

Extreme Gravity and Fundamental Physics

Astro2020 Science White Paper

EXTREME GRAVITY AND FUNDAMENTAL PHYSICS

Thematic Areas:

- Cosmology and Fundamental Physics
- Multi-Messenger Astronomy and Astrophysics

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Extreme Gravity and Fundamental Physics

In general relativity, gravitational waves are non-stationary solutions of Einstein's equations arising as a result of time-varying quadrupole and higher-order multipole moments that translate into freely propagating oscillations in the fabric of spacetime [1]. They emanate from regions of strong gravity and relativistic motion, yet the waves carry uncorrupted signature of their sources. They interact very weakly with matter and are hardly dispersed as they propagate from their sources to Earth, making them ideal for studying the dynamics of spacetime geometry [2, 3].

On September 14, 2015 the twin LIGO instruments at Hanford and Livingston made the first *direct* detection of gravitational waves [4]. Dubbed GW150914, the waves were observed for 200 milliseconds and came from the final stages of the inspiral and merger of a binary black hole system at a distance of ~ 450 Mpc. To date LIGO and Virgo in Italy have detected ten binary black hole mergers [5] that have helped to probe strong field gravity at unprecedented levels.

On August 17, 2017 LIGO and Virgo made another monumental discovery, this time the inspiral and coalescence of a pair of neutron stars [6]. Fermi Gamma-ray Space Telescope and the International Gamma-Ray Astrophysics Laboratory, both observed short gamma ray bursts 1.7 seconds after LIGO's discovery [7], thus confirming the long-held conjecture that merging binary neutron stars are progenitors of short gamma ray bursts.

Future gravitational-wave observations will enable unprecedented and unique science in *extreme gravity and fundamental physics*, that form the core topics of the Thematic Area 7 of Astro-2020 decadal survey.

- **The nature of gravity.** Can we prove Einstein wrong? What building-block principles and symmetries in nature invoked in the description of gravity can be challenged?
- **The nature of dark matter.** Is dark matter composed of particles, dark objects or modifications of gravitational interactions?
- **The nature of compact objects.** Are black holes and neutron stars the only astrophysical extreme compact objects in the Universe? What is the equation of state of densest matter?

These detections have ushered in a new era of fundamental physics. Gravitational-wave (GW) observations can be used for understanding not just the sky but also in testing general relativity in dynamical spacetimes [8–11] and in providing insights into the nature of matter under extreme physical conditions of gravity, density, and pressure [12–15]. Advanced LIGO and Virgo will only be first steps in this new endeavor that is guaranteed to change our perception of the Universe in the coming decades. Indeed, the next generation of GW observatories, such as the Einstein Telescope and Cosmic Explorer (referred to as 3G), will witness merging black holes and neutron stars when the Universe was still in its infancy assembling its first stars and black holes. At such sensitivity levels we can expect to measure extremely bright events that could reveal subtle signatures of new physics. 3G observatories promise to deliver data that could transform the landscape of physics, addressing some of the most pressing problems in fundamental physics and strong gravitational fields.

The nature of gravity.

Probing the nature of gravity and its possible implications on fundamental physics is a high-reward, even if uncertain, prospect of gravitational-wave observations. To our knowledge, astrophysical

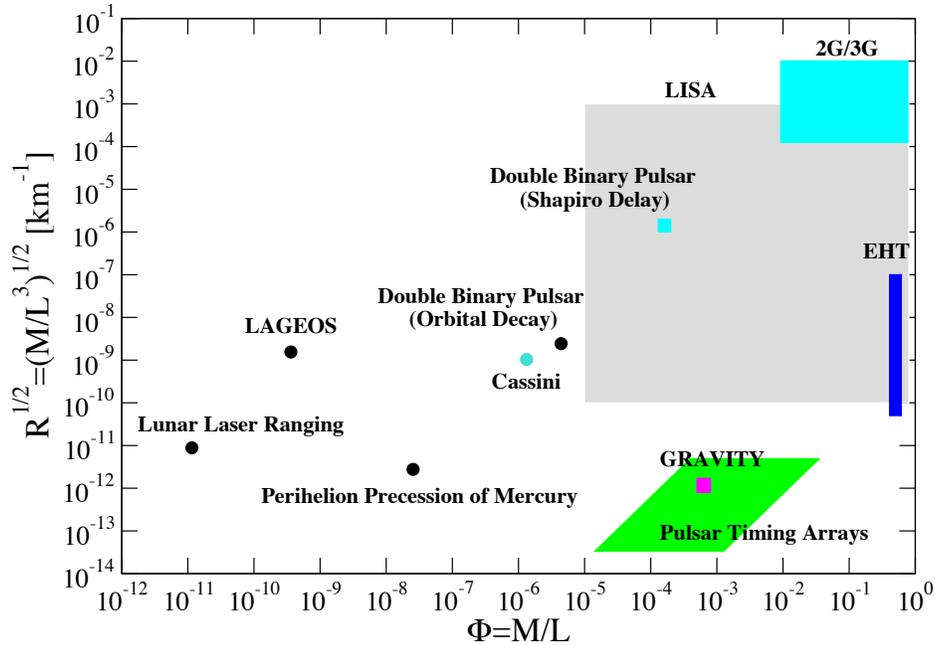


Figure 1: **Probing gravity at all scales:** Illustration of the reach in curvature scales vs potential scales targeted by different, representative, past/current/future missions. In this figure, M and L are the characteristic mass and length involved in the observable associated to each mission. For instance, in observables associated to binary systems M is the total mass and L the binary's separation, in this case M/L is related to v^2/c^2 through the virial theorem.

black holes and relativistic stars exhibit the largest curvature of spacetime accessible to us. They are, therefore, ideal systems to observe the behavior of spacetimes under the most extreme gravitational conditions. New physics indicative of departures from the basic tenants of General Relativity (GR) could reveal itself in high fidelity waveforms expected to be observed in the next generation of detectors.

Such signals would provide a unique access to extremely warped spacetimes and gain invaluable insights on GR or what might replace it as the theory of gravity governing such systems. The adjacent diagram provides a perspective of the reach of different missions/facilities and their target regime with respect to characteristic spacetime curvature (R) and gravitational potential Φ (which for binary systems can be traded with v^2/c^2 , where v is the binary's characteristic velocity and c the speed of light).

New fields, particles and polarizations Lovelock's uniqueness theorem in 4-dimensions [16] implies that departures from GR that preserve locality necessarily require the presence of extra degrees of freedom, which generically also arise from theories of quantum gravity in the low-energy limit. This often leads to violations of the strong equivalence principle through the fields' nonminimal coupling with matter. Among possible theories, those with an additional scalar field are relatively simple [17, 18] yet could give rise to exciting new strong-field phenomenology [19, 20]. Together with examples of strong-field GW signatures in more complicated scenarios inspired by the low-energy limit of quantum gravity theories [21, 22] they also serve as excellent proxies of the type of new physics we can hope to detect. In addition, if a binary's constituents can become

“dressed” with a scalar configuration [23–26], the system emits scalar waves in addition to tensorial ones, with the dominant component being dipolar emission [27] (although this may be suppressed for massive fields [28, 29]). Extra polarizations can be detected directly [10], and indirectly inferred from their effects on the system’s dynamics and consequent impact on GWs [27].

Graviton mass Recently, the possibility that gravitons could have a mass has resurfaced in theoretical physics within extensions of GR [30, 31]. The current best bound on the graviton mass from LIGO through modified dispersion relations is $m_g < 7.7 \times 10^{-23} \text{ eV}/c^2$ [10, 32] and improvements of two orders of magnitude would be possible with 3G detectors.

Lorentz violations Lorentz symmetry is regarded as a fundamental property of the Standard Model of particle physics, tested to spectacular accuracy in particle experiments [33]. In the gravitational sector, constraints are far less refined. Theories with Lorentz invariance violation (e.g., Horava-Lifschitz [34] and Einstein-Aether [35]) give rise to significant effects on black holes [36, 37], additional polarizations [38], and the propagation of GWs (e.g. through dispersion and birefringence [39]) which can be greatly constrained by 3G detectors that will observe sources at high redshifts of $z \sim 10\text{-}20$.

Parity violations Parity violations in gravity arise naturally within some flavors of string theory [40], loop quantum gravity [41] and inflationary models [42]. The associated phenomenologies are, to some degree, understood from effective theories [43]. For instance, they give rise to black holes with nontrivial pseudo-scalar configurations that violate spatial parity [44]. The resulting scalar dipole leads to a correction to the GWs produced through a of a binary inspiral and merger [21, 45, 46]. Additionally, parity violating theories can exhibit birefringence, thus impacting the characteristics of GWs tied to their handedness [47].

The nature of dark matter.

The exquisite ability of 3G detectors to probe the population and dynamics of electromagnetically dark objects throughout the Universe and harness deep insights on gravity can help reveal the nature of dark matter and answer key questions about its origin.

Black holes as dark matter candidates LIGO and Virgo discoveries have revived interest in the possibility that dark matter could be composed, in part, of black holes of masses $\sim 0.1\text{--}100 M_\odot$ [48–50]. Such black holes might have been produced from the collapse of large primordial density fluctuations in the very early Universe or during inflation [51, 52]. The exact distribution of masses depends on the model of inflation, and might be further affected by processes in the early Universe such as the quantum-chromodynamic phase transition [53].

The detection of GWs from binary systems composed of objects much lighter than stellar mass black holes, or with a mass distribution demonstrating an excess within a certain range, could point towards the existence of primordial black holes [54]. The detection of very high redshift sources would be another hint towards this formation channel [55]. With a sensitivity to observe stellar mass black holes at redshifts of $\sim 10\text{-}20$, 3G detectors will be uniquely positioned to determine their mass and spatial distribution, which will be crucial to test this hypothesis [56].

Detection of dark matter with compact objects Beyond probing whether dark matter can be partially made up of black holes, GWs can also scrutinize models where dark matter consists of particles beyond the standard model (e.g., weakly interacting massive particles [57], fuzzy dark matter [58] or axion-like particles [59]). Indeed, binary black holes evolving in a dark-matter rich environment will not only accrete the surrounding material, but also exert a gravitational drag on

the dark matter medium, which affects the inspiral dynamics [60–62]. Even though their magnitude is small, drag and accretion could have a cumulative effect over a large number of orbits that could be detected by a combination of observatories in space and 3G detectors [63].

Additionally, dark matter that interacts with standard model particles can scatter, lose energy, and be captured in astrophysical objects [64–67]. The dark-matter material eventually thermalizes with the star, and accumulates inside a finite-size core. The presence of this core might imprint a GW signature on the matter effects during the inspiral and merger of such objects in a binary system [68]. In certain models, asymmetric dark matter can accumulate and collapse to a black hole in the dense interiors of neutron stars. The core can grow by accumulating the remaining neutron star material, in effect turning neutron stars into light black holes in regions of high dark-matter density such as galactic centers [69, 70]. This provides a mechanism for creating light black holes that could be observed by 3G detectors.

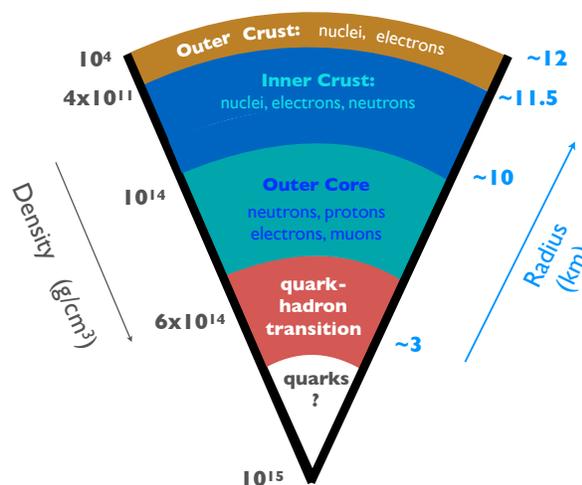
The nature of compact objects.

Observational evidence so far suggests that compact massive objects in the Universe exist in the form of black holes and neutron stars. Binary systems composed of such objects provide ideal scenarios to unravel both astrophysical and fundamental physics puzzles such as elucidating the connections of strong gravity with the most energetic phenomena in our Universe, exploring the “final state” conjecture [71] (i.e., that the end point of gravitational collapse is a Kerr black hole), and probing the existence of horizons.

Nature of black holes Black holes in isolation are the simplest objects in the Universe. Astrophysical black holes are electrically neutral and are described by just two parameters — their mass and spin angular momentum. A perturbed black hole returns to its equilibrium state by oscillating with its characteristic quasi-normal modes, whose frequency and decay time are uniquely determined by the two parameters. By detecting several quasi-normal modes 3G detectors can facilitate multiple null-hypothesis tests of the Kerr metric [9, 72, 73].

Nature of neutron stars General relativity, with input from nuclear physics, can describe the structure of ultra-dense neutron stars. However, the neutron star equation of state is currently poorly known [74]. Knowledge of the equation of state at supranuclear densities facilitated by 3G detectors will provide unprecedented insights on the properties of matter and fundamental interactions in regimes not accessible to laboratory experiments.

Signatures of matter in GWs from a binary inspiral result from a number of effects such as rotational deformations [75], various kinds of tidal effects including the excitation of internal oscillation modes of the star [76–80] and spin-tidal couplings [81, 82], and the presence of a surface instead of an event horizon [83–85].



Internal structure of a neutron star - predicted by theory. Phase transitions to states of matter containing de-confined quarks, hyperons and meson condensates are possible at the densities encountered in the inner core.

The most striking matter imprints in the waveform occur during the tidal disruption in a neutron star-black hole binary [86, 87], or the merger and post-merger epochs in binary neutron star collisions [88]. Signals from these regimes have high frequencies and are therefore very difficult to measure with current detectors. 3G detectors will improve current measurements of tidal deformability by a factor of ~ 10 and thus determine the cold equation of state significantly better, and enable unprecedented measurements of the new physics encountered during the coalescence and post-merger epochs.

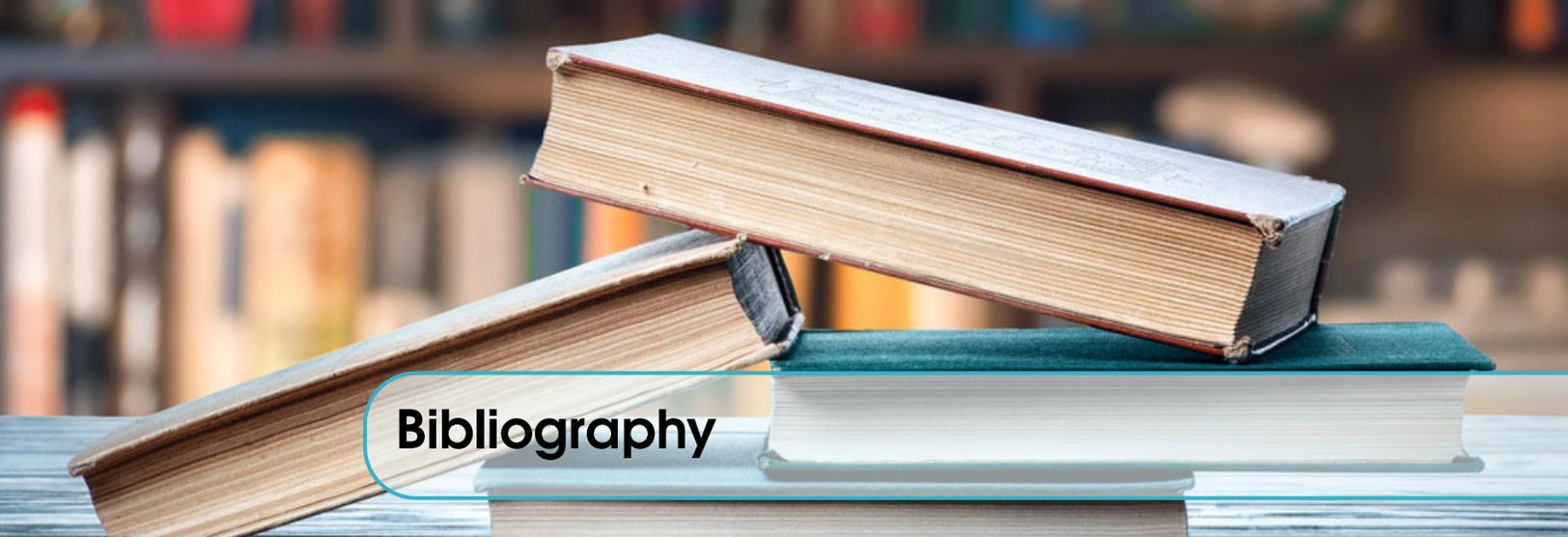
Beyond black holes and neutron stars From a phenomenological standpoint, black holes and neutron stars are just two “species” of a larger family of compact objects. More exotic species are theoretically predicted in extensions to GR, but also in particular scenarios within GR [63, 89]. For instance, exotic objects arise from beyond-standard model fundamental fields minimally coupled to gravity (e.g., boson stars [90]), in Grand Unified Theories in the early Universe (e.g., cosmic strings [91]), from exotic states of matter, as “dressed” compact objects with further structure stemming from quantum gravitational origin [92, 93] or new physics at the horizon scale (e.g., firewalls [94]), or as horizonless compact objects in a variety of scenarios (e.g., fuzzballs, gravastars, and dark stars [95–100]).

GW observations provide a unique discovery opportunity in this context, since exotic matter/dark matter might not interact electromagnetically or any electromagnetic signal from the surface of the compact object might be highly redshifted [89]. Example GW signatures from the inspiral epoch include dipole radiation as well as the variety of matter effects discussed above in the context of neutron stars [63].

Additionally, while the ringdown signal can be qualitatively similar to that of a black hole, quasi-normal modes of, e.g. gravastars, axion stars and boson stars, are different from Kerr black holes [9]. 3G detectors will have unprecedented ability to extract such modes. In addition to gravitational modes, matter modes might be excited in the ringdown of an extremely compact object, akin to fluid modes excited in a remnant neutron star [63]. In the case of certain black hole mimickers the prompt ringdown signal is identical to that of a black hole; however, these objects generically support quasi-bound trapped modes which produce a modulated train of pulses at late time. These modes appear after a delay time whose characteristics are key to test Planckian corrections at the horizon scale that could be explored with 3G detectors [89].

Bosonic clouds Ultralight bosons have been proposed in various extensions of the Standard Model [59]. When the Compton wavelength of such light bosons (masses of 10^{-21} - 10^{-11} eV) is comparable to the horizon size of a stellar or supermassive rotating black hole, superradiance can cause the spin to decay, populating bound Bohr orbits around the black hole with an exponentially large number of particles [101–103]. Such bound states, in effect “gravitational atoms”, have bosonic “clouds” with masses up to $\sim 10\%$ of the mass of the black hole [104–106]. Once formed, the clouds annihilate over a longer timescale through the emission of coherent, nearly-monochromatic, GWs [104, 107].

Alternatively, measuring the spin and mass distribution of binary black holes can provide evidence for characteristic spin down from superradiance [108–110], and explore the parameter space for ultralight bosons with 3G detectors. In addition, the presence of such clouds can be probed through the imprint of finite-size effects on the compact objects in a binary system [111]. GWs will, therefore, provide a unique window into the ultralight, weakly coupled regime of particle physics that cannot be easily probed with terrestrial experiments.



Bibliography

- [1] A. Einstein, “Über Gravitationswellen,” *Sitzungsber. K. Preuss. Akad. Wiss.* **1** (1918) 154–167.
- [2] K. S. Thorne, “Gravitational radiation,” in *Three hundred years of gravitation*, S. W. Hawking and W. Israel, eds., ch. 9, pp. 330–458. Cambridge University Press, Cambridge, 1987.
- [3] B. S. Sathyaprakash and B. F. Schutz, “Physics, Astrophysics and Cosmology with Gravitational Waves,” *Living Rev. Rel.* **12** (2009) 2.
- [4] **Virgo, LIGO Scientific** Collaboration, B. P. Abbott *et al.*, “Observation of Gravitational Waves from a Binary Black Hole Merger,” *Phys. Rev. Lett.* **116** no. 6, (2016) 061102.
- [5] B. P. Abbott *et al.*, “GWTC-1: A Gravitational-Wave Transient Catalog of Compact Binary Mergers Observed by LIGO and Virgo during the First and Second Observing Runs,” arXiv:1811.12907 [astro-ph.HE].
- [6] **Virgo, LIGO Scientific** Collaboration, B. P. Abbott *et al.*, “GW170817: Observation of Gravitational Waves from a Binary Neutron Star Inspiral,” *Phys. Rev. Lett.* **119** no. 16, (2017) 161101.
- [7] B. P. Abbott *et al.*, “Multi-messenger Observations of a Binary Neutron Star Merger,” *Astrophys. J.* **848** no. 2, (2017) L12.
- [8] N. Yunes, K. Yagi, and F. Pretorius, “Theoretical Physics Implications of the Binary Black-Hole Mergers GW150914 and GW151226,” *Phys. Rev.* **D94** no. 8, (2016) 084002.
- [9] E. Berti, K. Yagi, H. Yang, and N. Yunes, “Extreme Gravity Tests with Gravitational Waves from Compact Binary Coalescences: (II) Ringdown,” *Gen. Rel. Grav.* **50** no. 5, (2018) 49.
- [10] **LIGO Scientific, Virgo** Collaboration, B. P. Abbott *et al.*, “Tests of general relativity with GW150914,” *Phys. Rev. Lett.* **116** no. 22, (2016) 221101, arXiv:1602.03841 [gr-qc]. [Erratum: *Phys. Rev. Lett.* 121, no. 12, 129902 (2018)].
- [11] **LIGO Scientific, Virgo** Collaboration, B. P. Abbott *et al.*, “Tests of General Relativity with GW170817,” arXiv:1811.00364 [gr-qc].

-
- [12] S. De, D. Finstad, J. M. Lattimer, D. A. Brown, E. Berger, and C. M. Biwer, “Tidal Deformabilities and Radii of Neutron Stars from the Observation of GW170817,” *Phys. Rev. Lett.* **121** no. 9, (2018) 091102. [Erratum: *Phys. Rev. Lett.* 121, no. 25, 259902 (2018)].
- [13] I. Tews, J. Margueron, and S. Reddy, “Critical examination of constraints on the equation of state of dense matter obtained from GW170817,” *Phys. Rev.* **C98** no. 4, (2018) 045804.
- [14] E. Annala, T. Gorda, A. Kurkela, and A. Vuorinen, “Gravitational-wave constraints on the neutron-star-matter Equation of State,” *Phys. Rev. Lett.* **120** no. 17, (2018) 172703.
- [15] B. P. Abbott *et al.*, “GW170817: Measurements of neutron star radii and equation of state,” *Phys. Rev. Lett.* **121** no. 16, (2018) 161101.
- [16] D. Lovelock, “The Einstein tensor and its generalizations,” *J. Math. Phys.* **12** (1971) 498–501.
- [17] C. Brans and R. H. Dicke, “Mach’s principle and a relativistic theory of gravitation,” *Phys. Rev.* **124** (1961) 925–935.
- [18] Y. Fujii and K. Maeda, *The scalar-tensor theory of gravitation*. Cambridge University Press, 2007.
- [19] C. Palenzuela, E. Barausse, M. Ponce, and L. Lehner, “Dynamical scalarization of neutron stars in scalar-tensor gravity theories,” *Phys. Rev.* **D89** no. 4, (2014) 044024.
- [20] M. Shibata, K. Taniguchi, H. Okawa, and A. Buonanno, “Coalescence of binary neutron stars in a scalar-tensor theory of gravity,” *Phys. Rev.* **D89** no. 8, (2014) 084005.
- [21] M. Okounkova, L. C. Stein, M. A. Scheel, and D. A. Hemberger, “Numerical binary black hole mergers in dynamical Chern-Simons gravity: Scalar field,” *Phys. Rev.* **D96** no. 4, (2017) 044020.
- [22] H. Witek, L. Gualtieri, P. Pani, and T. P. Sotiriou, “Black holes and binary mergers in scalar Gauss-Bonnet gravity: scalar field dynamics,” [arXiv:1810.05177](https://arxiv.org/abs/1810.05177) [gr-qc].
- [23] T. Damour and G. Esposito-Farese, “Nonperturbative strong field effects in tensor - scalar theories of gravitation,” *Phys. Rev. Lett.* **70** (1993) 2220–2223.
- [24] P. Kanti, N. E. Mavromatos, J. Rizos, K. Tamvakis, and E. Winstanley, “Dilatonic black holes in higher curvature string gravity,” *Phys. Rev.* **D54** (1996) 5049–5058.
- [25] S. Mignemi and N. R. Stewart, “Charged black holes in effective string theory,” *Phys. Rev.* **D47** (1993) 5259–5269.
- [26] G. Antoniou, A. Bakopoulos, and P. Kanti, “Evasion of No-Hair Theorems and Novel Black-Hole Solutions in Gauss-Bonnet Theories,” *Phys. Rev. Lett.* **120** no. 13, (2018) 131102.
- [27] C. M. Will, “The Confrontation between General Relativity and Experiment,” *Living Rev. Rel.* **17** (2014) 4.

-
- [28] J. Alsing, E. Berti, C. M. Will, and H. Zaglauer, “Gravitational radiation from compact binary systems in the massive Brans-Dicke theory of gravity,” *Phys. Rev.* **D85** (2012) 064041.
- [29] L. Sagunski, J. Zhang, M. C. Johnson, L. Lehner, M. Sakellariadou, S. L. Liebling, C. Palenzuela, and D. Neilsen, “Neutron star mergers as a probe of modifications of general relativity with finite-range scalar forces,” *Phys. Rev.* **D97** no. 6, (2018) 064016.
- [30] C. de Rham, G. Gabadadze, and A. J. Tolley, “Resummation of Massive Gravity,” *Phys. Rev. Lett.* **106** (2011) 231101.
- [31] S. F. Hassan and R. A. Rosen, “Resolving the Ghost Problem in non-Linear Massive Gravity,” *Phys. Rev. Lett.* **108** (2012) 041101.
- [32] **VIRGO, LIGO Scientific** Collaboration, B. P. Abbott *et al.*, “GW170104: Observation of a 50-Solar-Mass Binary Black Hole Coalescence at Redshift 0.2,” *Phys. Rev. Lett.* **118** no. 22, (2017) 221101.
- [33] D. Mattingly, “Modern tests of Lorentz invariance,” *Living Rev. Rel.* **8** (2005) 5.
- [34] P. Horava, “Quantum Gravity at a Lifshitz Point,” *Phys. Rev.* **D79** (2009) 084008.
- [35] T. Jacobson and D. Mattingly, “Gravity with a dynamical preferred frame,” *Phys. Rev.* **D64** (2001) 024028.
- [36] C. Eling and T. Jacobson, “Black Holes in Einstein-Aether Theory,” *Class. Quant. Grav.* **23** (2006) 5643–5660. [Erratum: *Class. Quant. Grav.* 27,049802(2010)].
- [37] E. Barausse, T. Jacobson, and T. P. Sotiriou, “Black holes in Einstein-aether and Horava-Lifshitz gravity,” *Phys. Rev.* **D83** (2011) 124043.
- [38] T. P. Sotiriou, “Detecting Lorentz Violations with Gravitational Waves from Black Hole Binaries,” *Phys. Rev. Lett.* **120** no. 4, (2018) 041104.
- [39] V. A. Kostelecky and M. Mewes, “Testing local Lorentz invariance with gravitational waves,” *Phys. Lett.* **B757** (2016) 510–514.
- [40] M. B. Green, J. H. Schwarz, and E. Witten, *SUPERSTRING THEORY. VOL. 2: LOOP AMPLITUDES, ANOMALIES AND PHENOMENOLOGY*. 1988.
- [41] A. Ashtekar, A. P. Balachandran, and S. Jo, “The CP Problem in Quantum Gravity,” *Int. J. Mod. Phys.* **A4** (1989) 1493.
- [42] S. Weinberg, “Effective Field Theory for Inflation,” *Phys. Rev.* **D77** (2008) 123541.
- [43] R. Jackiw and S. Y. Pi, “Chern-Simons modification of general relativity,” *Phys. Rev.* **D68** (2003) 104012.
- [44] N. Yunes and F. Pretorius, “Dynamical Chern-Simons Modified Gravity. I. Spinning Black Holes in the Slow-Rotation Approximation,” *Phys. Rev.* **D79** (2009) 084043.

-
- [45] C. F. Sopuerta and N. Yunes, “Extreme and Intermediate-Mass Ratio Inspirals in Dynamical Chern-Simons Modified Gravity,” *Phys. Rev.* **D80** (2009) 064006.
- [46] K. Yagi, N. Yunes, and T. Tanaka, “Gravitational Waves from Quasi-Circular Black Hole Binaries in Dynamical Chern-Simons Gravity,” *Phys. Rev. Lett.* **109** (2012) 251105. [Erratum: *Phys. Rev. Lett.* 116, no. 16, 169902 (2016)].
- [47] K. Yagi and H. Yang, “Probing Gravitational Parity Violation with Gravitational Waves from Stellar-mass Black Hole Binaries,” *Phys. Rev.* **D97** no. 10, (2018) 104018.
- [48] S. Clesse and J. Garcia-Bellido, “The clustering of massive Primordial Black Holes as Dark Matter: measuring their mass distribution with Advanced LIGO,” *Phys. Dark Univ.* **15** (2017) 142–147.
- [49] S. Bird, I. Cholis, J. B. Munoz, Y. Ali-Haïmoud, M. Kamionkowski, E. D. Kovetz, A. Raccanelli, and A. G. Riess, “Did LIGO detect dark matter?,” *Phys. Rev. Lett.* **116** no. 20, (2016) 201301.
- [50] M. Sasaki, T. Suyama, T. Tanaka, and S. Yokoyama, “Primordial Black Hole Scenario for the Gravitational-Wave Event GW150914,” *Phys. Rev. Lett.* **117** no. 6, (2016) 061101. [erratum: *Phys. Rev. Lett.* 121, no. 5, 059901 (2018)].
- [51] B. J. Carr and S. W. Hawking, “Black holes in the early Universe,” *Mon. Not. Roy. Astron. Soc.* **168** (1974) 399–415.
- [52] M. Sasaki, T. Suyama, T. Tanaka, and S. Yokoyama, “Primordial black holes - perspectives in gravitational wave astronomy,” *Class. Quant. Grav.* **35** no. 6, (2018) 063001.
- [53] C. T. Byrnes, M. Hindmarsh, S. Young, and M. R. S. Hawkins, “Primordial black holes with an accurate QCD equation of state,” *JCAP* **1808** no. 08, (2018) 041.
- [54] **Virgo, LIGO Scientific** Collaboration, B. P. Abbott *et al.*, “Search for sub-solar mass ultracompact binaries in Advanced LIGO’s first observing run,” [arXiv:1808.04771](https://arxiv.org/abs/1808.04771) [astro-ph.CO].
- [55] S. M. Koushiappas and A. Loeb, “Maximum redshift of gravitational wave merger events,” *Phys. Rev. Lett.* **119** no. 22, (2017) 221104.
- [56] E. D. Kovetz, I. Cholis, P. C. Breyse, and M. Kamionkowski, “Black hole mass function from gravitational wave measurements,” *Phys. Rev.* **D95** no. 10, (2017) 103010.
- [57] G. Steigman and M. S. Turner, “Cosmological Constraints on the Properties of Weakly Interacting Massive Particles,” *Nucl. Phys.* **B253** (1985) 375–386.
- [58] L. Hui, J. P. Ostriker, S. Tremaine, and E. Witten, “Ultralight scalars as cosmological dark matter,” *Phys. Rev.* **D95** no. 4, (2017) 043541.
- [59] R. Essig *et al.*, “Working Group Report: New Light Weakly Coupled Particles,” in *Proceedings, 2013 Community Summer Study on the Future of U.S. Particle Physics: Snowmass on the Mississippi (CSS2013): Minneapolis, MN, USA, July 29-August 6, 2013.*

-
- [60] K. Eda, Y. Itoh, S. Kuroyanagi, and J. Silk, “New Probe of Dark-Matter Properties: Gravitational Waves from an Intermediate-Mass Black Hole Embedded in a Dark-Matter Minispikes,” *Phys. Rev. Lett.* **110** no. 22, (2013) 221101.
- [61] C. F. B. Macedo, P. Pani, V. Cardoso, and L. C. B. Crispino, “Into the lair: gravitational-wave signatures of dark matter,” *Astrophys. J.* **774** (2013) 48.
- [62] E. Barausse, V. Cardoso, and P. Pani, “Can environmental effects spoil precision gravitational-wave astrophysics?,” *Phys. Rev.* **D89** no. 10, (2014) 104059.
- [63] L. Barack *et al.*, “Black holes, gravitational waves and fundamental physics: a roadmap,”.
- [64] W. H. Press and D. N. Spergel, “Capture by the sun of a galactic population of weakly interacting massive particles,” *Astrophys. J.* **296** (1985) 679–684. [,277(1985)].
- [65] A. Gould, B. T. Draine, R. W. Romani, and S. Nussinov, “Neutron Stars: Graveyard of Charged Dark Matter,” *Phys. Lett.* **B238** (1990) 337–343.
- [66] I. Goldman and S. Nussinov, “Weakly Interacting Massive Particles and Neutron Stars,” *Phys. Rev.* **D40** (1989) 3221–3230.
- [67] G. Bertone and M. Fairbairn, “Compact Stars as Dark Matter Probes,” *Phys. Rev.* **D77** (2008) 043515.
- [68] J. Ellis, A. Hektor, G. Huetsi, K. Kannike, L. Marzola, M. Raidal, and V. Vaskonen, “Search for Dark Matter Effects on Gravitational Signals from Neutron Star Mergers,” *Phys. Lett.* **B781** (2018) 607–610.
- [69] J. Bramante, T. Linden, and Y.-D. Tsai, “Searching for dark matter with neutron star mergers and quiet kilonovae,” *Phys. Rev.* **D97** no. 5, (2018) 055016.
- [70] C. Kouvaris, P. Tinyakov, and M. H. G. Tytgat, “NonPrimordial Solar Mass Black Holes,” *Phys. Rev. Lett.* **121** no. 22, (2018) 221102.
- [71] R. Penrose, “Gravitational Collapse: the Role of General Relativity,” *Nuovo Cimento Rivista Serie* **1** (1969) .
- [72] O. Dreyer, B. J. Kelly, B. Krishnan, L. S. Finn, D. Garrison, and R. Lopez-Aleman, “Black hole spectroscopy: Testing general relativity through gravitational wave observations,” *Class. Quant. Grav.* **21** (2004) 787–804.
- [73] E. Berti, A. Sesana, E. Barausse, V. Cardoso, and K. Belczynski, “Spectroscopy of Kerr black holes with Earth- and space-based interferometers,” *Phys. Rev. Lett.* **117** no. 10, (2016) 101102.
- [74] G. Baym, T. Hatsuda, T. Kojo, P. D. Powell, Y. Song, and T. Takatsuka, “From hadrons to quarks in neutron stars: a review,” *Rept. Prog. Phys.* **81** no. 5, (2018) 056902, arXiv:1707.04966 [astro-ph.HE].

- [75] E. Poisson, “Gravitational waves from inspiraling compact binaries: The Quadrupole moment term,” *Phys. Rev.* **D57** (1998) 5287–5290.
- [76] K. D. Kokkotas and G. Schafer, “Tidal and tidal resonant effects in coalescing binaries,” *Mon. Not. Roy. Astron. Soc.* **275** (1995) 301.
- [77] D. Lai, “Resonant oscillations and tidal heating in coalescing binary neutron stars,” *Mon. Not. Roy. Astron. Soc.* **270** (1994) 611.
- [78] M. Shibata, “Effects of tidal resonances in coalescing compact binary systems,” *Prog. Theor. Phys.* **91** (1994) 871–884.
- [79] E. E. Flanagan and E. Racine, “Gravitomagnetic resonant excitation of Rossby modes in coalescing neutron star binaries,” *Phys. Rev.* **D75** (2007) 044001.
- [80] E. E. Flanagan and T. Hinderer, “Constraining neutron star tidal Love numbers with gravitational wave detectors,” *Phys. Rev.* **D77** (2008) 021502.
- [81] P. Landry, “Rotational-tidal phasing of the binary neutron star waveform,”.
- [82] T. Abdelsalhin, L. Gualtieri, and P. Pani, “Post-Newtonian spin-tidal couplings for compact binaries,” *Phys. Rev.* **D98** no. 10, (2018) 104046.
- [83] J. B. Hartle, “Tidal Friction in Slowly Rotating Black Holes,” *Phys. Rev.* **D8** (1973) 1010–1024.
- [84] K. Alvi, “Energy and angular momentum flow into a black hole in a binary,” *Phys. Rev.* **D64** (2001) 104020.
- [85] A. Maselli, P. Pani, V. Cardoso, T. Abdelsalhin, L. Gualtieri, and V. Ferrari, “Probing Planckian corrections at the horizon scale with LISA binaries,” *Phys. Rev. Lett.* **120** no. 8, (2018) 081101.
- [86] J. M. Lattimer and D. N. Schramm, “Black-hole-neutron-star collisions,” *Astrophys. J.* **192** (1974) L145.
- [87] M. Shibata and K. Taniguchi, “Coalescence of Black Hole-Neutron Star Binaries,” *Living Rev. Rel.* **14** (2011) 6.
- [88] L. Baiotti and L. Rezzolla, “Binary neutron star mergers: a review of Einstein’s richest laboratory,” *Rept. Prog. Phys.* **80** no. 9, (2017) 096901.
- [89] V. Cardoso and P. Pani, “Tests for the existence of black holes through gravitational wave echoes,” *Nat. Astron.* **1** no. 9, (2017) 586–591.
- [90] S. L. Liebling and C. Palenzuela, “Dynamical Boson Stars,” *Living Rev. Rel.* **15** (2012) 6. [Living Rev. Rel.20,no.1,5(2017)].
- [91] R. Jeannerot, J. Rocher, and M. Sakellariadou, “How generic is cosmic string formation in SUSY GUTs,” *Phys. Rev.* **D68** (2003) 103514.

-
- [92] S. B. Giddings, “Nonviolent information transfer from black holes: A field theory parametrization,” *Phys. Rev.* **D88** no. 2, (2013) 024018.
- [93] S. B. Giddings, “Nonviolent unitarization: basic postulates to soft quantum structure of black holes,” *JHEP* **12** (2017) 047.
- [94] A. Almheiri, D. Marolf, J. Polchinski, and J. Sully, “Black Holes: Complementarity or Firewalls?,” *JHEP* **02** (2013) 062.
- [95] S. D. Mathur, “The Fuzzball proposal for black holes: An Elementary review,” *Fortsch. Phys.* **53** (2005) 793–827.
- [96] P. O. Mazur and E. Mottola, “Gravitational vacuum condensate stars,” *Proc. Nat. Acad. Sci.* **101** (2004) 9545–9550.
- [97] C. Barcelo, S. Liberati, S. Sonego, and M. Visser, “Fate of gravitational collapse in semiclassical gravity,” *Phys. Rev.* **D77** (2008) 044032.
- [98] R. Carballo-Rubio, “Stellar equilibrium in semiclassical gravity,” *Phys. Rev. Lett.* **120** no. 6, (2018) 061102.
- [99] U. H. Danielsson, G. Dibitetto, and S. Giri, “Black holes as bubbles of AdS,” *JHEP* **10** (2017) 171.
- [100] C. Berthiere, D. Sarkar, and S. N. Solodukhin, “The fate of black hole horizons in semiclassical gravity,” *Phys. Lett.* **B786** (2018) 21–27.
- [101] A. Arvanitaki, S. Dimopoulos, S. Dubovsky, N. Kaloper, and J. March-Russell, “String Axiverse,” *Phys. Rev.* **D81** (2010) 123530.
- [102] P. Pani, V. Cardoso, L. Gualtieri, E. Berti, and A. Ishibashi, “Black hole bombs and photon mass bounds,” *Phys. Rev. Lett.* **109** (2012) 131102.
- [103] R. Brito, V. Cardoso, and P. Pani, “Massive spin-2 fields on black hole spacetimes: Instability of the Schwarzschild and Kerr solutions and bounds on the graviton mass,” *Phys. Rev.* **D88** no. 2, (2013) 023514.
- [104] A. Arvanitaki and S. Dubovsky, “Exploring the String Axiverse with Precision Black Hole Physics,” *Phys. Rev.* **D83** (2011) 044026.
- [105] R. Brito, V. Cardoso, and P. Pani, “Black holes as particle detectors: evolution of superradiant instabilities,” *Class. Quant. Grav.* **32** no. 13, (2015) 134001.
- [106] W. E. East and F. Pretorius, “Superradiant Instability and Backreaction of Massive Vector Fields around Kerr Black Holes,” *Phys. Rev. Lett.* **119** no. 4, (2017) 041101.
- [107] A. Arvanitaki, M. Baryakhtar, and X. Huang, “Discovering the QCD Axion with Black Holes and Gravitational Waves,” *Phys. Rev.* **D91** no. 8, (2015) 084011.

- [108] A. Arvanitaki, M. Baryakhtar, S. Dimopoulos, S. Dubovsky, and R. Lasenby, “Black Hole Mergers and the QCD Axion at Advanced LIGO,” *Phys. Rev.* **D95** no. 4, (2017) 043001.
- [109] M. Baryakhtar, R. Lasenby, and M. Teo, “Black Hole Superradiance Signatures of Ultralight Vectors,” *Phys. Rev.* **D96** no. 3, (2017) 035019.
- [110] R. Brito, V. Cardoso, and P. Pani, “Superradiance,” *Lect. Notes Phys.* **906** (2015) pp.1–237.
- [111] D. Baumann, H. S. Chia, and R. A. Porto, “Probing Ultralight Bosons with Binary Black Holes,” *Phys. Rev.* **D99** no. 4, (2019) 044001.