational Self-Force</i>
GW	<i>Gravitational Wave</i>
GW150914	Binary black hole merger event detected on 14 September 2015
GW151226	Binary black hole merger event detected on 26 December 2015
GW170104	Binary black hole merger event detected on 4 January 2017
GW170814	Binary black hole merger event detected on 14 August 2017
GW170817	Binary neutron star merger event detected on 17 August 2017
INTEGRAL	<i>International Gamma-Ray Astrophysics Laboratory</i>
LIGO	<i>Laser Interferometer Gravitational-Wave Observatory</i> , 4 km arm length interferometers in the US at Hanford WA and Livingston LA
LIGO-India	LIGO interferometer being built in India
LISA	<i>Laser Interferometer Space Antenna</i>
NR	<i>Numerical Relativity</i>
NS	<i>Neutron Star</i>
NSBH	<i>Neutron Star–Black Hole</i> , a binary system of one neutron star and one black hole
PM	<i>Post-Minkowskian</i>
PNS	<i>Proto-Neutron Star</i>
PN	<i>Post-Newtonian</i>
QCD	<i>Quantum Chromodynamics</i>
SASI	<i>Standing Accretion Shock Instability</i>
SM	the <i>Standard Model</i> of particle physics
SN, SNe	<i>Supernova, Supernovae</i>
SNR	<i>Signal-to-noise ratio</i>
Virgo	3 km arm length interferometer located in Cascina, Italy



Bibliography

- [1] LIGO SCIENTIFIC, VIRGO collaboration, *GW170817: Observation of Gravitational Waves from a Binary Neutron Star Inspiral*, *Phys. Rev. Lett.* **119** (2017) 161101 [1710.05832]. 6, 19
- [2] LIGO SCIENTIFIC, VIRGO, FERMI-GBM, INTEGRAL collaboration, *Gravitational Waves and Gamma-rays from a Binary Neutron Star Merger: GW170817 and GRB 170817A*, *Astrophys. J. Lett.* **848** (2017) L13 [1710.05834]. 6, 18
- [3] A. Goldstein et al., *An Ordinary Short Gamma-Ray Burst with Extraordinary Implications: Fermi-GBM Detection of GRB 170817A*, *Astrophys. J. Lett.* **848** (2017) L14 [1710.05446]. 6
- [4] D. A. Coulter et al., *Swope Supernova Survey 2017a (SSS17a), the Optical Counterpart to a Gravitational Wave Source*, *Science* (2017) [1710.05452]. 6
- [5] C. Freiburghaus, S. Rosswog and F.-K. Thielemann, *R-Process in Neutron Star Mergers*, *ApJ* **525** (1999) L121. 6
- [6] LIGO SCIENTIFIC, VIRGO collaboration, *Estimating the Contribution of Dynamical Ejecta in the Kilonova Associated with GW170817*, *Astrophys. J. Lett.* **850** (2017) L39 [1710.05836].
- [7] M. M. Kasliwal et al., *Illuminating Gravitational Waves: A Concordant Picture of Photons from a Neutron Star Merger*, *Science* **358** (2017) 1559 [1710.05436]. 6, 10
- [8] LIGO SCIENTIFIC, VIRGO, 1M2H, DARK ENERGY CAMERA GW-E, DES, DLT40, LAS CUMBRES OBSERVATORY, VINROUGE, MASTER collaboration, *A gravitational-wave standard siren measurement of the Hubble constant*, *Nature* **551** (2017) 85 [1710.05835]. 6, 18, 19
- [9] K. P. Mooley, A. T. Deller, O. Gottlieb, E. Nakar, G. Hallinan, S. Bourke et al., *Superluminal motion of a relativistic jet in the neutron-star merger GW170817*, *Nature* **561** (2018) 355 [1806.09693]. 6, 10
- [10] LIGO SCIENTIFIC, VIRGO collaboration, *GW170817: Measurements of neutron star radii and equation of state*, *Phys. Rev. Lett.* **121** (2018) 161101 [1805.11581]. 6, 15
- [11] E. Annala, T. Gorda, A. Kurkela and A. Vuorinen, *Gravitational-wave constraints on the neutron-star-matter Equation of State*, *Phys. Rev. Lett.* **120** (2018) 172703 [1711.02644]. 15
- [12] F. J. Fattoyev, J. Piekarewicz and C. J. Horowitz, *Neutron Skins and Neutron Stars in the Multimessenger Era*, *Phys. Rev. Lett.* **120** (2018) 172702 [1711.06615].
- [13] S. De, D. Finstad, J. M. Lattimer, D. A. Brown, E. Berger and C. M. Biwer, *Tidal Deformabilities and Radii of Neutron Stars from the Observation of GW170817*, *Phys. Rev. Lett.* **121** (2018) 091102 [1804.08583].
- [14] I. Tews, J. Margueron and S. Reddy, *Critical examination of constraints on the equation of state of dense matter obtained from GW170817*, *Phys. Rev.* **C98** (2018) 045804 [1804.02783].
- [15] C. D. Capano, I. Tews, S. M. Brown, B. Margalit, S. De, S. Kumar et al., *Stringent constraints on neutron-star radii from multimessenger observations and nuclear theory*, *Nature Astron.* **4** (2020) 625 [1908.10352]. 6

- [16] B. Margalit and B. D. Metzger, *Constraining the Maximum Mass of Neutron Stars From Multi-Messenger Observations of GW170817*, *Astrophys. J. Lett.* **850** (2017) L19 [1710.05938]. 6, 15
- [17] D. Radice, A. Perego, F. Zappa and S. Bernuzzi, *GW170817: Joint Constraint on the Neutron Star Equation of State from Multimessenger Observations*, *Astrophys. J. Lett.* **852** (2018) L29 [1711.03647].
- [18] D. Radice and L. Dai, *Multimessenger Parameter Estimation of GW170817*, *Eur. Phys. J.* **A55** (2019) 50 [1810.12917]. 6
- [19] S. Vitale, W. M. Farr, K. Ng and C. L. Rodriguez, *Measuring the star formation rate with gravitational waves from binary black holes*, *Astrophys. J. Lett.* **886** (2019) L1 [1808.00901]. 6, 13
- [20] J. M. Lattimer and M. Prakash, *The physics of neutron stars*, *Science* **304** (2004) 536 [astro-ph/0405262]. 7
- [21] G. Baym, T. Hatsuda, T. Kojo, P. D. Powell, Y. Song and T. Takatsuka, *From hadrons to quarks in neutron stars: a review*, *Rept. Prog. Phys.* **81** (2018) 056902 [1707.04966]. 7
- [22] A. L. Watts et al., *Colloquium : Measuring the neutron star equation of state using x-ray timing*, *Rev. Mod. Phys.* **88** (2016) 021001 [1602.01081]. 7, 8
- [23] F. Özel and P. Freire, *Masses, Radii, and the Equation of State of Neutron Stars*, *Ann. Rev. Astron. Astrophys.* **54** (2016) 401 [1603.02698]. 7, 8
- [24] S. Gandolfi, J. Carlson and S. Reddy, *The maximum mass and radius of neutron stars and the nuclear symmetry energy*, *Phys. Rev. C* **85** (2012) 032801 [1101.1921]. 7
- [25] K. Hebeler, J. Lattimer, C. Pethick and A. Schwenk, *Equation of state and neutron star properties constrained by nuclear physics and observation*, *Astrophys. J.* **773** (2013) 11 [1303.4662].
- [26] J. M. Lattimer, *The nuclear equation of state and neutron star masses*, *Ann. Rev. Nucl. Part. Sci.* **62** (2012) 485 [1305.3510].
- [27] M. Oertel, M. Hempel, T. Klähn and S. Typel, *Equations of state for supernovae and compact stars*, *Rev. Mod. Phys.* **89** (2017) 015007 [1610.03361].
- [28] I. Tews, J. Carlson, S. Gandolfi and S. Reddy, *Constraining the speed of sound inside neutron stars with chiral effective field theory interactions and observations*, *Astrophys. J.* **860** (2018) 149 [1801.01923]. 7
- [29] P. Demorest, T. Pennucci, S. Ransom, M. Roberts and J. Hessels, *Shapiro Delay Measurement of A Two Solar Mass Neutron Star*, *Nature* **467** (2010) 1081 [1010.5788]. 7
- [30] J. Antoniadis et al., *A Massive Pulsar in a Compact Relativistic Binary*, *Science* **340** (2013) 6131 [1304.6875].
- [31] H. T. Cromartie et al., *Relativistic Shapiro delay measurements of an extremely massive millisecond pulsar*, *Nature Astron.* **4** (2019) 72 [1904.06759]. 7
- [32] J. M. Lattimer and M. Prakash, *What a Two Solar Mass Neutron Star Really Means*, in *From Nuclei to Stars: Festschrift in Honor of Gerald E Brown*, S. Lee, ed., pp. 275–304, (2011), [1012.3208], DOI. 7
- [33] J. Nättilä, M. C. Miller, A. W. Steiner, J. J. E. Kajava, V. F. Suleimanov and J. Poutanen, *Neutron star mass and radius measurements from atmospheric model fits to X-ray burst cooling tail spectra*, *Astron. Astrophys.* **608** (2017) A31 [1709.09120]. 8
- [34] G. Raaijmakers et al., *A NICER view of PSR J0030+0451: Implications for the dense matter equation of state*, *Astrophys. J. Lett.* **887** (2019) L22 [1912.05703]. 8
- [35] S. Bogdanov et al., *Constraining the Neutron Star Mass-Radius Relation and Dense Matter Equation of State with NICER. II. Emission from Hot Spots on a Rapidly Rotating Neutron Star*, *Astrophys. J. Lett.* **887** (2019) L26 [1912.05707]. 8

- [36] K. D. Kokkotas and G. Schaefer, *Tidal and tidal resonant effects in coalescing binaries*, *Mon. Not. Roy. Astron. Soc.* **275** (1995) 301 [gr-qc/9502034]. 8
- [37] D. Lai, *Resonant oscillations and tidal heating in coalescing binary neutron stars*, *Mon. Not. Roy. Astron. Soc.* **270** (1994) 611 [astro-ph/9404062].
- [38] M. Shibata, *Effects of tidal resonances in coalescing compact binary systems*, *Prog. Theor. Phys.* **91** (1994) 871.
- [39] E. E. Flanagan and E. Racine, *Gravitomagnetic resonant excitation of Rossby modes in coalescing neutron star binaries*, *Phys. Rev.* **D75** (2007) 044001 [gr-qc/0601029]. 8
- [40] E. Poisson, *Gravitational waves from inspiraling compact binaries: The Quadrupole moment term*, *Phys. Rev.* **D57** (1998) 5287 [gr-qc/9709032]. 8
- [41] P. Landry, *Rotational-tidal phasing of the binary neutron star waveform*, [1805.01882]. 8
- [42] T. Abdelsalhin, L. Gualtieri and P. Pani, *Post-Newtonian spin-tidal couplings for compact binaries*, *Phys. Rev.* **D98** (2018) 104046 [1805.01487]. 8
- [43] J. M. Lattimer and D. N. Schramm, *Black-hole-neutron-star collisions*, *Astrophys. J. Lett.* **192** (1974) L145. 8
- [44] M. Shibata and K. Taniguchi, *Coalescence of Black Hole-Neutron Star Binaries*, *Living Rev. Rel.* **14** (2011) 6. 8
- [45] M. D. Duez, *Numerical relativity confronts compact neutron star binaries: a review and status report*, *Class. Quant. Grav.* **27** (2010) 114002 [0912.3529]. 8
- [46] J. A. Faber and F. A. Rasio, *Binary Neutron Star Mergers*, *Living Rev. Rel.* **15** (2012) 8 [1204.3858].
- [47] V. Paschalidis, *General relativistic simulations of compact binary mergers as engines for short gamma-ray bursts*, *Class. Quant. Grav.* **34** (2017) 084002 [1611.01519].
- [48] L. Baiotti and L. Rezzolla, *Binary neutron star mergers: a review of Einstein's richest laboratory*, *Rept. Prog. Phys.* **80** (2017) 096901 [1607.03540]. 8
- [49] Z.-G. Xing, J. M. Centrella and S. L. W. McMillan, *Gravitational radiation from coalescing binary neutron stars*, *Phys. Rev.* **D50** (1994) 6247 [gr-qc/9411029]. 8
- [50] M. Shibata and K. Uryu, *Simulation of merging binary neutron stars in full general relativity: Gamma = two case*, *Phys. Rev.* **D61** (2000) 064001 [gr-qc/9911058].
- [51] R. Oechslin, S. Rosswog and F. K. Thielemann, *Conformally flat smoothed particle hydrodynamics: application to neutron star mergers*, *Phys. Rev.* **D65** (2002) 103005 [gr-qc/0111005].
- [52] M. Shibata and K. Uryu, *Gravitational waves from the merger of binary neutron stars in a fully general relativistic simulation*, *Prog. Theor. Phys.* **107** (2002) 265 [gr-qc/0203037].
- [53] A. Bauswein and H. T. Janka, *Measuring neutron-star properties via gravitational waves from binary mergers*, *Phys. Rev. Lett.* **108** (2012) 011101 [1106.1616].
- [54] A. Bauswein, H. T. Janka, K. Hebeler and A. Schwenk, *Equation-of-state dependence of the gravitational-wave signal from the ring-down phase of neutron-star mergers*, *Phys. Rev.* **D86** (2012) 063001 [1204.1888].
- [55] K. Takami, L. Rezzolla and L. Baiotti, *Spectral properties of the post-merger gravitational-wave signal from binary neutron stars*, *Phys. Rev.* **D91** (2015) 064001 [1412.3240].
- [56] S. Bernuzzi, T. Dietrich and A. Nagar, *Modeling the complete gravitational wave spectrum of neutron star mergers*, *Phys. Rev. Lett.* **115** (2015) 091101 [1504.01764].
- [57] L. Lehner, S. L. Liebling, C. Palenzuela, O. L. Caballero, E. O'Connor, M. Anderson et al., *Unequal mass binary neutron star mergers and multimessenger signals*, *Class. Quant. Grav.* **33** (2016) 184002 [1603.00501].

- [58] F. Maione, R. De Pietri, A. Feo and F. Löffler, *Spectral analysis of gravitational waves from binary neutron star merger remnants*, *Phys. Rev.* **D96** (2017) 063011 [1707.03368]. 8
- [59] LIGO SCIENTIFIC, VIRGO, FERMI GBM, INTEGRAL, ICECUBE, ASTROSAT CADMIUM ZINC TELLURIDE IMAGER TEAM, IPN, INSIGHT-HXMT, ANTARES, SWIFT, AGILE TEAM, 1M2H TEAM, DARK ENERGY CAMERA GW-EM, DES, DLT40, GRAWITA, FERMI-LAT, ATCA, ASKAP, LAS CUMBRES OBSERVATORY GROUP, OZGRAV, DWF (DEEPER WIDER FASTER PROGRAM), AST3, CAASTRO, VINROUGE, MASTER, J-GEM, GROWTH, JAGWAR, CALTECHNRAO, TTU-NRAO, NUSTAR, PAN-STARRS, MAXI TEAM, TZAC CONSORTIUM, KU, NORDIC OPTICAL TELESCOPE, ePESSTO, GROND, TEXAS TECH UNIVERSITY, SALT GROUP, TOROS, BOOTES, MWA, CALET, IKI-GW FOLLOW-UP, H.E.S.S., LOFAR, LWA, HAWC, PIERRE AUGER, ALMA, EURO VLBI TEAM, PI OF SKY, CHANDRA TEAM AT MCGILL UNIVERSITY, DFN, ATLAS TELESCOPES, HIGH TIME RESOLUTION UNIVERSE SURVEY, RIMAS, RATIR, SKA SOUTH AFRICA/MEERKAT collaboration, *Multi-messenger Observations of a Binary Neutron Star Merger*, *Astrophys. J. Lett.* **848** (2017) L12 [1710.05833]. 9, 18, 29
- [60] D. Kasen, B. Metzger, J. Barnes, E. Quataert and E. Ramirez-Ruiz, *Origin of the heavy elements in binary neutron-star mergers from a gravitational wave event*, *Nature* (2017) [1710.05463]. 9
- [61] A. P. Ji, A. Frebel, A. Chiti and J. D. Simon, *R-process enrichment from a single event in an ancient dwarf galaxy*, *Nature* **531** (2016) 610 [1512.01558]. 9
- [62] H. K. Chaurasia and M. Bailes, *On the eccentricities and merger rates of double neutron star binaries and the creation of double supernovae*, *Astrophys. J.* **632** (2005) 1054 [astro-ph/0504021]. 10
- [63] J. S. Bloom, S. Sigurdsson and O. R. Pols, *The Spatial distribution of coalescing neutron star binaries: Implications for gamma-ray bursts*, *Mon. Not. Roy. Astron. Soc.* **305** (1999) 763 [astro-ph/9805222].
- [64] LIGO SCIENTIFIC, VIRGO collaboration, *On the Progenitor of Binary Neutron Star Merger GW170817*, *Astrophys. J. Lett.* **850** (2017) L40 [1710.05838]. 10
- [65] P. Madau and M. Dickinson, *Cosmic Star Formation History*, *Ann. Rev. Astron. Astrophys.* **52** (2014) 415 [1403.0007]. 10, 13, 15
- [66] G. Ghirlanda, O. S. Salafia, Z. Paragi, M. Giroletti, J. Yang, B. Marcote et al., *Compact radio emission indicates a structured jet was produced by a binary neutron star merger*, *Science* **363** (2019) 968 [1808.00469]. 10
- [67] R. Margutti et al., *The Binary Neutron Star Event LIGO/Virgo GW170817 160 Days after Merger: Synchrotron Emission across the Electromagnetic Spectrum*, *Astrophys. J. Lett.* **856** (2018) L18 [1801.03531].
- [68] G. P. Lamb et al., *The optical afterglow of GW170817 at one year post-merger*, *Astrophys. J. Lett.* **870** (2019) L15 [1811.11491]. 10
- [69] E. Nakar, O. Gottlieb, T. Piran, M. M. Kasliwal and G. Hallinan, *From γ to Radio - The Electromagnetic Counterpart of GW 170817*, *Astrophys. J.* **867** (2018) 18 [1803.07595]. 10
- [70] K. P. Mooley et al., *A mildly relativistic wide-angle outflow in the neutron star merger GW170817*, *Nature* **554** (2018) 207 [1711.11573]. 10
- [71] LIGO SCIENTIFIC, VIRGO collaboration, *Astrophysical Implications of the Binary Black-Hole Merger GW150914*, *Astrophys. J. Lett.* **818** (2016) L22 [1602.03846]. 12, 19
- [72] LIGO SCIENTIFIC, VIRGO collaboration, *Binary Black Hole Population Properties Inferred from the First and Second Observing Runs of Advanced LIGO and Advanced Virgo*, *Astrophys. J. Lett.* **882** (2019) L24 [1811.12940]. 13, 15, 16, 19
- [73] LIGO SCIENTIFIC, VIRGO collaboration, *GW190412: Observation of a Binary-Black-Hole Coalescence with Asymmetric Masses*, *Phys. Rev.* **D102** (2020) 043015 [2004.08342].

- [74] LIGO SCIENTIFIC, VIRGO collaboration, *Properties and astrophysical implications of the 150 Msun binary black hole merger GW190521*, *Astrophys. J. Lett.* **900** (2020) L13 [2009.01190].
- [75] LIGO SCIENTIFIC, VIRGO collaboration, *GW190814: Gravitational Waves from the Coalescence of a 23 Solar Mass Black Hole with a 2.6 Solar Mass Compact Object*, *Astrophys. J. Lett.* **896** (2020) L44 [2006.12611]. 12, 15
- [76] LIGO SCIENTIFIC, VIRGO collaboration, *Tests of general relativity with GW150914*, *Phys. Rev. Lett.* **116** (2016) 221101 [1602.03841]. 12, 25
- [77] N. Yunes, K. Yagi and F. Pretorius, *Theoretical Physics Implications of the Binary Black-Hole Mergers GW150914 and GW151226*, *Phys. Rev.* **D94** (2016) 084002 [1603.08955]. 23
- [78] LIGO SCIENTIFIC, VIRGO collaboration, *GW170104: Observation of a 50-Solar-Mass Binary Black Hole Coalescence at Redshift 0.2*, *Phys. Rev. Lett.* **118** (2017) 221101 [1706.01812]. 25
- [79] LIGO SCIENTIFIC, VIRGO collaboration, *GW170814: A Three-Detector Observation of Gravitational Waves from a Binary Black Hole Coalescence*, *Phys. Rev. Lett.* **119** (2017) 141101 [1709.09660].
- [80] LIGO SCIENTIFIC, VIRGO collaboration, *Tests of General Relativity with the Binary Black Hole Signals from the LIGO-Virgo Catalog GWTC-1*, *Phys. Rev.* **D100** (2019) 104036 [1903.04467]. 12, 23, 25
- [81] KAGRA collaboration, *KAGRA: 2.5 Generation Interferometric Gravitational Wave Detector*, *Nat. Astron.* **3** (2019) 35 [1811.08079]. 12
- [82] KAGRA, LIGO SCIENTIFIC, VIRGO collaboration, *Prospects for Observing and Localizing Gravitational-Wave Transients with Advanced LIGO, Advanced Virgo and KAGRA*, *Living Rev. Rel.* **21** (2018) 3 [1304.0670]. 12, 13, 31
- [83] I. Mandel and R. O’Shaughnessy, *Compact Binary Coalescences in the Band of Ground-based Gravitational-Wave Detectors*, *Class. Quant. Grav.* **27** (2010) 114007 [0912.1074]. 12, 13
- [84] S. Stevenson, C. P. L. Berry and I. Mandel, *Hierarchical analysis of gravitational-wave measurements of binary black hole spin-orbit misalignments*, *Mon. Not. Roy. Astron. Soc.* **471** (2017) 2801 [1703.06873].
- [85] C. Talbot and E. Thrane, *Determining the population properties of spinning black holes*, *Phys. Rev.* **D96** (2017) 023012 [1704.08370]. 13
- [86] J. W. Barrett, S. M. Gaebel, C. J. Neijssel, A. Vigna-Gómez, S. Stevenson, C. P. L. Berry et al., *Accuracy of inference on the physics of binary evolution from gravitational-wave observations*, *Mon. Not. Roy. Astron. Soc.* **477** (2018) 4685 [1711.06287]. 13, 15
- [87] M. Zevin, C. Pankow, C. L. Rodriguez, L. Sampson, E. Chase, V. Kalogera et al., *Constraining Formation Models of Binary Black Holes with Gravitational-Wave Observations*, *Astrophys. J.* **846** (2017) 82 [1704.07379]. 15
- [88] M. Fishbach, D. E. Holz and W. M. Farr, *Does the Black Hole Merger Rate Evolve with Redshift?*, *Astrophys. J. Lett.* **863** (2018) L41 [1805.10270]. 12, 13
- [89] B. Sathyaprakash et al., *Scientific Objectives of Einstein Telescope*, *Class. Quant. Grav.* **29** (2012) 124013 [1206.0331]. 13
- [90] LIGO SCIENTIFIC collaboration, *Exploring the Sensitivity of Next Generation Gravitational Wave Detectors*, *Class. Quant. Grav.* **34** (2017) 044001 [1607.08697]. 13, 21
- [91] J. Roulet, T. Venumadhav, B. Zackay, L. Dai and M. Zaldarriaga, *Binary Black Hole Mergers from LIGO/Virgo O1 and O2: Population Inference Combining Confident and Marginal Events*, [2008.07014]. 13

- [92] C. L. Rodriguez, M. Zevin, C. Pankow, V. Kalogera and F. A. Rasio, *Illuminating Black Hole Binary Formation Channels with Spins in Advanced LIGO*, *Astrophys. J. Lett.* **832** (2016) L2 [1609.05916]. 13
- [93] S. Vitale, R. Lynch, R. Sturani and P. Graff, *Use of gravitational waves to probe the formation channels of compact binaries*, *Class. Quant. Grav.* **34** (2017) 03LT01 [1503.04307].
- [94] C. Kimball, C. P. L. Berry and V. Kalogera, *What GW170729's exceptional mass and spin tells us about its family tree*, [1903.07813].
- [95] S. S. Bavera, T. Fragos, Y. Qin, E. Zapartas, C. J. Neijssel, I. Mandel et al., *The origin of spin in binary black holes: Predicting the distributions of the main observables of Advanced LIGO*, *Astron. Astrophys.* **635** (2020) A97 [1906.12257].
- [96] M. A. Sedda, M. Mapelli, M. Spera, M. Benacquista and N. Giacobbo, *Fingerprints of binary black hole formation channels encoded in the mass and spin of merger remnants*, *Astrophys. J.* **894** (2020) 133 [2003.07409]. 13
- [97] M. Dominik, K. Belczynski, C. Fryer, D. E. Holz, E. Berti, T. Bulik et al., *Double Compact Objects II: Cosmological Merger Rates*, *Astrophys. J.* **779** (2013) 72 [1308.1546]. 13
- [98] M. Mapelli, N. Giacobbo, E. Ripamonti and M. Spera, *The cosmic merger rate of stellar black hole binaries from the Illustris simulation*, *Mon. Not. Roy. Astron. Soc.* **472** (2017) 2422 [1708.05722].
- [99] M. Chruslinska, G. Nelemans and K. Belczynski, *The influence of the distribution of cosmic star formation at different metallicities on the properties of merging double compact objects*, *Mon. Not. Roy. Astron. Soc.* **482** (2019) 5012 [1811.03565].
- [100] M. Mapelli and N. Giacobbo, *The cosmic merger rate of neutron stars and black holes*, *Mon. Not. Roy. Astron. Soc.* **479** (2018) 4391 [1806.04866].
- [101] C. J. Neijssel, A. Vigna-Gómez, S. Stevenson, J. W. Barrett, S. M. Gaebel, F. Broekgaarden et al., *The effect of the metallicity-specific star formation history on double compact object mergers*, *Mon. Not. Roy. Astron. Soc.* **490** (2019) 3740 [1906.08136].
- [102] F. Santoliquido, M. Mapelli, N. Giacobbo, Y. Bouffanais and M. C. Artale, *The cosmic merger rate density of compact objects: impact of star formation, metallicity, initial mass function and binary evolution*, [2009.03911]. 13
- [103] M. Sasaki, T. Suyama, T. Tanaka and S. Yokoyama, *Primordial black holes—perspectives in gravitational wave astronomy*, *Class. Quant. Grav.* **35** (2018) 063001 [1801.05235]. 13, 27
- [104] G. Scelfo, N. Bellomo, A. Raccanelli, S. Matarrese and L. Verde, *GW×LSS: chasing the progenitors of merging binary black holes*, *JCAP* **1809** (2018) 039 [1809.03528]. 13
- [105] M. Dominik, K. Belczynski, C. Fryer, D. Holz, E. Berti, T. Bulik et al., *Double Compact Objects I: The Significance of the Common Envelope on Merger Rates*, *Astrophys. J.* **759** (2012) 52 [1202.4901]. 13
- [106] T. Kinugawa, A. Miyamoto, N. Kanda and T. Nakamura, *The detection rate of inspiral and quasi-normal modes of Population III binary black holes which can confirm or refute the general relativity in the strong gravity region*, *Mon. Not. Roy. Astron. Soc.* **456** (2016) 1093 [1505.06962].
- [107] I. Mandel and S. E. de Mink, *Merging binary black holes formed through chemically homogeneous evolution in short-period stellar binaries*, *Mon. Not. Roy. Astron. Soc.* **458** (2016) 2634 [1601.00007].
- [108] P. Marchant, N. Langer, P. Podsiadlowski, T. M. Tauris and T. J. Moriya, *A new route towards merging massive black holes*, *Astron. Astrophys.* **588** (2016) A50 [1601.03718].
- [109] C. L. Rodriguez, C.-J. Haster, S. Chatterjee, V. Kalogera and F. A. Rasio, *Dynamical Formation of the GW150914 Binary Black Hole*, *Astrophys. J. Lett.* **824** (2016) L8 [1604.04254].

- [110] G. Fragione and B. Kocsis, *Black hole mergers from an evolving population of globular clusters*, *Phys. Rev. Lett.* **121** (2018) 161103 [1806.02351].
- [111] Y. Qin, T. Fragos, G. Meynet, J. Andrews, M. Sørensen and H. F. Song, *The spin of the second-born black hole in coalescing binary black holes*, *Astron. Astrophys.* **616** (2018) A28 [1802.05738].
- [112] C. L. Rodriguez and A. Loeb, *Redshift Evolution of the Black Hole Merger Rate from Globular Clusters*, *Astrophys. J. Lett.* **866** (2018) L5 [1809.01152].
- [113] N. Choksi, M. Volonteri, M. Colpi, O. Y. Gnedin and H. Li, *The star clusters that make black hole binaries across cosmic time*, *Astrophys. J.* **873** (2019) 100 [1809.01164].
- [114] M. Safarzadeh and K. Hotokezaka, *Being Careful with the Field Formation Interpretation of GW190412*, *Astrophys. J. Lett.* **897** (2020) L7 [2005.06519]. 13
- [115] L. Barsotti, L. McCuller, M. Evans and P. Fritschel, *The A+ design curve*, Tech. Rep. LIGO-T1800042, LIGO, Pasadena, CA, 2018. 13
- [116] D. J. Whalen, C. L. Fryer, D. E. Holz, A. Heger, S. E. Woosley, M. Stiavelli et al., *Seeing the First Supernovae at the Edge of the Universe with JWST*, *Astrophys. J. Lett.* **762** (2013) L6 [1209.3457]. 14
- [117] L. V. E. Koopmans et al., *The Cosmic Dawn and Epoch of Reionization with the Square Kilometre Array*, *PoS AASKA14* (2015) 001 [1505.07568].
- [118] R. Cassano et al., *SKA-Athena Synergy White Paper*, [1807.09080].
- [119] J. Kalirai, *Scientific Discovery with the James Webb Space Telescope*, *Contemp. Phys.* **59** (2018) 251 [1805.06941].
- [120] H. Katz, T. P. Galligan, T. Kimm, J. Rosdahl, M. G. Haehnelt, J. Blaizot et al., *Probing Cosmic Dawn with Emission Lines: Predicting Infrared and Nebular Line Emission for ALMA and JWST*, *Mon. Not. Roy. Astron. Soc.* **487** (2019) 5902 [1901.01272]. 14
- [121] LISA collaboration, *Laser Interferometer Space Antenna*, [1702.00786]. 14
- [122] A. Sesana, *Prospects for Multiband Gravitational-Wave Astronomy after GW150914*, *Phys. Rev. Lett.* **116** (2016) 231102 [1602.06951]. 14, 16
- [123] P. Amaro-Seoane and L. Santamaria, *Detection of IMBHs with ground-based gravitational wave observatories: A biography of a binary of black holes, from birth to death*, *Astrophys. J.* **722** (2010) 1197 [0910.0254]. 14
- [124] A. Klein et al., *Science with the space-based interferometer eLISA: Supermassive black hole binaries*, *Phys. Rev.* **D93** (2016) 024003 [1511.05581]. 14
- [125] S. Babak, J. Gair, A. Sesana, E. Barausse, C. F. Sopuerta, C. P. L. Berry et al., *Science with the space-based interferometer LISA. V: Extreme mass-ratio inspirals*, *Phys. Rev.* **D95** (2017) 103012 [1703.09722]. 14
- [126] D. Barret et al. in *Athena+: The first Deep Universe X-ray Observatory*, 2013, [1310.3814]. 14
- [127] LYNX TEAM collaboration, *The Lynx Mission Concept Study Interim Report*, [1809.09642]. 14
- [128] LIGO SCIENTIFIC, VIRGO collaboration, *GW190521: A Binary Black Hole Merger with a Total Mass of 150 M_{\odot}* , *Phys. Rev. Lett.* **125** (2020) 101102 [2009.01075]. 14, 15, 23
- [129] M. Volonteri, *The Formation and Evolution of Massive Black Holes*, *Science* **337** (2012) 544 [1208.1106]. 14
- [130] M. A. Latif and A. Ferrara, *Formation of supermassive black hole seeds*, *Publ. Astron. Soc. Austral.* **33** (2016) e051 [1605.07391].

- [131] J. L. Bernal, A. Raccanelli, L. Verde and J. Silk, *Signatures of primordial black holes as seeds of supermassive black holes*, *JCAP* **1805** (2018) 017 [1712.01311].
- [132] T. E. Woods et al., *Titans of the Early Universe: The Prato Statement on the Origin of the First Supermassive Black Holes*, *Publ. Astron. Soc. Austral.* **36** (2019) e027 [1810.12310]. 14
- [133] E. Banados et al., *An 800-million-solar-mass black hole in a significantly neutral Universe at redshift 7.5*, *Nature* **553** (2018) 473 [1712.01860]. 14
- [134] A. Merloni, *Observing Supermassive Black Holes across cosmic time: from phenomenology to physics*, *Lect. Notes Phys.* **905** (2016) 101 [1505.04940]. 14
- [135] J. Kormendy and L. C. Ho, *Coevolution (Or Not) of Supermassive Black Holes and Host Galaxies*, *Ann. Rev. Astron. Astrophys.* **51** (2013) 511 [1304.7762]. 14
- [136] D. D. Nguyen, A. C. Seth, M. den Brok, N. Neumayer, M. Cappellari, A. J. Barth et al., *Improved Dynamical Constraints on the Mass of the Central Black Hole in NGC 404*, *ApJ* **836** (2017) 237 [1610.09385]. 14
- [137] V. F. Baldassare, A. E. Reines, E. Gallo and J. E. Greene, *A 50,000 Solar Mass Black Hole in the Nucleus of RGG 118*, *ApJ* **809** (2015) L14 [1506.07531].
- [138] A. W. Graham, R. Soria and B. L. Davis, *Expected intermediate-mass black holes in the Virgo cluster – II. Late-type galaxies*, *Mon. Not. Roy. Astron. Soc.* **484** (2019) 814 [1811.03232].
- [139] A. W. Graham and R. Soria, *Expected intermediate-mass black holes in the Virgo cluster – I. Early-type galaxies*, *Mon. Not. Roy. Astron. Soc.* **484** (2019) 794 [1812.01231]. 14
- [140] N. J. McConnell, C.-P. Ma, K. Gebhardt, S. A. Wright, J. D. Murphy, T. R. Lauer et al., *Two ten-billion-solar-mass black holes at the centres of giant elliptical galaxies*, *Nature* **480** (2011) 215 [1112.1078]. 14
- [141] X.-B. Wu, F. Wang, X. Fan, W. Yi, W. Zuo, F. Bian et al., *An ultraluminous quasar with a twelve-billion-solar-mass black hole at redshift 6.30*, *Nature* **518** (2015) 512 [1502.07418].
- [142] M. Mezcuca, J. Hlavacek-Larrondo, J. R. Lucey, M. T. Hogan, A. C. Edge and B. R. McNamara, *The most massive black holes on the Fundamental Plane of black hole accretion*, *MNRAS* **474** (2018) 1342 [1710.10268]. 14
- [143] S. Hirano, T. Hosokawa, N. Yoshida, K. Omukai and H. W. Yorke, *Primordial star formation under the influence of far ultraviolet radiation: 1540 cosmological haloes and the stellar mass distribution*, *Mon. Not. Roy. Astron. Soc.* **448** (2015) 568 [1501.01630]. 14
- [144] M. A. Latif, D. R. G. Schleicher, W. Schmidt and J. C. Niemeyer, *The characteristic black hole mass resulting from direct collapse in the early universe*, *Mon. Not. Roy. Astron. Soc.* **436** (2013) 2989 [1309.1097]. 14
- [145] M. Dijkstra, A. Ferrara and A. Mesinger, *Feedback-regulated supermassive black hole seed formation*, *MNRAS* **442** (2014) 2036 [1405.6743].
- [146] H. Umeda, T. Hosokawa, K. Omukai and N. Yoshida, *The Final Fates of Accreting Supermassive Stars*, *Astrophys. J. Lett.* **830** (2016) L34 [1609.04457].
- [147] M. Habouzit, M. Volonteri, M. Latif, Y. Dubois and S. Peirani, *On the number density of 'direct collapse' black hole seeds*, *Mon. Not. Roy. Astron. Soc.* **463** (2016) 529 [1601.00557].
- [148] R. Valiante, R. Schneider, M. Volonteri and K. Omukai, *From The First Stars To The First Black Holes*, in *Active Galactic Nuclei 12: A Multi-Messenger Perspective (AGN12)*, p. 4, Oct., 2016, DOI.
- [149] J. Regan, E. Visbal, J. H. Wise, Z. Haiman, P. H. Johansson and G. L. Bryan, *Rapid Formation of Massive Black Holes in close proximity to Embryonic Proto-Galaxies*, [1703.03805]. 14

- [150] K. Belczynski et al., *The Effect of Pair-Instability Mass Loss on Black Hole Mergers*, *Astron. Astrophys.* **594** (2016) A97 [1607.03116]. 14
- [151] M. Spera and M. Mapelli, *Very massive stars, pair-instability supernovae and intermediate-mass black holes with the SEVN code*, *Mon. Not. Roy. Astron. Soc.* **470** (2017) 4739 [1706.06109].
- [152] S. E. Woosley, *Pulsational Pair-Instability Supernovae*, *Astrophys. J.* **836** (2017) 244 [1608.08939]. 15
- [153] N. Giacobbo, M. Mapelli and M. Spera, *Merging black hole binaries: the effects of progenitor's metallicity, mass-loss rate and Eddington factor*, *Mon. Not. Roy. Astron. Soc.* **474** (2018) 2959 [1711.03556].
- [154] P. Marchant, M. Renzo, R. Farmer, K. M. W. Pappas, R. E. Taam, S. de Mink et al., *Pulsational pair-instability supernovae in very close binaries*, [1810.13412]. 15
- [155] R. Farmer, M. Renzo, S. de Mink, P. Marchant and S. Justham, *Mind the gap: The location of the lower edge of the pair instability supernovae black hole mass gap*, [1910.12874]. 14, 15
- [156] T. Hartwig, M. Volonteri, V. Bromm, R. S. Klessen, E. Barausse, M. Magg et al., *Gravitational Waves from the Remnants of the First Stars*, *Mon. Not. Roy. Astron. Soc.* **460** (2016) L74 [1603.05655]. 14
- [157] J. M. Ezquiaga and D. E. Holz, *Jumping the gap: searching for LIGO's biggest black holes*, [2006.02211]. 14
- [158] L. Ferrarese and D. Merritt, *A Fundamental relation between supermassive black holes and their host galaxies*, *Astrophys. J. Lett.* **539** (2000) L9 [astro-ph/0006053]. 14
- [159] C. Y. Peng, *How Mergers May Affect The Mass Scaling Relations Between Black Holes, Galaxies, and Other Gravitationally Bound Systems*, *Astrophys. J.* **671** (2007) 1098 [0704.1860].
- [160] M. Volonteri and P. Natarajan, *Journey to the $M_{\text{BH}} - \sigma$ relation: the fate of low mass black holes in the Universe*, *Mon. Not. Roy. Astron. Soc.* **400** (2009) 1911 [0903.2262].
- [161] R. Valiante, R. Schneider, M. Volonteri and K. Omukai, *From the first stars to the first black holes*, *MNRAS* **457** (2016) 3356 [1601.07915].
- [162] J. Silk, *Feedback by Massive Black Holes in Gas-rich Dwarf Galaxies*, *Astrophys. J. Lett.* **839** (2017) L13 [1703.08553]. 14
- [163] S. Vitale and M. Evans, *Parameter estimation for binary black holes with networks of third generation gravitational-wave detectors*, *Phys. Rev.* **D95** (2017) 064052 [1610.06917]. 15, 25
- [164] S. Vitale and C. Whittle, *Characterization of binary black holes by heterogeneous gravitational-wave networks*, *Phys. Rev.* **D98** (2018) 024029 [1804.07866]. 19
- [165] E. D. Hall and M. Evans, *Metrics for next-generation gravitational-wave detectors*, *Class. Quant. Grav.* **36** (2019) 225002 [1902.09485]. 15
- [166] C. Cutler and E. E. Flanagan, *Gravitational waves from merging compact binaries: How accurately can one extract the binary's parameters from the inspiral wave form?*, *Phys. Rev.* **D49** (1994) 2658 [gr-qc/9402014]. 15
- [167] C. K. Mishra, K. G. Arun, B. R. Iyer and B. S. Sathyaprakash, *Parametrized tests of post-Newtonian theory using Advanced LIGO and Einstein Telescope*, *Phys. Rev.* **D82** (2010) 064010 [1005.0304]. 15
- [168] S. Gossan, J. Veitch and B. S. Sathyaprakash, *Bayesian model selection for testing the no-hair theorem with black hole ringdowns*, *Phys. Rev.* **D85** (2012) 124056 [1111.5819]. 26
- [169] S. Bhagwat, D. A. Brown and S. W. Ballmer, *Spectroscopic analysis of stellar mass black-hole mergers in our local universe with ground-based gravitational wave detectors*, *Phys. Rev.* **D94** (2016) 084024 [1607.07845].

- [170] E. Berti, K. Yagi, H. Yang and N. Yunes, *Extreme Gravity Tests with Gravitational Waves from Compact Binary Coalescences: (II) Ringdown*, *Gen. Rel. Grav.* **50** (2018) 49 [1801.03587]. 23, 27
- [171] M. Isi, M. Giesler, W. M. Farr, M. A. Scheel and S. A. Teukolsky, *Testing the no-hair theorem with GW150914*, *Phys. Rev. Lett.* **123** (2019) 111102 [1905.00869].
- [172] G. Carullo, W. Del Pozzo and J. Veitch, *Observational Black Hole Spectroscopy: A time-domain multimode analysis of GW150914*, *Phys. Rev. D* **99** (2019) 123029 [1902.07527]. 15
- [173] S. E. Woosley, S. Blinnikov and A. Heger, *Pulsational pair instability as an explanation for the most luminous supernovae*, *Nature* **450** (2007) 390 [0710.3314]. 15
- [174] F. Ozel, D. Psaltis, R. Narayan and J. E. McClintock, *The Black Hole Mass Distribution in the Galaxy*, *Astrophys. J.* **725** (2010) 1918 [1006.2834]. 15
- [175] W. M. Farr, N. Sravan, A. Cantrell, L. Kreidberg, C. D. Bailyn, I. Mandel et al., *The Mass Distribution of Stellar-Mass Black Holes*, *Astrophys. J.* **741** (2011) 103 [1011.1459].
- [176] L. Kreidberg, C. D. Bailyn, W. M. Farr and V. Kalogera, *Mass Measurements of Black Holes in X-Ray Transients: Is There a Mass Gap?*, *Astrophys. J.* **757** (2012) 36 [1205.1805].
- [177] L. Wyrzykowski and I. Mandel, *Constraining the masses of microlensing black holes and the mass gap with Gaia DR2*, *Astron. Astrophys.* **636** (2020) A20 [1904.07789]. 15
- [178] K. Belczynski, G. Wiktorowicz, C. Fryer, D. Holz and V. Kalogera, *Missing Black Holes Unveil The Supernova Explosion Mechanism*, *Astrophys. J.* **757** (2012) 91 [1110.1635]. 15
- [179] C. L. Fryer, K. Belczynski, G. Wiktorowicz, M. Dominik, V. Kalogera and D. E. Holz, *Compact Remnant Mass Function: Dependence on the Explosion Mechanism and Metallicity*, *Astrophys. J.* **749** (2012) 91 [1110.1726].
- [180] C. L. Fryer, S. Andrews, W. Even, A. Heger and S. Safi-Harb, *Parameterizing the Supernova Engine and Its Effect on Remnants and Basic Yields*, *Astrophys. J.* **856** (2018) 63 [1712.03415].
- [181] M. Zevin, M. Spera, C. P. Berry and V. Kalogera, *Exploring the Lower Mass Gap and Unequal Mass Regime in Compact Binary Evolution*, *Astrophys. J. Lett.* **899** (2020) L1 [2006.14573]. 15
- [182] B. Kiziltan, A. Kottas, M. De Yoreo and S. E. Thorsett, *The Neutron Star Mass Distribution*, *Astrophys. J.* **778** (2013) 66 [1309.6635]. 15
- [183] J. Alsing, H. O. Silva and E. Berti, *Evidence for a maximum mass cut-off in the neutron star mass distribution and constraints on the equation of state*, *Mon. Not. Roy. Astron. Soc.* **478** (2018) 1377 [1709.07889].
- [184] E. R. Most, L. J. Papenfort, L. R. Weih and L. Rezzolla, *A lower bound on the maximum mass if the secondary in GW190814 was once a rapidly spinning neutron star*, [2006.14601].
- [185] F. J. Fattoyev, C. J. Horowitz, J. Piekarewicz and B. Reed, *GW190814: Impact of a 2.6 solar mass neutron star on nucleonic equations of state*, [2007.03799].
- [186] A. Tsokaros, M. Ruiz and S. L. Shapiro, *GW190814: Spin and equation of state of a neutron star companion*, [2007.05526]. 15
- [187] R. Farmer, M. Renzo, S. de Mink, M. Fishbach and S. Justham, *Constraints from gravitational wave detections of binary black hole mergers on the $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$ rate*, [2006.06678]. 15
- [188] T. A. Apostolatos, C. Cutler, G. J. Sussman and K. S. Thorne, *Spin induced orbital precession and its modulation of the gravitational wave forms from merging binaries*, *Phys. Rev.* **D49** (1994) 6274. 16
- [189] L. Blanchet, *Gravitational Radiation from Post-Newtonian Sources and Inspiralling Compact Binaries*, *Living Rev. Rel.* **17** (2014) 2 [1310.1528]. 16

- [190] P. C. Peters, *Gravitational Radiation and the Motion of Two Point Masses*, *Phys. Rev.* **136** (1964) B1224. 16
- [191] M. E. Lower, E. Thrane, P. D. Lasky and R. Smith, *Measuring eccentricity in binary black hole inspirals with gravitational waves*, *Phys. Rev.* **D98** (2018) 083028 [1806.05350]. 16
- [192] I. M. Romero-Shaw, P. D. Lasky and E. Thrane, *Searching for Eccentricity: Signatures of Dynamical Formation in the First Gravitational-Wave Transient Catalogue of LIGO and Virgo*, *Mon. Not. Roy. Astron. Soc.* **490** (2019) 5210 [1909.05466].
- [193] A. K. Lenon, A. H. Nitz and D. A. Brown, *Measuring the eccentricity of GW170817 and GW190425*, [2005.14146]. 16
- [194] K. Kyutoku and N. Seto, *Concise estimate of the expected number of detections for stellar-mass binary black holes by eLISA*, *MNRAS* **462** (2016) 2177 [1606.02298]. 16
- [195] A. Nishizawa, E. Berti, A. Klein and A. Sesana, *eLISA eccentricity measurements as tracers of binary black hole formation*, *Phys. Rev. D* **94** (2016) 064020 [1605.01341]. 16
- [196] K. W. K. Wong, E. D. Kovetz, C. Cutler and E. Berti, *Expanding the LISA Horizon from the Ground*, *Phys. Rev. Lett.* **121** (2018) 251102 [1808.08247]. 16
- [197] K. Breivik, C. L. Rodriguez, S. L. Larson, V. Kalogera and F. A. Rasio, *Distinguishing between Formation Channels for Binary Black Holes with LISA*, *ApJ* **830** (2016) L18 [1606.09558]. 16
- [198] A. Nishizawa, A. Sesana, E. Berti and A. Klein, *Constraining stellar binary black hole formation scenarios with eLISA eccentricity measurements*, *MNRAS* **465** (2017) 4375 [1606.09295].
- [199] J. Samsing, D. J. D’Orazio, A. Askar and M. Giersz, *Black Hole Mergers from Globular Clusters Observable by LISA and LIGO: Results from post-Newtonian Binary-Single Scatterings*, [1802.08654]. 16
- [200] S. Vitale, *Multiband Gravitational-Wave Astronomy: Parameter Estimation and Tests of General Relativity with Space- and Ground-Based Detectors*, *Phys. Rev. Lett.* **117** (2016) 051102 [1605.01037]. 16
- [201] Z. Carson and K. Yagi, *Multi-band gravitational wave tests of general relativity*, *Class. Quant. Grav.* **37** (2020) 02LT01 [1905.13155].
- [202] G. Gnocchi, A. Maselli, T. Abdelsalhin, N. Giacobbo and M. Mapelli, *Bounding alternative theories of gravity with multiband GW observations*, *Phys. Rev. D* **100** (2019) 064024 [1905.13460].
- [203] C. Liu, L. Shao, J. Zhao and Y. Gao, *Multiband Observation of LIGO/Virgo Binary Black Hole Mergers in the Gravitational-wave Transient Catalog GWTC-1*, *Mon. Not. Roy. Astron. Soc.* **496** (2020) 182 [2004.12096]. 16
- [204] E. Barausse, N. Yunes and K. Chamberlain, *Theory-Agnostic Constraints on Black-Hole Dipole Radiation with Multiband Gravitational-Wave Astrophysics*, *Phys. Rev. Lett.* **116** (2016) 241104 [1603.04075]. 16
- [205] R. Valiante, M. Colpi, R. Schneider, A. Mangiagli, M. Bonetti, G. Cerini et al., *Unveiling early black hole growth with multifrequency gravitational wave observations*, *Mon. Not. Roy. Astron. Soc.* **500** (2020) 4095 [2010.15096]. 16, 17
- [206] P. Amaro-Seoane and L. Santamaría, *Detection of IMBHs with Ground-based Gravitational Wave Observatories: A Biography of a Binary of Black Holes, from Birth to Death*, *ApJ* **722** (2010) 1197 [0910.0254]. 17
- [207] P. Amaro-Seoane, *Detecting Intermediate-Mass Ratio Inspirals From The Ground And Space*, *Phys. Rev.* **D98** (2018) 063018 [1807.03824]. 17
- [208] B. F. Schutz, *Determining the Hubble Constant from Gravitational Wave Observations*, *Nature* **323** (1986) 310. 18, 19
- [209] LIGO SCIENTIFIC, VIRGO collaboration, *GW170817: Implications for the Stochastic Gravitational-Wave Background from Compact Binary Coalescences*, *Phys. Rev. Lett.* **120** (2018) 091101 [1710.05837]. 19, 22

- [210] L. P. Grishchuk, *Amplification of gravitational waves in an isotropic universe*, *Sov. Phys. JETP* **40** (1975) 409. 19, 20
- [211] A. A. Starobinsky, *Spectrum of relict gravitational radiation and the early state of the universe*, *JETP Lett.* **30** (1979) 682.
- [212] V. A. Rubakov, M. V. Sazhin and A. V. Veryaskin, *Graviton Creation in the Inflationary Universe and the Grand Unification Scale*, *Phys. Lett.* **115B** (1982) 189. 20
- [213] M. Giovannini, *Gravitational waves constraints on postinflationary phases stiffer than radiation*, *Phys. Rev.* **D58** (1998) 083504 [hep-ph/9806329]. 20
- [214] L. A. Boyle and A. Buonanno, *Relating gravitational wave constraints from primordial nucleosynthesis, pulsar timing, laser interferometers, and the CMB: Implications for the early Universe*, *Phys. Rev.* **D78** (2008) 043531 [0708.2279].
- [215] D. G. Figueroa and C. T. Byrnes, *The Standard Model Higgs as the origin of the hot Big Bang*, *Phys. Lett.* **B767** (2017) 272 [1604.03905]. 20
- [216] C. Caprini and D. G. Figueroa, *Cosmological Backgrounds of Gravitational Waves*, *Class. Quant. Grav.* **35** (2018) 163001 [1801.04268].
- [217] A. Kosowsky, M. S. Turner and R. Watkins, *Gravitational radiation from colliding vacuum bubbles*, *Phys. Rev.* **D45** (1992) 4514. 21
- [218] M. Kamionkowski, A. Kosowsky and M. S. Turner, *Gravitational radiation from first order phase transitions*, *Phys. Rev.* **D49** (1994) 2837 [astro-ph/9310044].
- [219] M. Hindmarsh, S. J. Huber, K. Rummukainen and D. J. Weir, *Shape of the acoustic gravitational wave power spectrum from a first order phase transition*, *Phys. Rev.* **D96** (2017) 103520 [1704.05871]. 21
- [220] C. Caprini, R. Durrer and G. Servant, *The stochastic gravitational wave background from turbulence and magnetic fields generated by a first-order phase transition*, *JCAP* **0912** (2009) 024 [0909.0622]. 19, 21
- [221] S. Nissanke, D. E. Holz, S. A. Hughes, N. Dalal and J. L. Sievers, *Exploring Short Gamma-ray Bursts as Gravitational-wave Standard Sirens*, *ApJ* **725** (2010) 496 [0904.1017]. 19
- [222] H.-Y. Chen, M. Fishbach and D. E. Holz, *A two per cent Hubble constant measurement from standard sirens within five years*, *Nature* **562** (2018) 545 [1712.06531].
- [223] S. M. Feeney, H. V. Peiris, A. R. Williamson, S. M. Nissanke, D. J. Mortlock, J. Alsing et al., *Prospects for Resolving the Hubble Constant Tension with Standard Sirens*, *Phys. Rev. Lett.* **122** (2019) 061105 [1802.03404]. 19
- [224] S. Vitale and H.-Y. Chen, *Measuring the Hubble constant with neutron star black hole mergers*, *Phys. Rev. Lett.* **121** (2018) 021303 [1804.07337]. 19
- [225] P. B. Graff, A. Buonanno and B. S. Sathyaprakash, *Missing Link: Bayesian detection and measurement of intermediate-mass black-hole binaries*, *Phys. Rev. D* **92** (2015) 022002 [1504.04766]. 19
- [226] S. Borhanian, A. Dhani, A. Gupta, K. Arun and B. Sathyaprakash, *Dark Sirens to Resolve the Hubble-Lemaître Tension*, [2007.02883]. 19
- [227] K. Hotokezaka, E. Nakar, O. Gottlieb, S. Nissanke, K. Masuda, G. Hallinan et al., *A Hubble constant measurement from superluminal motion of the jet in GW170817*, *Nat. Astron.* **3** (2019) 940 [1806.10596]. 19
- [228] W. Del Pozzo, *Inference of the cosmological parameters from gravitational waves: application to second generation interferometers*, *Phys. Rev.* **D86** (2012) 043011 [1108.1317]. 19

- [229] M. Soares-Santos, D. E. Holz, J. Annis, R. Chornock, K. Herner, E. Berger et al., *The Electromagnetic Counterpart of the Binary Neutron Star Merger LIGO/Virgo GW170817. I. Discovery of the Optical Counterpart Using the Dark Energy Camera*, *ApJ* **848** (2017) L16 [1710.05459]. 19
- [230] N. Dalal, D. E. Holz, S. A. Hughes and B. Jain, *Short grb and binary black hole standard sirens as a probe of dark energy*, *Phys. Rev.* **D74** (2006) 063006 [astro-ph/0601275]. 19
- [231] S. Nissanke, D. E. Holz, S. A. Hughes, N. Dalal and J. L. Sievers, *Exploring short gamma-ray bursts as gravitational-wave standard sirens*, *Astrophys. J.* **725** (2010) 496 [0904.1017].
- [232] B. S. Sathyaprakash, B. F. Schutz and C. Van Den Broeck, *Cosmography with the Einstein Telescope*, *Class. Quant. Grav.* **27** (2010) 215006 [0906.4151].
- [233] W. Zhao, C. Van Den Broeck, D. Baskaran and T. G. F. Li, *Determination of Dark Energy by the Einstein Telescope: Comparing with CMB, BAO and SNIa Observations*, *Phys. Rev.* **D83** (2011) 023005 [1009.0206].
- [234] S. R. Taylor and J. R. Gair, *Cosmology with the lights off: standard sirens in the Einstein Telescope era*, *Phys. Rev.* **D86** (2012) 023502 [1204.6739].
- [235] S. Camera and A. Nishizawa, *Beyond Concordance Cosmology with Magnification of Gravitational-Wave Standard Sirens*, *Phys. Rev. Lett.* **110** (2013) 151103 [1303.5446].
- [236] R.-G. Cai and T. Yang, *Estimating cosmological parameters by the simulated data of gravitational waves from the Einstein Telescope*, *Phys. Rev.* **D95** (2017) 044024 [1608.08008].
- [237] E. Belgacem, Y. Dirian, S. Foffa and M. Maggiore, *Gravitational-wave luminosity distance in modified gravity theories*, *Phys. Rev.* **D97** (2018) 104066 [1712.08108]. 19
- [238] E. Belgacem, Y. Dirian, S. Foffa and M. Maggiore, *Modified gravitational-wave propagation and standard sirens*, *Phys. Rev.* **D98** (2018) 023510 [1805.08731]. 19, 20
- [239] C. Deffayet and K. Menou, *Probing Gravity with Spacetime Sirens*, *Astrophys. J. Lett.* **668** (2007) L143 [0709.0003]. 19
- [240] I. D. Saltas, I. Sawicki, L. Amendola and M. Kunz, *Anisotropic Stress as a Signature of Nonstandard Propagation of Gravitational Waves*, *Phys. Rev. Lett.* **113** (2014) 191101 [1406.7139]. 19
- [241] L. Lombriser and A. Taylor, *Breaking a Dark Degeneracy with Gravitational Waves*, *JCAP* **1603** (2016) 031 [1509.08458].
- [242] S. Arai and A. Nishizawa, *Generalized framework for testing gravity with gravitational-wave propagation. II. Constraints on Horndeski theory*, *Phys. Rev.* **D97** (2018) 104038 [1711.03776].
- [243] L. Amendola, I. Sawicki, M. Kunz and I. D. Saltas, *Direct detection of gravitational waves can measure the time variation of the Planck mass*, *JCAP* **1808** (2018) 030 [1712.08623].
- [244] E. V. Linder, *No Slip Gravity*, *JCAP* **1803** (2018) 005 [1801.01503]. 19
- [245] M. Maggiore, *Phantom dark energy from nonlocal infrared modifications of general relativity*, *Phys. Rev.* **D89** (2014) 043008 [1307.3898]. 19
- [246] M. Maggiore and M. Mancarella, *Nonlocal gravity and dark energy*, *Phys. Rev.* **D90** (2014) 023005 [1402.0448].
- [247] E. Belgacem, Y. Dirian, S. Foffa and M. Maggiore, *Nonlocal gravity. Conceptual aspects and cosmological predictions*, *JCAP* **1803** (2018) 002 [1712.07066]. 19
- [248] S. Sachdev, T. Regimbau and B. Sathyaprakash, *Subtracting compact binary foreground sources to reveal primordial gravitational-wave backgrounds*, *Phys. Rev. D* **102** (2020) 024051 [2002.05365]. 20, 22

- [249] LIGO SCIENTIFIC, VIRGO collaboration, *Upper Limits on the Stochastic Gravitational-Wave Background from Advanced LIGO's First Observing Run*, *Phys. Rev. Lett.* **118** (2017) 121101 [1612.02029]. 21
- [250] P. D. Lasky et al., *Gravitational-wave cosmology across 29 decades in frequency*, *Phys. Rev.* **X6** (2016) 011035 [1511.05994]. 21
- [251] E. Thrane and J. D. Romano, *Sensitivity curves for searches for gravitational-wave backgrounds*, *Phys. Rev.* **D88** (2013) 124032 [1310.5300]. 21
- [252] LIGO SCIENTIFIC, VIRGO collaboration, *GW150914: Implications for the stochastic gravitational wave background from binary black holes*, *Phys. Rev. Lett.* **116** (2016) 131102 [1602.03847]. 21
- [253] "Lisa mission proposal." Available at https://www.elisascience.org/files/publications/LISA_L3_20170120.pdf (2017). 21
- [254] M. Fitz Axen, S. Banagiri, A. Matas, C. Caprini and V. Mandic, *Multiwavelength observations of cosmological phase transitions using LISA and Cosmic Explorer*, *Phys. Rev.* **D98** (2018) 103508 [1806.02500]. 21
- [255] B. P. Abbott, R. Abbott, T. D. Abbott, F. Acernese, K. Ackley, C. Adams et al., *GW170817: Implications for the Stochastic Gravitational-Wave Background from Compact Binary Coalescences*, *Phys. Rev. Lett.* **120** (2018) 091101 [1710.05837]. 21
- [256] N. Barnaby and M. Peloso, *Large Nongaussianity in Axion Inflation*, *Phys. Rev. Lett.* **106** (2011) 181301 [1011.1500]. 20
- [257] J. L. Cook and L. Sorbo, *Particle production during inflation and gravitational waves detectable by ground-based interferometers*, *Phys. Rev.* **D85** (2012) 023534 [1109.0022]. 20
- [258] J. L. Cook and L. Sorbo, *An inflationary model with small scalar and large tensor nongaussianities*, *JCAP* **1311** (2013) 047 [1307.7077]. 20
- [259] M. Giovannini, *Production and detection of relic gravitons in quintessential inflationary models*, *Phys. Rev.* **D60** (1999) 123511 [astro-ph/9903004]. 20
- [260] D. J. Weir, *Revisiting the envelope approximation: gravitational waves from bubble collisions*, *Phys. Rev.* **D93** (2016) 124037 [1604.08429]. 21
- [261] T. Kahniashvili, L. Campanelli, G. Gogoberidze, Y. Maravin and B. Ratra, *Gravitational Radiation from Primordial Helical Inverse Cascade MHD Turbulence*, *Phys. Rev.* **D78** (2008) 123006 [0809.1899]. 21
- [262] R. Jeannerot, J. Rocher and M. Sakellariadou, *How generic is cosmic string formation in SUSY GUTs*, *Phys. Rev.* **D68** (2003) 103514 [hep-ph/0308134]. 21, 26
- [263] T. Vachaspati and A. Vilenkin, *Gravitational Radiation from Cosmic Strings*, *Phys. Rev.* **D31** (1985) 3052. 21
- [264] M. Sakellariadou, *Gravitational waves emitted from infinite strings*, *Phys. Rev.* **D42** (1990) 354. 21
- [265] J. J. Blanco-Pillado and K. D. Olum, *Stochastic gravitational wave background from smoothed cosmic string loops*, *Phys. Rev.* **D96** (2017) 104046 [1709.02693]. 21
- [266] C. Ringeval and T. Suyama, *Stochastic gravitational waves from cosmic string loops in scaling*, *JCAP* **1712** (2017) 027 [1709.03845].
- [267] LIGO SCIENTIFIC, VIRGO collaboration, *Constraints on cosmic strings using data from the first Advanced LIGO observing run*, *Phys. Rev.* **D97** (2018) 102002 [1712.01168].
- [268] A. C. Jenkins and M. Sakellariadou, *Anisotropies in the stochastic gravitational-wave background: Formalism and the cosmic string case*, *Phys. Rev. D* **98** (2018) 063509 [1802.06046]. 21

- [269] A. Pierce, K. Riles and Y. Zhao, *Searching for Dark Photon Dark Matter with Gravitational Wave Detectors*, *Phys. Rev. Lett.* **121** (2018) 061102 [1801.10161]. 21
- [270] T. Regimbau, M. Evans, N. Christensen, E. Katsavounidis, B. Sathyaprakash and S. Vitale, *Digging deeper: Observing primordial gravitational waves below the binary black hole produced stochastic background*, *Phys. Rev. Lett.* **118** (2017) 151105 [1611.08943]. 22
- [271] V. Mandic, S. Bird and I. Cholis, *Stochastic Gravitational-Wave Background due to Primordial Binary Black Hole Mergers*, *Phys. Rev. Lett.* **117** (2016) 201102 [1608.06699].
- [272] I. Dvorkin, J.-P. Uzan, E. Vangioni and J. Silk, *Synthetic model of the gravitational wave background from evolving binary compact objects*, *Phys. Rev.* **D94** (2016) 103011 [1607.06818].
- [273] K. Nakazato, Y. Niino and N. Sago, *Gravitational-Wave Background from Binary Mergers and Metallicity Evolution of Galaxies*, *Astrophys. J.* **832** (2016) 146 [1605.02146].
- [274] I. Dvorkin, E. Vangioni, J. Silk, J.-P. Uzan and K. A. Olive, *Metallicity-constrained merger rates of binary black holes and the stochastic gravitational wave background*, *Mon. Not. Roy. Astron. Soc.* **461** (2016) 3877 [1604.04288].
- [275] E. F. D. Evangelista and J. C. N. Araujo, *The Gravitational Wave Background from Coalescing Compact Binaries: A New Method*, *Braz. J. Phys.* **44** (2014) 824 [1504.06605]. 22
- [276] M. Abernathy, F. Acernese, P. Ajith, B. Allen, P. Amaro-Seoane, N. Andersson et al., *Einstein Gravitational Wave Telescope Conceptual Design Study*. 2011. 22
- [277] C. Cutler and J. Harms, *BBO and the neutron-star-binary subtraction problem*, *Phys. Rev.* **D73** (2006) 042001 [gr-qc/0511092]. 22
- [278] G. Cusin, I. Dvorkin, C. Pitrou and J.-P. Uzan, *First predictions of the angular power spectrum of the astrophysical gravitational wave background*, *Phys. Rev. Lett.* **120** (2018) 231101 [1803.03236]. 22
- [279] A. C. Jenkins, M. Sakellariadou, T. Regimbau and E. Slezak, *Anisotropies in the astrophysical gravitational-wave background: Predictions for the detection of compact binaries by LIGO and Virgo*, *Phys. Rev.* **D98** (2018) 063501 [1806.01718]. 22
- [280] R. Smith and E. Thrane, *Optimal Search for an Astrophysical Gravitational-Wave Background*, *Phys. Rev.* **X8** (2018) 021019 [1712.00688]. 22
- [281] M. Surace, K. D. Kokkotas and P. Pnigouras, *The stochastic background of gravitational waves due to the f-mode instability in neutron stars*, *Astron. Astrophys.* **586** (2016) A86 [1512.02502]. 22
- [282] D. Talukder, E. Thrane, S. Bose and T. Regimbau, *Measuring neutron-star ellipticity with measurements of the stochastic gravitational-wave background*, *Phys. Rev.* **D89** (2014) 123008 [1404.4025].
- [283] P. D. Lasky, M. F. Bennett and A. Melatos, *Stochastic gravitational wave background from hydrodynamic turbulence in differentially rotating neutron stars*, *Phys. Rev.* **D87** (2013) 063004 [1302.6033]. 22
- [284] K. Crocker, T. Prestegard, V. Mandic, T. Regimbau, K. Olive and E. Vangioni, *Systematic study of the stochastic gravitational-wave background due to stellar core collapse*, *Phys. Rev.* **D95** (2017) 063015 [1701.02638]. 22
- [285] K. Crocker, V. Mandic, T. Regimbau, K. Belczynski, W. Gladysz, K. Olive et al., *Model of the stochastic gravitational-wave background due to core collapse to black holes*, *Phys. Rev.* **D92** (2015) 063005 [1506.02631]. 22
- [286] I. Kowalska, T. Bulik and K. Belczynski, *Gravitational wave background from population III binaries*, *Astron. Astrophys.* **541** (2012) A120 [1202.3346]. 22

- [287] C. M. Will, *The Confrontation between general relativity and experiment*, *Living Rev. Rel.* **9** (2006) 3 [gr-qc/0510072]. 23
- [288] D. Psaltis, *Probes and Tests of Strong-Field Gravity with Observations in the Electromagnetic Spectrum*, *Living Rev. Rel.* **11** (2008) 9 [0806.1531].
- [289] N. Yunes and X. Siemens, *Gravitational-Wave Tests of General Relativity with Ground-Based Detectors and Pulsar Timing-Arrays*, *Living Rev. Rel.* **16** (2013) 9 [1304.3473]. 23
- [290] C. Rham, *Introduction to Massive Gravity*, *Lect. Notes Phys.* **892** (2015) 139. 23
- [291] G. Chapline, E. Hohlfeld, R. B. Laughlin and D. I. Santiago, *Quantum phase transitions and the breakdown of classical general relativity*, *Int. J. Mod. Phys.* **A18** (2003) 3587 [gr-qc/0012094].
- [292] T. Baker, D. Psaltis and C. Skordis, *Linking Tests of Gravity On All Scales: from the Strong-Field Regime to Cosmology*, *Astrophys. J.* **802** (2015) 63 [1412.3455].
- [293] M. Ishak, *Testing General Relativity in Cosmology*, *Living Rev. Rel.* **22** (2019) 1 [1806.10122]. 23
- [294] E. Bianchi and R. C. Myers, *On the Architecture of Spacetime Geometry*, *Class. Quant. Grav.* **31** (2014) 214002 [1212.5183]. 23, 27
- [295] S. W. Hawking, M. J. Perry and A. Strominger, *Soft Hair on Black Holes*, *Phys. Rev. Lett.* **116** (2016) 231301 [1601.00921]. 23
- [296] LIGO SCIENTIFIC, VIRGO collaboration, *Observation of Gravitational Waves from a Binary Black Hole Merger*, *Phys. Rev. Lett.* **116** (2016) 061102 [1602.03837]. 23
- [297] LIGO SCIENTIFIC, VIRGO collaboration, *GWTC-1: A Gravitational-Wave Transient Catalog of Compact Binary Mergers Observed by LIGO and Virgo during the First and Second Observing Runs*, *Phys. Rev.* **X9** (2019) 031040 [1811.12907]. 23
- [298] K. S. Thorne, *Gravitational radiation*, in *Three hundred years of gravitation*, S. W. Hawking and W. Israel, eds., (Cambridge), pp. 330–458, Cambridge University Press, (1987). 23
- [299] B. S. Sathyaprakash and B. F. Schutz, *Physics, Astrophysics and Cosmology with Gravitational Waves*, *Living Rev. Rel.* **12** (2009) 2 [0903.0338]. 23
- [300] C. M. Will, *The Confrontation between General Relativity and Experiment*, *Living Rev. Rel.* **17** (2014) 4 [1403.7377]. 24, 25
- [301] E. Berti et al., *Testing General Relativity with Present and Future Astrophysical Observations*, *Class. Quant. Grav.* **32** (2015) 243001 [1501.07274]. 24
- [302] J. F. Donoghue, *Introduction to the effective field theory description of gravity*, in *Advanced School on Effective Theories Almunecar, Spain, June 25-July 1, 1995*, 1995, [gr-qc/9512024]. 24
- [303] C. P. Burgess, *Introduction to Effective Field Theory*, *Ann. Rev. Nucl. Part. Sci.* **57** (2007) 329 [hep-th/0701053].
- [304] S. Weinberg, *Effective Field Theory for Inflation*, *Phys. Rev.* **D77** (2008) 123541 [0804.4291]. 25
- [305] W. D. Goldberger and I. Z. Rothstein, *An Effective field theory of gravity for extended objects*, *Phys. Rev.* **D73** (2006) 104029 [hep-th/0409156]. 24
- [306] N. Sennett, L. Shao and J. Steinhoff, *Effective action model of dynamically scalarizing binary neutron stars*, *Phys. Rev.* **D96** (2017) 084019 [1708.08285]. 24
- [307] S. Endlich, V. Gorbenko, J. Huang and L. Senatore, *An effective formalism for testing extensions to General Relativity with gravitational waves*, *JHEP* **09** (2017) 122 [1704.01590].

- [308] V. Cardoso, M. Kimura, A. Maselli and L. Senatore, *Black Holes in an Effective Field Theory Extension of General Relativity*, *Phys. Rev. Lett.* **121** (2018) 251105 [1808.08962]. 24
- [309] D. Lovelock, *The Einstein tensor and its generalizations*, *J. Math. Phys.* **12** (1971) 498. 24
- [310] C. Brans and R. H. Dicke, *Mach's principle and a relativistic theory of gravitation*, *Phys. Rev.* **124** (1961) 925. 24
- [311] Y. Fujii and K. Maeda, *The scalar-tensor theory of gravitation*, Cambridge Monographs on Mathematical Physics. Cambridge University Press, 2007, 10.1017/CBO9780511535093. 24
- [312] C. Palenzuela, E. Barausse, M. Ponce and L. Lehner, *Dynamical scalarization of neutron stars in scalar-tensor gravity theories*, *Phys. Rev.* **D89** (2014) 044024 [1310.4481]. 25
- [313] M. Shibata, K. Taniguchi, H. Okawa and A. Buonanno, *Coalescence of binary neutron stars in a scalar-tensor theory of gravity*, *Phys. Rev.* **D89** (2014) 084005 [1310.0627]. 25
- [314] T. Damour and G. Esposito-Farese, *Nonperturbative strong field effects in tensor - scalar theories of gravitation*, *Phys. Rev. Lett.* **70** (1993) 2220. 25
- [315] P. Kanti, N. E. Mavromatos, J. Rizos, K. Tamvakis and E. Winstanley, *Dilatonic black holes in higher curvature string gravity*, *Phys. Rev.* **D54** (1996) 5049 [hep-th/9511071].
- [316] S. Mignemi and N. R. Stewart, *Charged black holes in effective string theory*, *Phys. Rev.* **D47** (1993) 5259 [hep-th/9212146].
- [317] G. Antoniou, A. Bakopoulos and P. Kanti, *Evasion of No-Hair Theorems and Novel Black-Hole Solutions in Gauss-Bonnet Theories*, *Phys. Rev. Lett.* **120** (2018) 131102 [1711.03390]. 25
- [318] LIGO SCIENTIFIC, VIRGO collaboration, *Tests of General Relativity with GW170817*, *Phys. Rev. Lett.* **123** (2019) 011102 [1811.00364]. 25
- [319] C. de Rham, G. Gabadadze and A. J. Tolley, *Resummation of Massive Gravity*, *Phys. Rev. Lett.* **106** (2011) 231101 [1011.1232]. 25
- [320] S. F. Hassan and R. A. Rosen, *Resolving the Ghost Problem in non-Linear Massive Gravity*, *Phys. Rev. Lett.* **108** (2012) 041101 [1106.3344]. 25
- [321] D. Mattingly, *Modern tests of Lorentz invariance*, *Living Rev. Rel.* **8** (2005) 5 [gr-qc/0502097]. 25
- [322] P. Horava, *Quantum Gravity at a Lifshitz Point*, *Phys. Rev.* **D79** (2009) 084008 [0901.3775]. 25
- [323] T. Jacobson and D. Mattingly, *Gravity with a dynamical preferred frame*, *Phys. Rev.* **D64** (2001) 024028 [gr-qc/0007031]. 25
- [324] C. Eling and T. Jacobson, *Black Holes in Einstein-Aether Theory*, *Class. Quant. Grav.* **23** (2006) 5643 [gr-qc/0604088]. 25
- [325] E. Barausse, T. Jacobson and T. P. Sotiriou, *Black holes in Einstein-aether and Horava-Lifshitz gravity*, *Phys. Rev.* **D83** (2011) 124043 [1104.2889]. 25
- [326] T. P. Sotiriou, *Detecting Lorentz Violations with Gravitational Waves from Black Hole Binaries*, *Phys. Rev. Lett.* **120** (2018) 041104 [1709.00940]. 25
- [327] V. A. Kostelecký and M. Mewes, *Testing local Lorentz invariance with gravitational waves*, *Phys. Lett.* **B757** (2016) 510 [1602.04782]. 25
- [328] M. B. Green, J. H. Schwarz and E. Witten, *SUPERSTRING THEORY. VOL. 2: LOOP AMPLITUDES, ANOMALIES AND PHENOMENOLOGY*. Cambridge, Uk: Univ. Pr. (1987) 596 P. (Cambridge Monographs On Mathematical Physics), 1988. 25

- [329] A. Ashtekar, A. P. Balachandran and S. Jo, *The CP Problem in Quantum Gravity*, *Int. J. Mod. Phys. A* **4** (1989) 1493. 25
- [330] R. Jackiw and S. Y. Pi, *Chern-Simons modification of general relativity*, *Phys. Rev.* **D68** (2003) 104012 [gr-qc/0308071]. 25
- [331] N. Yunes and F. Pretorius, *Dynamical Chern-Simons Modified Gravity. I. Spinning Black Holes in the Slow-Rotation Approximation*, *Phys. Rev.* **D79** (2009) 084043 [0902.4669]. 25
- [332] C. F. Sopuerta and N. Yunes, *Extreme and Intermediate-Mass Ratio Inspirals in Dynamical Chern-Simons Modified Gravity*, *Phys. Rev.* **D80** (2009) 064006 [0904.4501]. 25
- [333] K. Yagi, N. Yunes and T. Tanaka, *Gravitational Waves from Quasi-Circular Black Hole Binaries in Dynamical Chern-Simons Gravity*, *Phys. Rev. Lett.* **109** (2012) 251105 [1208.5102].
- [334] M. Okounkova, L. C. Stein, M. A. Scheel and D. A. Hemberger, *Numerical binary black hole mergers in dynamical Chern-Simons gravity: Scalar field*, *Phys. Rev.* **D96** (2017) 044020 [1705.07924]. 25
- [335] K. Yagi and H. Yang, *Probing Gravitational Parity Violation with Gravitational Waves from Stellar-mass Black Hole Binaries*, *Phys. Rev.* **D97** (2018) 104018 [1712.00682]. 25
- [336] R. Essig et al., *Working Group Report: New Light Weakly Coupled Particles*, in *Proceedings, 2013 Community Summer Study on the Future of U.S. Particle Physics: Snowmass on the Mississippi (CSS2013): Minneapolis, MN, USA, July 29-August 6, 2013*, 2013, [1311.0029], <http://www.slac.stanford.edu/econf/C1307292/docs/IntensityFrontier/NewLight-17.pdf>. 25, 27
- [337] A. Arvanitaki, S. Dimopoulos, S. Dubovsky, N. Kaloper and J. March-Russell, *String Axiverse*, *Phys. Rev.* **D81** (2010) 123530 [0905.4720]. 25
- [338] P. Pani, V. Cardoso, L. Gualtieri, E. Berti and A. Ishibashi, *Black hole bombs and photon mass bounds*, *Phys. Rev. Lett.* **109** (2012) 131102 [1209.0465].
- [339] R. Brito, V. Cardoso and P. Pani, *Massive spin-2 fields on black hole spacetimes: Instability of the Schwarzschild and Kerr solutions and bounds on the graviton mass*, *Phys. Rev.* **D88** (2013) 023514 [1304.6725]. 25
- [340] A. Arvanitaki and S. Dubovsky, *Exploring the String Axiverse with Precision Black Hole Physics*, *Phys. Rev.* **D83** (2011) 044026 [1004.3558]. 25
- [341] R. Brito, V. Cardoso and P. Pani, *Black holes as particle detectors: evolution of superradiant instabilities*, *Class. Quant. Grav.* **32** (2015) 134001 [1411.0686].
- [342] W. E. East and F. Pretorius, *Superradiant Instability and Backreaction of Massive Vector Fields around Kerr Black Holes*, *Phys. Rev. Lett.* **119** (2017) 041101 [1704.04791]. 25
- [343] A. Arvanitaki, M. Baryakhtar and X. Huang, *Discovering the QCD Axion with Black Holes and Gravitational Waves*, *Phys. Rev.* **D91** (2015) 084011 [1411.2263]. 25
- [344] A. Arvanitaki, M. Baryakhtar, S. Dimopoulos, S. Dubovsky and R. Lasenby, *Black Hole Mergers and the QCD Axion at Advanced LIGO*, *Phys. Rev.* **D95** (2017) 043001 [1604.03958]. 25
- [345] M. Baryakhtar, R. Lasenby and M. Teo, *Black Hole Superradiance Signatures of Ultralight Vectors*, *Phys. Rev.* **D96** (2017) 035019 [1704.05081]. 25
- [346] R. Brito, S. Ghosh, E. Barausse, E. Berti, V. Cardoso, I. Dvorkin et al., *Gravitational wave searches for ultralight bosons with LIGO and LISA*, *Phys. Rev.* **D96** (2017) 064050 [1706.06311]. 25
- [347] R. Brito, S. Ghosh, E. Barausse, E. Berti, V. Cardoso, I. Dvorkin et al., *Stochastic and resolvable gravitational waves from ultralight bosons*, *Phys. Rev. Lett.* **119** (2017) 131101 [1706.05097]. 25

- [348] H. Yoshino and H. Kodama, *Gravitational radiation from an axion cloud around a black hole: Superradiant phase*, *PTEP* **2014** (2014) 043E02 [1312.2326]. 25
- [349] W. E. East, *Superradiant instability of massive vector fields around spinning black holes in the relativistic regime*, *Phys. Rev.* **D96** (2017) 024004 [1705.01544]. 25
- [350] D. Baumann, H. S. Chia and R. A. Porto, *Probing Ultralight Bosons with Binary Black Holes*, *Phys. Rev.* **D99** (2019) 044001 [1804.03208]. 25
- [351] L. Hui, J. P. Ostriker, S. Tremaine and E. Witten, *Ultralight scalars as cosmological dark matter*, *Phys. Rev.* **D95** (2017) 043541 [1610.08297]. 25, 27
- [352] S. D. Mathur, *The Fuzzball proposal for black holes: An Elementary review*, *Fortsch. Phys.* **53** (2005) 793 [hep-th/0502050]. 26, 27
- [353] S. D. Mathur, *Fuzzballs and the information paradox: A Summary and conjectures*, [0810.4525]. 26
- [354] P. O. Mazur and E. Mottola, *Gravitational vacuum condensate stars*, *Proc. Nat. Acad. Sci.* **101** (2004) 9545 [gr-qc/0407075]. 26, 27
- [355] C. Barcelo, S. Liberati, S. Sonogo and M. Visser, *Fate of gravitational collapse in semiclassical gravity*, *Phys. Rev.* **D77** (2008) 044032 [0712.1130]. 26
- [356] R. Carballo-Rubio, *Stellar equilibrium in semiclassical gravity*, *Phys. Rev. Lett.* **120** (2018) 061102 [1706.05379]. 26
- [357] U. H. Danielsson, G. Dibitetto and S. Giri, *Black holes as bubbles of AdS*, *JHEP* **10** (2017) 171 [1705.10172]. 26
- [358] C. Berthiere, D. Sarkar and S. N. Solodukhin, *The fate of black hole horizons in semiclassical gravity*, *Phys. Lett.* **B786** (2018) 21 [1712.09914]. 26
- [359] A. Almheiri, D. Marolf, J. Polchinski and J. Sully, *Black Holes: Complementarity or Firewalls?*, *JHEP* **02** (2013) 062 [1207.3123]. 26, 27
- [360] S. B. Giddings, *Nonviolent information transfer from black holes: A field theory parametrization*, *Phys. Rev.* **D88** (2013) 024018 [1302.2613]. 26
- [361] S. B. Giddings, *Modulated Hawking radiation and a nonviolent channel for information release*, *Phys. Lett.* **B738** (2014) 92 [1401.5804].
- [362] S. B. Giddings, *Nonviolent unitarization: basic postulates to soft quantum structure of black holes*, *JHEP* **12** (2017) 047 [1701.08765]. 26
- [363] V. Cardoso and P. Pani, *Tests for the existence of black holes through gravitational wave echoes*, *Nat. Astron.* **1** (2017) 586 [1709.01525]. 26, 27
- [364] R. Penrose, *Gravitational Collapse: the Role of General Relativity*, *Nuovo Cimento Rivista Serie* **1** (1969) 252. 26
- [365] O. Dreyer, B. J. Kelly, B. Krishnan, L. S. Finn, D. Garrison and R. Lopez-Aleman, *Black hole spectroscopy: Testing general relativity through gravitational wave observations*, *Class. Quant. Grav.* **21** (2004) 787 [gr-qc/0309007]. 26
- [366] J. Meidam, M. Agathos, C. Van Den Broeck, J. Veitch and B. S. Sathyaprakash, *Testing the no-hair theorem with black hole ringdowns using TIGER*, *Phys. Rev.* **D90** (2014) 064009 [1406.3201].
- [367] E. Berti, A. Sesana, E. Barausse, V. Cardoso and K. Belczynski, *Spectroscopy of Kerr black holes with Earth- and space-based interferometers*, *Phys. Rev. Lett.* **117** (2016) 101102 [1605.09286].

- [368] S. Bhagwat, M. Okounkova, S. W. Ballmer, D. A. Brown, M. Giesler, M. A. Scheel et al., *On choosing the start time of binary black hole ringdowns*, *Phys. Rev.* **D97** (2018) 104065 [1711.00926].
- [369] H. Yang, K. Yagi, J. Blackman, L. Lehner, V. Paschalidis, F. Pretorius et al., *Black hole spectroscopy with coherent mode stacking*, *Phys. Rev. Lett.* **118** (2017) 161101 [1701.05808].
- [370] R. Brito, A. Buonanno and V. Raymond, *Black-hole Spectroscopy by Making Full Use of Gravitational-Wave Modeling*, *Phys. Rev.* **D98** (2018) 084038 [1805.00293].
- [371] G. Carullo et al., *Empirical tests of the black hole no-hair conjecture using gravitational-wave observations*, *Phys. Rev.* **D98** (2018) 104020 [1805.04760]. 26
- [372] L. Barack et al., *Black holes, gravitational waves and fundamental physics: a roadmap*, *Class. Quant. Grav.* **36** (2019) 143001 [1806.05195]. 26, 27
- [373] S. L. Liebling and C. Palenzuela, *Dynamical Boson Stars*, *Living Rev. Rel.* **15** (2012) 6 [1202.5809]. 26
- [374] M. A. Abramowicz, W. Kluzniak and J.-P. Lasota, *No observational proof of the black hole event-horizon*, *Astron. Astrophys.* **396** (2002) L31 [astro-ph/0207270]. 27
- [375] V. Cardoso and P. Pani, *Testing the nature of dark compact objects: a status report*, *Living Rev. Rel.* **22** (2019) 4 [1904.05363]. 27
- [376] S. Clesse and J. García-Bellido, *The clustering of massive Primordial Black Holes as Dark Matter: measuring their mass distribution with Advanced LIGO*, *Phys. Dark Univ.* **15** (2017) 142 [1603.05234]. 27
- [377] S. Bird, I. Cholis, J. B. Muñoz, Y. Ali-Haïmoud, M. Kamionkowski, E. D. Kovetz et al., *Did LIGO detect dark matter?*, *Phys. Rev. Lett.* **116** (2016) 201301 [1603.00464].
- [378] M. Sasaki, T. Suyama, T. Tanaka and S. Yokoyama, *Primordial Black Hole Scenario for the Gravitational-Wave Event GW150914*, *Phys. Rev. Lett.* **117** (2016) 061101 [1603.08338]. 27
- [379] B. J. Carr and S. W. Hawking, *Black holes in the early Universe*, *Mon. Not. Roy. Astron. Soc.* **168** (1974) 399. 27
- [380] C. T. Byrnes, M. Hindmarsh, S. Young and M. R. S. Hawkins, *Primordial black holes with an accurate QCD equation of state*, *JCAP* **1808** (2018) 041 [1801.06138]. 27, 28
- [381] LIGO SCIENTIFIC, VIRGO collaboration, *Search for Substellar-Mass Ultracompact Binaries in Advanced LIGO's First Observing Run*, *Phys. Rev. Lett.* **121** (2018) 231103 [1808.04771]. 27, 28
- [382] S. M. Koushiappas and A. Loeb, *Maximum redshift of gravitational wave merger events*, *Phys. Rev. Lett.* **119** (2017) 221104 [1708.07380]. 27
- [383] E. D. Kovetz, I. Cholis, P. C. Breysse and M. Kamionkowski, *Black hole mass function from gravitational wave measurements*, *Phys. Rev.* **D95** (2017) 103010 [1611.01157]. 27
- [384] G. Steigman and M. S. Turner, *Cosmological Constraints on the Properties of Weakly Interacting Massive Particles*, *Nucl. Phys.* **B253** (1985) 375. 27
- [385] K. Eda, Y. Itoh, S. Kuroyanagi and J. Silk, *New Probe of Dark-Matter Properties: Gravitational Waves from an Intermediate-Mass Black Hole Embedded in a Dark-Matter Minispike*, *Phys. Rev. Lett.* **110** (2013) 221101 [1301.5971]. 27
- [386] C. F. B. Macedo, P. Pani, V. Cardoso and L. C. B. Crispino, *Into the lair: gravitational-wave signatures of dark matter*, *Astrophys. J.* **774** (2013) 48 [1302.2646].
- [387] E. Barausse, V. Cardoso and P. Pani, *Can environmental effects spoil precision gravitational-wave astrophysics?*, *Phys. Rev.* **D89** (2014) 104059 [1404.7149]. 27

- [388] EROS-2 collaboration, *Limits on the Macho Content of the Galactic Halo from the EROS-2 Survey of the Magellanic Clouds*, *Astron. Astrophys.* **469** (2007) 387 [astro-ph/0607207]. 28
- [389] Ł. Wyrzykowski, S. Kozłowski, J. Skowron, A. Udalski, M. K. Szymański, M. Kubiak et al., *The OGLE view of microlensing towards the Magellanic Clouds - III. Ruling out subsolar MACHOs with the OGLE-III LMC data*, *MNRAS* **413** (2011) 493 [1012.1154]. 28
- [390] M. Zumalacarregui and U. Seljak, *Limits on stellar-mass compact objects as dark matter from gravitational lensing of type Ia supernovae*, *Phys. Rev. Lett.* **121** (2018) 141101 [1712.02240]. 28
- [391] J. M. Diego et al., *Dark Matter under the Microscope: Constraining Compact Dark Matter with Caustic Crossing Events*, *Astrophys. J.* **857** (2018) 25 [1706.10281]. 28
- [392] T. Venumadhav, L. Dai and J. Miralda-Escudé, *Microlensing of Extremely Magnified Stars near Caustics of Galaxy Clusters*, *Astrophys. J.* **850** (2017) 49 [1707.00003].
- [393] M. Oguri, J. M. Diego, N. Kaiser, P. L. Kelly and T. Broadhurst, *Understanding caustic crossings in giant arcs: characteristic scales, event rates, and constraints on compact dark matter*, *Phys. Rev.* **D97** (2018) 023518 [1710.00148]. 28
- [394] Y. Ali-Haïmoud and M. Kamionkowski, *Cosmic microwave background limits on accreting primordial black holes*, *Phys. Rev.* **D95** (2017) 043534 [1612.05644]. 28
- [395] J. Luis Bernal, N. Bellomo, A. Raccanelli and L. Verde, *Cosmological implications of Primordial Black Holes*, *JCAP* **1710** (2017) 052 [1709.07465]. 28
- [396] T. D. Brandt, *Constraints on MACHO Dark Matter from Compact Stellar Systems in Ultra-Faint Dwarf Galaxies*, *Astrophys. J. Lett.* **824** (2016) L31 [1605.03665]. 28
- [397] DES collaboration, *Farthest Neighbor: The Distant Milky Way Satellite Eridanus II*, *Astrophys. J.* **838** (2017) 8 [1611.05052]. 28
- [398] Y. Ali-Haïmoud, E. D. Kovetz and M. Kamionkowski, *Merger rate of primordial black-hole binaries*, *Phys. Rev.* **D96** (2017) 123523 [1709.06576]. 28
- [399] W. H. Press and D. N. Spergel, *Capture by the sun of a galactic population of weakly interacting massive particles*, *Astrophys. J.* **296** (1985) 679. 27
- [400] A. Gould, B. T. Draine, R. W. Romani and S. Nussinov, *Neutron Stars: Graveyard of Charged Dark Matter*, *Phys. Lett.* **B238** (1990) 337.
- [401] I. Goldman and S. Nussinov, *Weakly Interacting Massive Particles and Neutron Stars*, *Phys. Rev.* **D40** (1989) 3221.
- [402] G. Bertone and M. Fairbairn, *Compact Stars as Dark Matter Probes*, *Phys. Rev.* **D77** (2008) 043515 [0709.1485]. 27
- [403] J. Ellis, A. Hektor, G. Hütsi, K. Kannike, L. Marzola, M. Raidal et al., *Search for Dark Matter Effects on Gravitational Signals from Neutron Star Mergers*, *Phys. Lett.* **B781** (2018) 607 [1710.05540]. 27
- [404] J. Bramante, T. Linden and Y.-D. Tsai, *Searching for dark matter with neutron star mergers and quiet kilonovae*, *Phys. Rev.* **D97** (2018) 055016 [1706.00001]. 27
- [405] C. Kouvaris, P. Tinyakov and M. H. G. Tytgat, *NonPrimordial Solar Mass Black Holes*, *Phys. Rev. Lett.* **121** (2018) 221102 [1804.06740]. 27
- [406] B. Müller, *The Status of Multi-Dimensional Core-Collapse Supernova Models*, *Publ. Astron. Soc. Austral.* **33** (2016) e048 [1608.03274]. 30

- [407] D. Radice, V. Morozova, A. Burrows, D. Vartanyan and H. Nagakura, *Characterizing the Gravitational Wave Signal from Core-Collapse Supernovae*, *Astrophys. J. Lett.* **876** (2019) L9 [1812.07703]. 30
- [408] K. N. Yakunin et al., *Gravitational wave signatures of ab initio two-dimensional core collapse supernova explosion models for 12–25 solar mass stars*, *Phys. Rev.* **D92** (2015) 084040 [1505.05824]. 30, 31
- [409] V. Morozova, D. Radice, A. Burrows and D. Vartanyan, *The gravitational wave signal from core-collapse supernovae*, *Astrophys. J.* **861** (2018) 10 [1801.01914]. 30, 31
- [410] J. W. Murphy, C. D. Ott and A. Burrows, *A Model for Gravitational Wave Emission from Neutrino-Driven Core-Collapse Supernovae*, *Astrophys. J.* **707** (2009) 1173 [0907.4762]. 30, 31
- [411] H. Dimmelmeier, C. D. Ott, H.-T. Janka, A. Marek and E. Müller, *Generic Gravitational Wave Signals from the Collapse of Rotating Stellar Cores*, *Phys. Rev. Lett.* **98** (2007) 251101 [astro-ph/0702305]. 30
- [412] H. Dimmelmeier, C. D. Ott, A. Marek and H. T. Janka, *The Gravitational Wave Burst Signal from Core Collapse of Rotating Stars*, *Phys. Rev.* **D78** (2008) 064056 [0806.4953].
- [413] E. Müller and H. T. Janka, *Gravitational radiation from convective instabilities in Type II supernova explosions.*, *A&A* **317** (1997) 140. 30, 31
- [414] S. Richers, C. D. Ott, E. Abdikamalov, E. O’Connor and C. Sullivan, *Equation of State Effects on Gravitational Waves from Rotating Core Collapse*, *Phys. Rev.* **D95** (2017) 063019 [1701.02752]. 30
- [415] T. Zwerger and E. Müller, *Dynamics and gravitational wave signature of axisymmetric rotational core collapse*, *Astron. Astrophys.* **320** (1997) 209. 30
- [416] E. Abdikamalov, S. Gossan, A. M. DeMaio and C. D. Ott, *Measuring the Angular Momentum Distribution in Core-Collapse Supernova Progenitors with Gravitational Waves*, *Phys. Rev.* **D90** (2014) 044001 [1311.3678]. 30
- [417] J. Fuller, H. Klion, E. Abdikamalov and C. D. Ott, *Supernova Seismology: Gravitational Wave Signatures of Rapidly Rotating Core Collapse*, *Mon. Not. Roy. Astron. Soc.* **450** (2015) 414 [1501.06951]. 30
- [418] K. Kotake, W. Iwakami, N. Ohnishi and S. Yamada, *Ray-Tracing Analysis of Anisotropic Neutrino Radiation for Estimating Gravitational Waves in Core-Collapse Supernovae*, *Astrophys. J.* **704** (2009) 951 [0909.3622]. 30
- [419] E. Müller, H. T. Janka and A. Wongwathanarat, *Parametrized 3D models of neutrino-driven supernova explosions: Neutrino emission asymmetries and gravitational-wave signals*, *Astron. Astrophys.* **537** (2012) A63 [1106.6301]. 30, 31
- [420] H. Andresen, B. Müller, E. Müller and H.-T. Janka, *Gravitational Wave Signals from 3D Neutrino Hydrodynamics Simulations of Core-Collapse Supernovae*, *Mon. Not. Roy. Astron. Soc.* **468** (2017) 2032 [1607.05199]. 30, 31
- [421] E. Müller, M. Rampp, R. Buras, H.-T. Janka and D. H. Shoemaker, *Towards gravitational wave signals from realistic core collapse supernova models*, *Astrophys. J.* **603** (2004) 221 [astro-ph/0309833].
- [422] K. N. Yakunin et al., *Gravitational Waves from Core Collapse Supernovae*, *Class. Quant. Grav.* **27** (2010) 194005 [1005.0779]. 30, 31
- [423] B. Müller and H.-T. Janka, *A New Multi-Dimensional General Relativistic Neutrino Hydrodynamics Code for Core-Collapse Supernovae IV. The Neutrino Signal*, *Astrophys. J.* **788** (2014) 82 [1402.3415]. 30, 31
- [424] A. Marek, H. T. Janka and E. Müller, *Equation-of-State Dependent Features in Shock-Oscillation Modulated Neutrino and Gravitational-Wave Signals from Supernovae*, *Astron. Astrophys.* **496** (2009) 475 [0808.4136]. 30
- [425] A. Torres-Forné, P. Cerdá-Durán, A. Passamonti and J. A. Font, *Towards asteroseismology of core-collapse supernovae with gravitational-wave observations – I. Cowling approximation*, *Mon. Not. Roy. Astron. Soc.* **474** (2018) 5272 [1708.01920]. 30

- [426] J. W. Murphy and A. Burrows, *Criteria for Core-Collapse Supernova Explosions by the Neutrino Mechanism*, *Astrophys. J.* **688** (2008) 1159 [0805.3345].
- [427] P. Cerdá-Durán, N. DeBrye, M. A. Aloy, J. A. Font and M. Obergaulinger, *Gravitational wave signatures in black-hole-forming core collapse*, *Astrophys. J. Lett.* **779** (2013) L18 [1310.8290]. 30, 31
- [428] J. M. Blondin and A. Mezzacappa, *Pulsar spins from an instability in the accretion shock of supernovae*, *Nature* **445** (2007) 58 [astro-ph/0611680]. 30
- [429] T. Kuroda, K. Kotake and T. Takiwaki, *A new Gravitational-wave Signature From Standing Accretion Shock Instability in Supernovae*, *Astrophys. J. Lett.* **829** (2016) L14 [1605.09215]. 30, 31
- [430] T. Kuroda, K. Kotake, K. Hayama and T. Takiwaki, *Correlated Signatures of Gravitational-Wave and Neutrino Emission in Three-Dimensional General-Relativistic Core-Collapse Supernova Simulations*, *Astrophys. J.* **851** (2017) 62 [1708.05252]. 30, 31
- [431] C. D. Ott, C. Reisswig, E. Schnetter, E. O'Connor, U. Sperhake, F. Loffler et al., *Dynamics and Gravitational Wave Signature of Collapsar Formation*, *Phys. Rev. Lett.* **106** (2011) 161103 [1012.1853]. 31
- [432] K.-C. Pan, M. Liebendörfer, S. M. Couch and F.-K. Thielemann, *Equation of State Dependent Dynamics and Multi-messenger Signals from Stellar-mass Black Hole Formation*, *Astrophys. J.* **857** (2018) 13 [1710.01690]. 31
- [433] I. S. Heng, *Rotating stellar core-collapse waveform decomposition: A principal component analysis approach*, *Class. Quant. Grav.* **26** (2009) 105005 [0810.5707]. 31
- [434] J. Powell, S. E. Gossan, J. Logue and I. S. Heng, *Inferring the core-collapse supernova explosion mechanism with gravitational waves*, *Phys. Rev.* **D94** (2016) 123012 [1610.05573]. 31
- [435] K. N. Yakunin, A. Mezzacappa, P. Marronetti, E. J. Lentz, S. W. Bruenn, W. R. Hix et al., *The Gravitational Wave Signal of a Core Collapse Supernova Explosion of a 15 Solar Mass Star*, [1701.07325]. 31
- [436] S. Hild et al., *Sensitivity Studies for Third-Generation Gravitational Wave Observatories*, *Class. Quant. Grav.* **28** (2011) 094013 [1012.0908]. 31
- [437] S. Dwyer, D. Sigg, S. W. Ballmer, L. Barsotti, N. Mavalvala and M. Evans, *Gravitational wave detector with cosmological reach*, *Phys. Rev.* **D91** (2015) 082001 [1410.0612]. 31
- [438] S. Zha, E. P. O'Connor, M.-c. Chu, L.-M. Lin and S. M. Couch, *Gravitational-wave Signature of a First-order Quantum Chromodynamics Phase Transition in Core-Collapse Supernovae*, *Phys. Rev. Lett.* **125** (2020) 051102 [2007.04716]. 31
- [439] K. Kotake, *Multiple physical elements to determine the gravitational-wave signatures of core-collapse supernovae*, *Comptes Rendus Physique* **14** (2013) 318 [1110.5107]. 31
- [440] SUPER-KAMIOKANDE collaboration, *Real-Time Supernova Neutrino Burst Monitor at Super-Kamiokande*, *Astropart. Phys.* **81** (2016) 39 [1601.04778]. 31
- [441] A. Ankowski et al., *Supernova Physics at DUNE*, in *Supernova Physics at DUNE Blacksburg, Virginia, USA, March 11-12, 2016*, 2016, [1608.07853]. 31
- [442] J.-S. Lu, J. Cao, Y.-F. Li and S. Zhou, *Constraining Absolute Neutrino Masses via Detection of Galactic Supernova Neutrinos at JUNO*, *JCAP* **1505** (2015) 044 [1412.7418]. 31
- [443] ICECUBE collaboration, *IceCube Sensitivity for Low-Energy Neutrinos from Nearby Supernovae*, *Astron. Astrophys.* **535** (2011) A109 [1108.0171]. 31
- [444] N. Yu. Agafonova et al., *On-line recognition of supernova neutrino bursts in the LVD detector*, *Astropart. Phys.* **28** (2008) 516 [0710.0259]. 31

- [445] L. Cadonati, F. P. Calaprice and M. C. Chen, *Supernova neutrino detection in borexino*, *Astropart. Phys.* **16** (2002) 361 [hep-ph/0012082]. 31
- [446] KAMLAND collaboration, *Supernova detection with KamLAND*, *Nucl. Phys. Proc. Suppl.* **221** (2011) 355. 31
- [447] C. D. Ott, E. Abdikamalov, E. O'Connor, C. Reisswig, R. Haas, P. Kalmus et al., *Correlated Gravitational Wave and Neutrino Signals from General-Relativistic Rapidly Rotating Iron Core Collapse*, *Phys. Rev.* **D86** (2012) 024026 [1204.0512]. 31
- [448] T. Yokozawa, M. Asano, T. Kayano, Y. Suwa, N. Kanda, Y. Koshio et al., *Probing the Rotation of Core-collapse Supernova With a Concurrent Analysis of Gravitational Waves and Neutrinos*, *Astrophys. J.* **811** (2015) 86 [1410.2050]. 31
- [449] M. Sieniawska and M. Bejger, *Continuous gravitational waves from neutron stars: current status and prospects*, *Universe* **5** (2019) 217 [1909.12600]. 32
- [450] K. Riles, *Recent searches for continuous gravitational waves*, *Modern Physics Letters A* **32** (2017) 1730035 [1712.05897]. 32
- [451] N. Andersson, *A New Class of Unstable Modes of Rotating Relativistic Stars*, *ApJ* **502** (1998) 708 [gr-qc/9706075]. 32
- [452] J. L. Friedman and S. M. Morsink, *Axial Instability of Rotating Relativistic Stars*, *ApJ* **502** (1998) 714 [gr-qc/9706073]. 32
- [453] H. Spruit and E. Phinney, *Birth kicks as the origin of pulsar rotation*, *Nature* **393** (1998) 139 [astro-ph/9803201]. 32
- [454] E. P. J. v. d. Heuvel, *Formation of Double Neutron Stars, Millisecond Pulsars and Double Black Holes*, *J. Astrophys. Astron.* **38** (2017) 45 [1709.07636]. 32
- [455] M. Ruderman, *Neutron starquakes and pulsar periods*, *Nature (London)* **223** (1969) 597. 32
- [456] F. Douchin and P. Haensel, *A unified equation of state of dense matter and neutron star structure*, *Astron. Astrophys.* **380** (2001) 151 [astro-ph/0111092]. 32
- [457] G. Woan, M. D. Pitkin, B. Haskell, D. I. Jones and P. D. Lasky, *Evidence for a Minimum Ellipticity in Millisecond Pulsars*, *Astrophys. J. Lett.* **863** (2018) L40 [1806.02822]. 32
- [458] A. Melatos and D. J. B. Payne, *Gravitational Radiation from an Accreting Millisecond Pulsar with a Magnetically Confined Mountain*, *ApJ* **623** (2005) 1044 [astro-ph/0503287]. 32
- [459] L. Bildsten, *Gravitational Radiation and Rotation of Accreting Neutron Stars*, *ApJ* **501** (1998) L89 [astro-ph/9804325]. 32
- [460] N. Andersson, D. I. Jones and W. C. G. Ho, *Implications of an r mode in XTE J1751-305: mass, radius and spin evolution*, *MNRAS* **442** (2014) 1786 [1403.0860]. 32
- [461] U. Lee, *Excitation of a non-radial mode in a millisecond X-ray pulsar XTE J1751-305*, *MNRAS* **442** (2014) 3037 [1403.3476]. 32
- [462] D. Chakrabarty, E. H. Morgan, M. P. Muno, D. K. Galloway, R. Wijnands, M. van der Klis et al., *Nuclear-powered millisecond pulsars and the maximum spin frequency of neutron stars*, *Nature* **424** (2003) 42 [astro-ph/0307029]. 32
- [463] D. Chakrabarty, *The spin distribution of millisecond X-ray pulsars*, in *American Institute of Physics Conference Series*, R. Wijnands, D. Altamirano, P. Soleri, N. Degenaar, N. Rea, P. Casella et al., eds., vol. 1068 of *American Institute of Physics Conference Series*, pp. 67–74, Oct., 2008, [0809.4031], DOI.

- [464] A. Patruno, B. Haskell and N. Andersson, *The Spin Distribution of Fast-spinning Neutron Stars in Low-mass X-Ray Binaries: Evidence for Two Subpopulations*, *ApJ* **850** (2017) 106 [1705.07669]. 32
- [465] LIGO SCIENTIFIC, VIRGO collaboration, *All-sky search for continuous gravitational waves from isolated neutron stars using Advanced LIGO O2 data*, *Phys. Rev.* **D100** (2019) 024004 [1903.01901]. 32
- [466] LIGO SCIENTIFIC, VIRGO collaboration, *Searches for Gravitational Waves from Known Pulsars at Two Harmonics in 2015-2017 LIGO Data*, *Astrophys. J.* **879** (2019) 10 [1902.08507]. 32
- [467] V. M. Kaspi and A. Beloborodov, *Magnetars*, *Ann. Rev. Astron. Astrophys.* **55** (2017) 261 [1703.00068]. 32
- [468] R. Turolla, S. Zane and A. Watts, *Magnetars: the physics behind observations. A review*, *Rept. Prog. Phys.* **78** (2015) 116901 [1507.02924]. 32
- [469] G. Israel, T. Belloni, L. Stella, Y. Rephaeli, D. Gruber, P. G. Casella et al., *Discovery of rapid x-ray oscillations in the tail of the SGR 1806-20 hyperflare*, *Astrophys. J. Lett.* **628** (2005) L53 [astro-ph/0505255]. 32
- [470] T. E. Strohmayer and A. L. Watts, *Discovery of fast x-ray oscillations during the 1998 giant flare from SGR 1900+14*, *Astrophys. J. Lett.* **632** (2005) L111 [astro-ph/0508206]. 32
- [471] T. E. Strohmayer and A. L. Watts, *The 2004 Hyperflare from SGR 1806-20: Further Evidence for Global Torsional Vibrations*, *Astrophys. J.* **653** (2006) 593 [astro-ph/0608463]. 32
- [472] D. Huppenkothen, L. M. Heil, A. L. Watts and E. Göğüş, *Quasi-Periodic Oscillations in Short Recurring Bursts of the magnetars SGR 1806-20 and SGR 1900+14 Observed With RXTE*, *Astrophys. J.* **795** (2014) 114 [1409.7642]. 32
- [473] D. Huppenkothen et al., *Quasi-periodic Oscillations in Short Recurring Bursts of the Soft Gamma Repeater J1550–5418*, *Astrophys. J.* **787** (2014) 128 [1404.2756]. 32
- [474] P. D. Lasky, A. Melatos, V. Ravi and G. Hobbs, *Pulsar timing noise and the minimum observation time to detect gravitational waves with pulsar timing arrays*, *Mon. Not. Roy. Astron. Soc.* **449** (2015) 3293 [1503.03298]. 33
- [475] K. Glampedakis and L. Gualtieri, *Gravitational waves from single neutron stars: an advanced detector era survey*, *Astrophys. Space Sci. Libr.* **457** (2018) 673 [1709.07049]. 33
- [476] C. M. Espinoza, A. G. Lyne, B. W. Stappers and M. Kramer, *A study of 315 glitches in the rotation of 102 pulsars*, *Mon. Not. Roy. Astron. Soc.* **414** (2011) 1679 [1102.1743]. 33
- [477] A. Melatos, J. A. Douglass and T. P. Simula, *Persistent Gravitational Radiation From Glitching Pulsars*, *Astrophys. J.* **807** (2015) 132. 33
- [478] M. F. Bennett, C. A. van Eysden and A. Melatos, *Continuous-wave gravitational radiation from pulsar glitch recovery*, *Mon. Not. Roy. Astron. Soc.* **409** (2010) 1705 [1008.0236]. 33
- [479] R. Prix, S. Giampanis and C. Messenger, *Search method for long-duration gravitational-wave transients from neutron stars*, *Phys. Rev.* **D84** (2011) 023007 [1104.1704]. 33
- [480] T. Sidery, A. Passamonti and N. Andersson, *The dynamics of pulsar glitches: Contrasting phenomenology with numerical evolutions*, *Mon. Not. Roy. Astron. Soc.* **405** (2010) 1061 [0910.3918]. 33
- [481] L. C. Keer and D. I. Jones, *Neutron Star Oscillations from Starquakes*, in *Proceedings, 13th Marcel Grossmann Meeting on Recent Developments in Theoretical and Experimental General Relativity, Astrophysics, and Relativistic Field Theories (MG13): Stockholm, Sweden, July 1-7, 2012*, pp. 1996–1997, 2015, DOI. 33