

A Virtual Event Designed For The Masses

Now Available On-Demand!

Scale up your research and translate your results more rapidly and simply than ever before. Welcome to vLC-MS.com - the event for Orbitrap Exploris mass spectrometers and much more!

Tune in to:

- Explore the LC-MS portfolio and meet the expanded Orbitrap Exploris MS system in our staffed Exhibit Hall.
- Learn from mass spectrometry experts, such as Professor Alexander Makarov himself, about Orbitrap mass spectrometry technology and the applications it enables.
- Browse posters and short presentations in our application area.

Event Highlights:

Prof. Alexander Makarov



Dr. Christian Münch



Thomas Moehring



RESEARCH ARTICLE

Ligand-protein interactions in lysozyme investigated through a dual-resolution model

Raffaele Fiorentini¹ | Kurt Kremer¹ | Raffaello Potestio^{2,3} ¹Max Planck Institute for Polymer Research, Mainz, Germany²Physics Department, University of Trento, Trento, Italy³INFN-TIFPA, Trento Institute for Fundamental Physics and Applications, Trento, Italy**Correspondence**Raffaello Potestio, Physics Department, University of Trento, via Sommarive, 14 I-38123 Trento, Italy.
Email: raffaello.potestio@unitn.it**Funding information**

European Research Council, Grant/Award Numbers: 340906, 758588

Abstract

A fully atomistic (AT) modeling of biological macromolecules at relevant length- and time-scales is often cumbersome or not even desirable, both in terms of computational effort required and a posteriori analysis. This difficulty can be overcome with the use of multiresolution models, in which different regions of the same system are concurrently described at different levels of detail. In enzymes, computationally expensive AT detail is crucial in the modeling of the active site in order to capture, for example, the chemically subtle process of ligand binding. In contrast, important yet more collective properties of the remainder of the protein can be reproduced with a coarser description. In the present work, we demonstrate the effectiveness of this approach through the calculation of the binding free energy of hen egg white lysozyme with the inhibitor di-N-acetylchitotriose. Particular attention is paid to the impact of the mapping, that is, the selection of AT and coarse-grained residues, on the binding free energy. It is shown that, in spite of small variations of the binding free energy with respect to the active site resolution, the separate contributions coming from different energetic terms (such as electrostatic and van der Waals interactions) manifest a stronger dependence on the mapping, thus pointing to the existence of an optimal level of intermediate resolution.

KEYWORDS

coarse-graining, dual-resolution modeling, free energy calculation, multiscale modeling, protein-ligand binding

1 | INTRODUCTION

One of the most relevant challenges of computational biochemistry and biophysics is the accurate calculation of binding free energies,¹⁻³ which represents one of the key steps in the identification of pharmaceutical targets as well as in the development of new drugs.⁴⁻⁶ However, the large sizes of the proteins under examination (often above the hundreds of residues), as well as the necessity to screen through large datasets of potential candidate drugs they can interact with, make this effort onerous in terms of time and computational resources.

A promising way to mitigate these limitations is the use of multiple-resolution models of the protein, that is, representations in which different parts of the molecule are concurrently described at different levels of accuracy.⁷⁻¹² The chemically relevant part of the protein, for example, the active site, is modeled at a higher level of detail, typically atomistic (AT). For the remainder, on the contrary, a simplified representation is used, where several atoms are lumped together in effective interaction sites. The working hypothesis underlying these methods is that only a relatively small part of the molecule requires an explicitly AT treatment; the remainder, in fact, is mainly

This is an open access article under the terms of the Creative Commons Attribution License, which permits use, distribution and reproduction in any medium, provided the original work is properly cited.

© 2020 The Authors. *Proteins: Structure, Function, and Bioinformatics* published by Wiley Periodicals, Inc.

responsible for large-scale, collective fluctuations whose function-oriented role is well recognized and prominent,^{12–16} however also prone to be accurately reproduced by lower-resolution representations.^{17–22} Hence, the resulting model favorably joins the accuracy of an AT description where needed and the computational efficiency of a coarse-grained (CG) one where possible.

In order to take full advantage of the dual-resolution approach to protein modeling, though, one has to solve a few key open issues: first, the definition of the appropriate CG model to employ in the low-resolution part^{22–30}; second, the coupling between high- and low-resolution models, which has to be performed so as to guarantee that the observables of interest are reproduced with respect to the reference provided for example by a fully AT simulation. This issue entails a further one, namely the identification of observables apt to quantify the fidelity with which the behavior of the system is reproduced by the dual-resolution model; third, the selection of the subpart of the molecule that *requires* a high-resolution modeling. In the present work, we will focus specifically on this third aspect.

Various methods and approaches have been developed in the past few years to describe proteins in dual resolution.^{8–11,31,32} In general, the high-resolution part is modeled at the all-atom level, making use of one of the several AT force fields available. The CG representations range from simple bead-spring elastic networks^{12,17,20} to more sophisticated Gō-type models.⁸ Other approaches maintain the high-resolution description for the solute while employing a simplified model for the solvent, with varying degrees of detail depending on the specific systems and applications^{33–38}; among these, some treat the solvent with an adaptive resolution approach, that is, solvent molecules are AT in proximity of the solute and smoothly blend in a CG representation away from it.^{39–47}

Recently, we have proposed a dual-resolution model¹² where, in the CG part, only the C_α carbons of the protein chain are retained and connected one with the other by harmonic bonds. This model has been employed in the present work with the aim of assessing the accuracy of a hybrid AT/CG description of a protein for binding free energy calculations. The system under examination is hen egg-white lysozyme (HEWL) in explicit water, bound to a sugar substrate, di-N-acetylchitotriose. We carried out calculations of the binding free energy of the ligand in the active site, with a 2-fold objective. In fact, not only we aimed at verifying that the computed quantity in the dual-resolution model matches a reference, all-atom calculation; but rather we also investigated the impact of different choices in the definition of the high-resolution subdomain. This aspect bears the highest prominence, as it is becoming increasingly more evident that a crucial component in the construction of accurate and effective low-resolution models for biological and soft matter systems is represented by the mapping,^{12,29,30} that is, the particular selection of collective variables employed to describe the system. Here, we provide novel evidence of this general property in the context of a dual-resolution model of a biomolecule, and describe a broadly applicable strategy to tackle this issue.

2 | METHODS

The system under examination in the present work is HEWL in aqueous solution. In this model, the binding site of the enzyme and the substrate molecule, the inhibitor di-N-acetylchitotriose, are represented with AT detail. The protein model employed is not adaptive, that is, the resolution of a given residue is fixed—either AT or CG—and does not change throughout a simulation. However, at difference with other works,^{7,8,46} several values of the number of protein residues treated at high resolution have been explored and employed in independent calculations. The impact of choosing different numbers of active site residues to model at the AT level is a central aspect of this study. The CG model employed to describe the low-resolution part of the protein is a simple bead-spring representation where the selected sites (namely the C_α atoms) are connected by elastic bonds penalizing the deviations from the distances that interacting atoms have in the reference conformation. Two values of elastic constants are employed, one for C_α's along the chain, and one for all other bonds. Water molecules are described in AT detail throughout the whole simulation box: the interaction with the high-resolution part of the protein takes place through the standard all-atom force field, while the interaction with the CG beads is mediated by a purely repulsive potential acting on the sole oxygen atom.

Hereafter we provide a detailed description of the model. We first discuss the calculation of the binding free energy ΔG_{bind} , then we outline the dual-resolution model and its coupling to the AT part, and finally report information about the simulation setup. Further details are made available in the Supporting Information.

2.1 | Binding free energy calculation

One of the key points of this work is the calculation of the protein-ligand binding free energy ΔG_{bind} , which quantifies the affinity of a molecule toward a protein.^{1–3} As such, it plays a prominent role in the investigation of the biochemical function and activity of enzymes and similar biomolecules, and in the development of effective drugs.

ΔG_{bind} is defined as the difference between the free energy of the system in the configuration in which the ligand is bound to the active site (G_{b}) and the corresponding value when the ligand is absent (G_{ub}):

$$\Delta G_{\text{bind}} = G_{\text{b}} - G_{\text{ub}}. \quad (1)$$

This value, in the specific case under examination, varies according to the number of active site residues modeled with AT resolution, as we will see in Section 3.

The free energy difference between two states is here computed by means of thermodynamic integration (TI).⁴⁸ Specifically, a scalar $\lambda \in [0, 1]$ is defined that parameterizes the potential energy of the system as $U_{\lambda}(r) = \lambda U_A(r) + (1 - \lambda)U_B(r)$ connecting the states A and B. The sought quantity is given by:

contribute to the total binding free energy with terms of the same order of magnitude, as shown in the first and second column. On the other hand, the third term of Equation (3), that is ΔG_{r_off} , is one order of magnitude smaller than the previous two (as shown later in Section 3); however, it is not negligible for the calculation of the overall binding free energy.

2.2 | Dual-resolution protein model

Proteins undergo both high frequency, localized fluctuations about transient conformational substates, and slower, more global transitions between them.^{49,50} In the present molecular modeling approach, those local fluctuations that can play an important role in the biological function of the protein of interest are allowed by the all-atom description of the binding site. The set of these protein residues that are modeled with AT detail does not change during the simulation, that is, the protein has a fixed, position- and time-independent dual resolution. The rest of the protein is described through a CG, lower-resolution model. If, on the one hand, it is reasonable to expect that regions of the molecule far away from the active site have a negligible direct impact on the latter, on the other hand the collective fluctuations that they determine are important to modulate the structure of those residues involved in the binding.^{22,51}

Hence, to ensure the correct structure and conformational fluctuations of the binding site, it is necessary to provide a representation of the remainder of the molecule that, albeit lower-resolution, is nonetheless capable of reproducing the appropriate large-scale dynamics.

To describe the lower-resolution part we thus employ an elastic network model (ENM),^{12,17} in which each residue is mapped onto a bead whose position corresponds to the C_α atom in the AT

description. These beads are connected by harmonic springs as shown in Figure 2.

The potential energy is given by:

$$E = \sum_i \sum_j k_{ij} (r_{ij} - r_{ij}^0)^2 \theta(r_c - r_{ij}^0) \quad (4)$$

with spring constants k_{ij} , equilibrium distance r_{ij}^0 , a cutoff distance r_c , i and j are the node index, and $\theta(r)$ is a Heaviside function taking value 1 if $r > 0$ and 0 otherwise. In this model we made use of two different elastic constants: a very stiff spring (k_b) for consecutive beads, represented in blue in Figure 2; and a weaker spring k_{nb} for not consecutive beads whose distance in the reference (native) conformation lies below a fixed cutoff (in green).

The ENM used here is parameterized to reproduce the conformational fluctuations of the reference all-atom model, these being quantified by the root mean square fluctuations (RMSF) of the all C_α atoms of the system.¹² The residues in direct contact (H-bonding or hydrophobic contact) with the substrate are modeled with all-atom detail; in order to select the other binding site residues to be described at the AT level, we sorted them by increasing distance of their center of mass from the closest ligand atom. The solvent is treated with all-atom detail and it surrounds the dual-resolution protein. The water-CG protein interaction consists in a simple excluded volume term, modeled via a Weeks-Chandler-Anderson (WCA) potential.⁵² The details about the procedure followed to determine the ENM elastic constants and the excluded volume interaction are provided in the Supporting Information, while the numerical values of the resulting parameters are reported hereafter.

As anticipated, the focus of the present work lies in the analysis of the impact that a modulation of the resolution of a protein in proximity of the active site can have on physical and mechanical properties

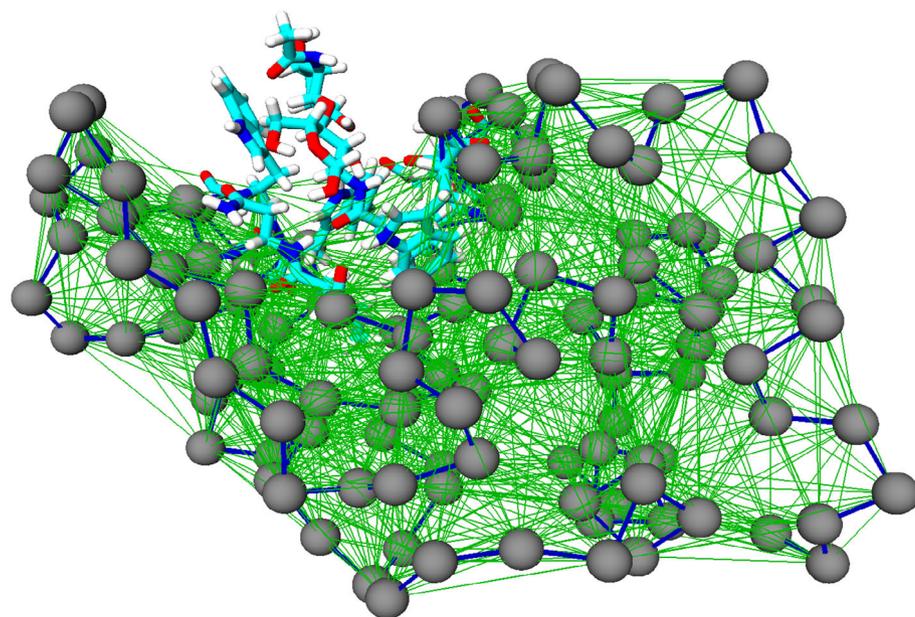


FIGURE 2 Visualization of the dual-resolution protein.¹² The residues included in atomistic detail are shown in red, blue, cyan, and white (O, N, C, and H atoms). The gray spheres are elastic network model nodes, the stiff backbone springs are shown as dark blue lines and all others (weaker) springs are shown in green [Color figure can be viewed at wileyonlinelibrary.com]

of the latter, as well as on the information that the study of this impact can reveal. However, the multiresolution description can, in principle, also provide a valuable computational advantage. In fact, a dual-resolution model can be significantly faster than the equivalent fully AT one. The speedup, which depends primarily on the fraction of atoms retained as such,⁵³ is about 2 for the system investigated here: this value is relatively low, due to the fact that lysozyme, albeit a relevant, nontrivial protein, is still relatively small. In this dual-resolution model, up to 10 residues out of 129 are described at the all-atom level, and the degree of coarse-graining of the low-resolution part is not drastic (one interaction site per residue). Additionally, it has to be kept in mind that a considerable fraction (actually the majority) of the degrees of freedom of the whole setup is due to the water modeled with all-atom detail.

A much more relevant speedup can be achieved in larger systems, for example, high molecular weight proteins, antibodies, or viral capsids, for which lower degrees of detail are allowed in the CG region. The main advantage of a dual-resolution treatment of these macromolecules, possibly in combination with an adaptive resolution model of the solvent, is indeed that the computational gain increases with the system size, that is, precisely for those systems for which an all-atom description becomes challenging.

2.3 | Simulation details

The reference model is given by the 2 ns equilibrated PDB structure 1HEW in the NPT ensemble (the Parrinello-Rahman barostat⁵⁴ with a time constant of 2.0 ps and a pressure of 1 bar was used). Both fully AT and dual-resolution models of HEWL are solvated in water and placed in a cubic simulation box of 7.06 nm side. The force field employed is Amber99SB,⁵⁵ whereas the water model is TIP3P.⁵⁶ The inhibitor, which was always AT, had GLYCAM force field parameters consistent with Amber99SB.⁵⁷ The TI binding free energy calculation consists of three different steps: ΔG_{compl} , $\Delta G_{\text{r-off}}$, ΔG_{lig} :

- 1 The protein-ligand complex free energy (ΔG_{compl}) calculation uses 11 λ values per $\Delta G_{\text{restr-on, c}}$, 5 evenly spaced λ values per $\Delta G_{\text{LJ, c}}$ (with separation 0.20) and 15 λ values per $\Delta G_{\text{coul, c}}$, with 600 ps of simulation per λ in the fully AT case, and 4000 ps in the dual-resolution case to improve the statistics.
- 2 The restraint removal free energy ($\Delta G_{\text{r-off}}$) calculation.
- 3 The ligand solvation free energy (ΔG_{lig}) calculation uses 5 evenly spaced λ values per $\Delta G_{\text{coul, e}}$ (with separation 0.20) and 16 λ values per $\Delta G_{\text{LJ, e}}$, with 600 ps of simulation of each λ -value.

In the TI, we employ the soft-core potential of Reference 58 with parameters $\alpha = 0.5$ and $P = 1.0$ to avoid possible singularities in the Lennard-Jones terms from atoms overlapping during the alchemical change. The temperature is kept constant at 298 K by means of a Langevin thermostat with a friction constant $\gamma = 15 \text{ ps}^{-1}$.

The integration step is 1 fs. The calculation of electrostatic interaction is performed using the reaction field method with a dielectric constant $\epsilon = 80$ and a cutoff of 1.2 nm. These parameters are a good compromise between speed and accuracy, as verified in Reference 59. The SETTLE⁶⁰ and RATTLE⁶¹ algorithms for rigid water and rigid bonds to hydrogen have been used. Each system is prepared using fully AT minimization with steepest descent and 6 ns of equilibration in NVT (for both ligand-free and ligand-bound systems). All simulations (both fully AT and dual-resolution) are carried out with the ESPResSo++ simulation package,^{62,63} in which we have implemented TI (except in case of annihilation, for which all steps are performed in both ESPResSo++ and GROMACS⁶⁴). Some preliminary fully AT equilibration simulations use GROMACS. The error bars shown are calculated using the Student *t* at 95% confidence limit,⁶⁵ via standard deviations obtained using block averaging in which all trajectories are divided into four blocks of equal length.

The parameterization of the dual-resolution model is consistent with the work in Reference 12: the spring constant between consecutive C_{α} nodes along the backbone (k_b) has a stiff value of $5 \times 10^4 \text{ kJ mol}^{-1} \text{ nm}^{-2}$, while all the other ones (k_{nb}) have a value of $160 \text{ kJ mol}^{-1} \text{ nm}^{-2}$, until 1.2 nm as cutoff, parameterized by minimizing the average root mean square error in the C_{α} RMSF. Moreover, a WCA interaction is applied between C_{α} nodes and all solvent molecules' center of mass. In the WCA potential, ϵ has a value of 0.34 kJ mol^{-1} arbitrarily chosen as the value for carbon in the AT force field, and $\sigma_i = R_{g, i} \cdot c$ where $R_{g, i}$ is the radius of gyration of a given residue i where c is the same for all amino acids. The value of c is tuned to give the correct bulk water density of reference for a protein-water system. The c value found is 0.658. Further explanations about c can be found in the Supporting Information. The raw data about the simulations and analyses performed in this work are freely available on the Zenodo repository <https://zenodo.org/record/3665677>.

3 | RESULTS AND DISCUSSION

We performed the calculation of ΔG_{bind} of lysozyme modeled in dual-resolution, varying the number of AT residues constituting the binding site and comparing the results with a fully AT reference simulation. Recall that the binding free energy calculation consists of three steps: restraint removal, ligand ΔG , and ligand-complex ΔG ; of these, only the latter depends on protein resolution, that is, only ΔG_{compl} assumes different values for different numbers of active site residues described at the all-atom level.

As explained in the previous section, the contribution coming from the restraints can be analytically computed and amounts to $\Delta G_{\text{r-off}} = -31.3 \text{ kJ mol}^{-1}$. Likewise, the Coulomb and Lennard-Jones contributions to the ligand free energy ΔG_{lig} are the following:

$$\begin{aligned} \Delta G_{\text{coul, l}} &= -142.8 \pm 1.7 \text{ kJ mol}^{-1} \\ \Delta G_{\text{LJ, l}} &= -9.1 \pm 6.3 \text{ kJ mol}^{-1}. \end{aligned}$$

TABLE 2 The resulting values of the complex free energy (fourth column) and its components (Coulomb, Lennard-Jones, and restraints, respectively, in the first three columns) in fully atomistic system and varying the number of atomistic residues

At res	$\Delta G_{\text{Coul,c}}$	$\Delta G_{\text{LJ,c}}$	$\Delta G_{\text{Restr,on,c}}$	ΔG_{compl}
Fully-at	145.2 ± 3.5	44.2 ± 5.2	3.6 ± 0.4	193.0 ± 9.1
aa-3	125.5 ± 7.0	50.4 ± 6.3	8.3 ± 1.1	184.2 ± 14.4
aa-4	141.4 ± 4.9	39.7 ± 9.4	7.2 ± 1.0	188.3 ± 15.3
aa-5	140.2 ± 2.8	48.7 ± 4.5	7.5 ± 1.2	196.4 ± 8.5
aa-6	147.0 ± 1.9	41.7 ± 5.4	5.1 ± 0.5	193.8 ± 7.8
aa-7	144.5 ± 0.8	38.4 ± 3.8	5.0 ± 0.2	187.9 ± 4.8
aa-8	148.0 ± 1.4	33.6 ± 1.9	6.4 ± 1.8	188.0 ± 5.1
aa-9	143.4 ± 4.7	38.1 ± 5.3	5.1 ± 0.3	186.6 ± 10.3
aa-10	145.9 ± 2.2	38.2 ± 1.0	4.4 ± 0.3	188.5 ± 3.5

Note: All the values are in kJ mol^{-1} and performed with thermodynamic integration. Moreover, all simulations are carried out in ESPResSo++. In particular, for each value of λ , the dual-resolution simulations with different number of atomistic residues last 4 ns; the atomistic simulation, instead, lasts 0.6 ns (600 ps).

Hence,

$$\Delta G_{\text{lig}} = -151.9 \pm 8.0 \text{ kJ mol}^{-1}.$$

The final step is the calculation of ΔG_{compl} , whose results, including the comparison between dual-resolution model and fully AT reference, are shown in Table 2 and illustrated in Figure 3.

The first three columns of the table describe the Coulomb, Lennard-Jones, restraints contributions to free energy, respectively, while the last one corresponds to the value of the total ligand-protein complex free energy. All the values are expressed in kJ mol^{-1} . In Figure 3, the AT reference is represented with a dashed black line with its error bar. In particular, panels A-C show the three components that contribute to the total complex free energy, reported in panel D. Looking at these values as a function of the number of all-atom active site residues, we notice that there are important deviations of the free energy from the reference, especially in the case of 3 and 4 AT residues. On the contrary, the total value of the binding free energy agrees with the reference within the error bar in all cases.

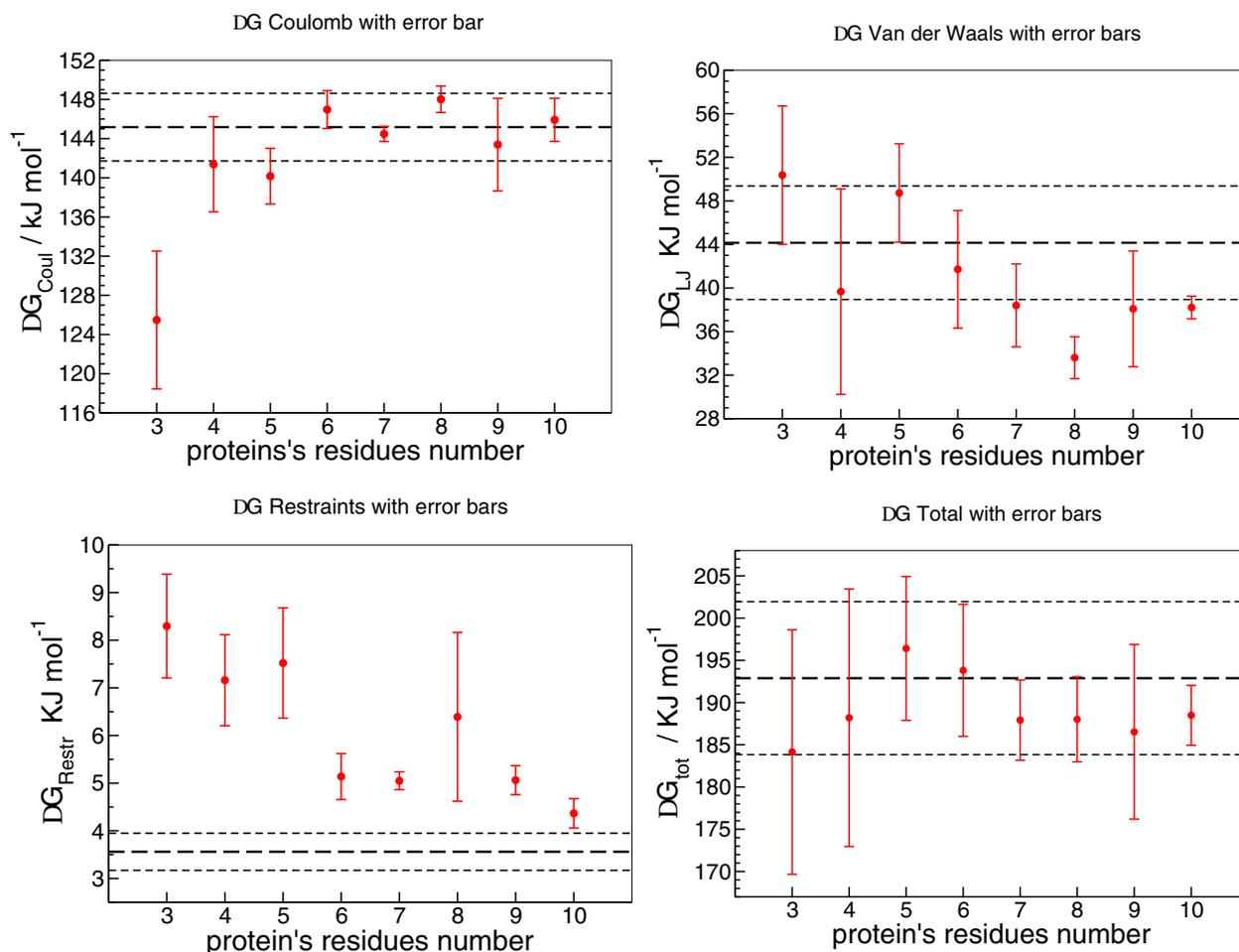


FIGURE 3 A, Coulomb; B, Lennard-Jones; C, restraint; and D, total free energies in the protein-ligand complex, as a function of protein's residues number included in atomistic detail in the multiresolution setup. The heavy dashed black horizontal lines are the reference values from fully atomistic simulations, and the lighter dotted black horizontal lines are the error bars for those values. These simulations use decoupling, not annihilation. Y-axes do not cover the same energy range [Color figure can be viewed at wileyonlinelibrary.com]

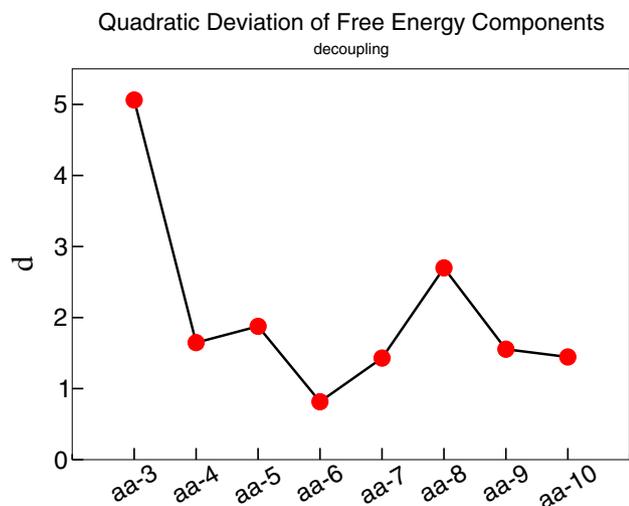


FIGURE 4 Square root of the quadratic deviation δ^2 vs the number of atomistic residues chosen. The plot shows that in the case of six atomistic residues, the value of quadratic deviation is the lowest one and hence it means that such a number leads the best result of free energy. Moreover, the black line shows the trend of free energy values as discussed in Section 3

Furthermore, we observe that the trend of free energy values, in comparison to the reference, is essentially the same: starting from 3 amino acids, it approaches the reference until reaching 6, both in its components and in total. In contrast, going from 6 to 8 AT residues the free energy value deviates from the reference, even though the total remains close to it. Finally, from 8 to 10, ΔG converges again. Hence, increasing the number of AT residues does not introduce necessarily an improvement of the computed free energy, at least as long as the various free energy components are considered separately.

In order to gain further, quantitative insight into these results, we computed the quadratic deviation from the reference, δ^2 , defined as:

$$\delta_i^2 = \delta_{i-Coul}^2 + \delta_{i-LJ}^2 + \delta_{i-Restr}^2 = (\Delta G_{Coul,i} - \Delta G_{Coul-at})^2 + (\Delta G_{LJ,i} - \Delta G_{LJ-at})^2 + (\Delta G_{Restr,i} - \Delta G_{Restr-at})^2, \quad (5)$$

where the index $i = 3, \dots, 10$ runs over AT residues. Figure 4 reports δ^2 as a function of the number of active site amino acids modeled with AT detail.

The plot shows that the binding free energy computed in the dual-resolution model approaches the reference as the number of AT

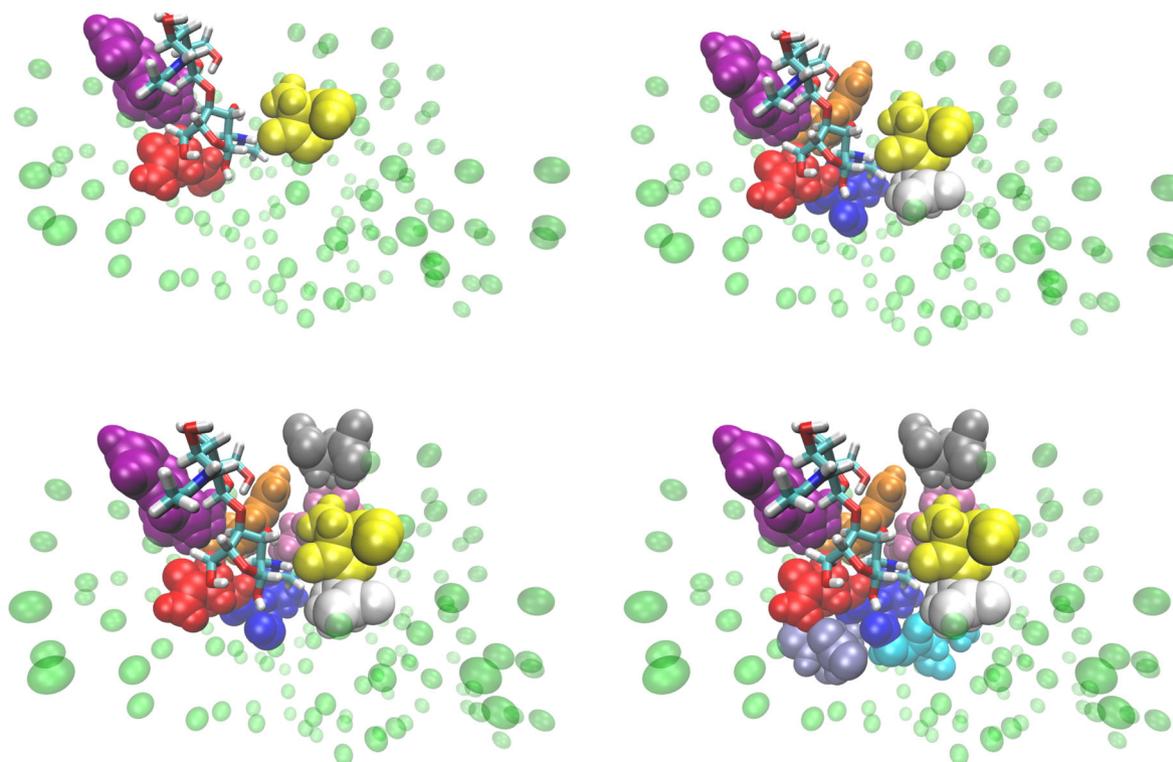


FIGURE 5 VMD representation of lysozyme and ligand in different resolution: A, 3; B, 6; C, 8; and D, 10 atomistic residues. The complete set can be found in Supporting Information. The ligand is always atomistic and it is represented in Licorice. In green are represented the elastic network model beads. With the other colors are represented, instead, the various atomistic residues that surround the ligand [Color figure can be viewed at wileyonlinelibrary.com]

active site residues increases, and most importantly, this trend persists for each component up to 6 residues. Beyond this value, though, the trend stops and the deviation becomes larger, peaking at 8 residues and decreasing when further AT amino acids are added. These results highlight a nonmonotonic dependence of the free energy on the mapping, that is, the number of retained AT residues. If, on the one hand, the overall value of the binding free energy (Figure 3D) levels to the

TABLE 3 Representation of free energies values computed in ESPResSo++ and GROMACS (respectively *espp* and *grom* using a short notation on the table) in case of annihilation and decoupling

	Ligand	Complex	Binding
Annihilation			
<i>atom, espp</i>	-1275.3 ± 11.2	1315.2 ± 16.3	8.6 ± 27.5
<i>atom, grom</i>	-1259.0 ± 5.9	1314.8 ± 13.2	24.5 ± 19.1
Decoupling			
<i>atom, espp</i>	-151.9 ± 8.0	193.0 ± 9.1	9.8 ± 17.1
<i>aa-3, espp</i>	-151.9 ± 8.0	184.2 ± 14.4	1.0 ± 22.4
<i>aa-4, espp</i>	-151.9 ± 8.0	188.3 ± 15.3	5.1 ± 23.3
<i>aa-5, espp</i>	-151.9 ± 8.0	196.4 ± 8.5	13.2 ± 16.5
<i>aa-6, espp</i>	-151.9 ± 8.0	193.8 ± 7.8	10.6 ± 15.8
<i>aa-7, espp</i>	-151.9 ± 8.0	187.9 ± 4.8	4.7 ± 12.8
<i>aa-8, espp</i>	-151.9 ± 8.0	188.0 ± 5.1	4.8 ± 13.1
<i>aa-9, espp</i>	-151.9 ± 8.0	186.6 ± 10.3	3.4 ± 18.3
<i>aa-10, espp</i>	-151.9 ± 8.0	188.5 ± 3.5	5.3 ± 11.5

Note: The table is divided in three column: from left to right are represented the ligand, protein-ligand complex and binding FE. The latter is the algebraic sum of ΔG_{comp} , $\Delta G_{\text{r-off}}$, and ΔG_{lig} . Results are in kJ mol^{-1} .

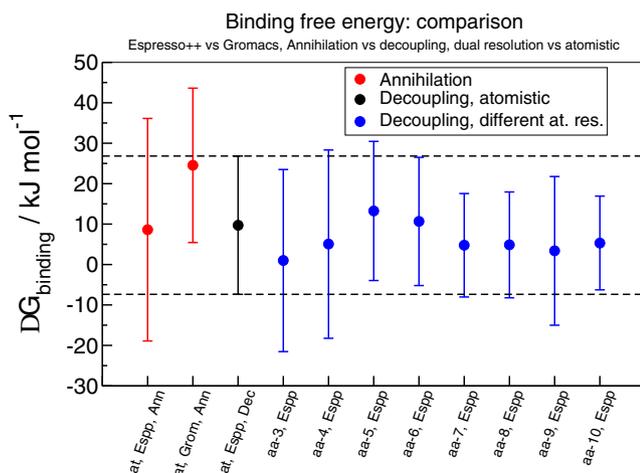


FIGURE 6 Binding free energies as a function of the number of protein residues included in atomistic detail in the multiresolution setup, as well as in a fully atomistic setup. The heavy dashed black horizontal lines and black point are the reference values from fully atomistic simulations obtained in ESPResSo++ with decoupling, and the lighter dotted black horizontal lines are the error bars for those values. Binding free energies values in ESPResSo++ and GROMACS in case of annihilation are represented in red. The binding free energy value in dual resolution simulation changing the number of atomistic residues is represented in blue

reference with as few all-atom residues as 4, the separate components oscillate and reach the plateau only for larger numbers. The existence of a minimum in the standard deviation of all three contributions pinpoints a particular number of AT active site residues for which the accuracy of the computed free energy is the highest and the economy of the high-resolution subpart the largest. Including more than 6 AT residues counterintuitively worsens the results -when the various contributions are looked at- and the previous accuracy is only recovered when more residues are included. This behavior suggests that the total free energy undergoes an error cancellation that hides the deviations of the separate terms.

A possible explanation for this nontrivial behavior is that when 6 active site residues are modeled with all-atom accuracy (Figure 5B) the ligand is stable in the catalytic site, namely it is surrounded by a complete shell of AT residues. The addition or deletion of other residues (Figure 5C,A, respectively) leads to a worsening of ΔG : in the first case, the two added residues (in pink and gray) are located behind the first shell of amino acids (far away from the ligand) and start to form a second, incomplete shell; in the second case, only three AT amino acids take part in the direct interaction with the ligand: therefore, the first layer is still incomplete and important interactions are missing; in order to improve the free energy value one has to add further amino acids in order to complete the second shell. We emphasize that the impact on the deviation from the reference is inversely proportional to the distance of the added/removed amino acid. Thus, the farther the AT amino acid is from the ligand, the more negligible its effect is. In the Supporting Information, we provide detail about the other numbers of all-atom residues not reported here. Finally, the values of binding free energy (also for the case of annihilation whose calculations are reported in the Supporting Information) are summarized in Table 3 and illustrated in Figure 6.

4 | CONCLUSIONS

In this work, we have shown how the dual resolution model employed, constituted by an all-atom subregion coupled to an ENM remainder, can be used to calculate the binding free energy of an enzyme-substrate complex with AT accuracy. Furthermore, and most importantly, we have highlighted the impact that different choices of the model resolution can have. Specifically, we have computed the total value of the binding free energy as well as that of its various energetic components, and quantitatively inspected how these change when different selections are performed for the subgroup of amino acids, ranging from 3 to 10 in total, to be modeled at the fully AT level.

At first sight, one can appreciate that the binding free energy value rapidly converges to the AT reference when as few as 4 amino acids constituting the active site are described all-atom. This comforting result, however, unveils a greater complexity when the different terms constituting the free energy are looked at separately. These show an oscillating behavior as the number of all-atom residues in the active site is increased, with a decreasing difference from the

reference followed by a sudden jump to larger values, which dampens upon further addition of AT amino acids. The rationale in this behavior is identified in the structure of the active site, which is constituted by a first shell of the six residues exposed to the solvent and closest to the ligand; when further amino acids beyond these are modeled with AT resolution, they interact with the substrate affecting the binding free energy components and shifting them away from the reference, with a steadily lowering impact as the model's resolution is increased—as one can expect. Surprisingly, very little if no signal of this behavior is observed in the value of the binding free energy as a whole, rather it becomes visible only upon inspection of its separate contributions.

The results of this work thus highlight the importance of the mapping in the construction of multiscale and multiresolution models, as a higher degree of detail does not necessarily correlate with a higher accuracy of the quantities of interest. The implications of these observations should serve as a warning in the realization of CG models concurrently employing various levels of detail for different regions of the same system, whose range of application spans from fundamental understating of a molecule's properties to real-life pharmaceutical applications.

ACKNOWLEDGMENTS

The authors are grateful to Robinson Cortes-Huerto and Thomas Tarenzi for a critical reading of the manuscript. This work has been supported by the European Research Council under the European Union's Seventh Framework Programme (FP7/2007-2013)/ERC Grant Agreement No. 340906-MOLPROCOMP. This project received funding from the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation program (Grant 758588).

ORCID

Raffaello Potestio  <https://orcid.org/0000-0001-6408-9380>

REFERENCES

- Boyce SE, Mobley DL, Rocklin GJ, Graves AP, Dill KA, Shoichet BK. Predicting ligand binding affinity with alchemical free energy methods in a polar model binding site. *J Mol Biol.* 2009;394:747-763.
- Aldeghi M, Bluck JP, Biggin PC. Absolute alchemical free energy calculations for ligand binding: a beginner's guide. *Comput Drug Discov Des.* 2018;1762:199-232.
- Boresch S, Tettinger F, Leitgeb M, Karplus M. Absolute binding free energies: a quantitative approach for their calculation. *J Phys Chem B.* 2003;107(35):9535-9551.
- Cournia Z, Allen B, Sherman W. Relative binding free energy calculations in drug discovery: recent advances and practical considerations. *J Chem Inf Model.* 2017;57(12):2911, 29243483-2937.
- Abel R, Wang L, Harder ED, Berne BJ, Friesner RA. Advancing drug discovery through enhanced free energy calculations. *Acc Chem Res.* 2017;50(7):1625-1632.
- Dominy BN. Molecular recognition and binding free energy calculations in drug development. *Curr Pharm Biotechnol.* 2008;9(2):87-95.
- Potestio R, Peter C, Kremer K. Computer simulations of soft matter: linking the scales. *Entropy.* 2014;16(8):4199-4245.
- Neri M, Anselmi C, Cascella M, Maritan A, Carloni P. Coarse-grained model of proteins incorporating atomistic detail of the active site. *Phys Rev Lett.* 2005;95:218102.
- Neri M, Baaden M, Carnevale V, Anselmi C, Maritan A, Carloni P. Microseconds dynamics simulations of the outer-membrane protease T. *Biophys J.* 2008;94(1):71-78.
- Machado MR, Dans PD, Pantano S. A hybrid all-atom/coarse grain model for multiscale simulations of DNA. *Phys Chem Chem Phys.* 2011;13:18134-18144.
- Machado MR, Pantano S. Exploring LacI-DNA dynamics by multiscale simulations using the sirah force field. *J Chem Theory Comput.* 2015;11(10):5012-5023.
- Fogarty AC, Potestio R, Kremer K. A multi-resolution model to capture both global fluctuations of an enzyme and molecular recognition in the ligand-binding site. *Proteins Struct Funct Bioinf.* 2016;84(12):1902-1913.
- Amadei A, Linssen ABM, Berendsen HJC. Essential dynamics of proteins. *Proteins Struct Funct Bioinf.* 1993;17(4):412-425.
- Carnevale V, Raugei S, Micheletti C, Carloni P. Convergent dynamics in the protease enzymatic superfamily. *J Am Chem Soc.* 2006;2:173-181.
- Zen A, Carnevale V, Lesk AM, Micheletti C. Correspondences between low-energy modes in enzymes: dynamics-based alignment of enzymatic functional families. *Protein Sci.* 2008;17:918-929.
- Pontiggia F, Zen A, Micheletti C. Small and large scale conformational changes of adenylate kinase: a molecular dynamics study of the sub-domain motion and mechanics. *Biophys J.* 2008;95(12):5901-5912.
- Tirion MM. Large amplitude elastic motions in proteins from a single-parameter, atomic analysis. *Phys Rev Lett.* 1996;77:1905-1908.
- Hinsen K. Analysis of domain motions by approximate normal mode calculations. *Proteins.* 1998;33:417-429.
- Delarue M, Sanejouand YH. Simplified normal mode analysis of conformational transitions in DNA-dependent polymerases: the elastic network model. *J Mol Biol.* 2002;320(5):1011-1024.
- Micheletti C, Carloni P, Maritan A. Accurate and efficient description of protein vibrational dynamics: comparing molecular dynamics and gaussian models. *Proteins.* 2004;55(3):635-645.
- Romo TD, Grossfield A. Validating and improving elastic network models with molecular dynamics simulations. *Proteins Struct Funct Bioinf.* 2011;79(1):23-34.
- Potestio R, Pontiggia F, Micheletti C. Coarse-grained description of proteins' internal dynamics: an optimal strategy for decomposing proteins in rigid subunits. *Biophys J.* 2009;96:4993-5002.
- Gohlke H, Thorpe MF. A natural coarse graining for simulating large biomolecular motion. *Biophys J.* 2006;91:2115-2120.
- Zhang Z, Lu L, Noid WG, Krishna V, Pfaendtner J, Voth GA. A systematic methodology for defining coarse-grained sites in large biomolecules. *Biophys J.* 2008;95(11):5073-5083.
- Zhang Z, Pfaendtner J, Grafmüller A, Voth GA. Defining coarse-grained representations of large biomolecules and biomolecular complexes from elastic network models. *Biophys J.* 2009;97(8):2327-2337.
- Zhang Z, Voth GA. Coarse-grained representations of large biomolecular complexes from low-resolution structural data. *J Chem Theory Comput.* 2010;6(9):2990-3002.
- Sinitskiy AV, Saunders MG, Voth GA. Optimal number of coarse-grained sites in different components of large biomolecular complexes. *J Phys Chem B.* 2012;116(29):8363-8374.
- Polles G, Indelicato G, Potestio R, Cermelli P, Twarock R, Micheletti C. Mechanical and assembly units of viral capsids identified via quasi-rigid domain decomposition. *PLoS Comput Biol.* 2013;9(11):1-13.
- Foley TT, Shell MS, Noid WG. The impact of resolution upon entropy and information in coarse-grained models. *J Chem Phys.* 2015;143(24):243104.
- Patrick Diggins IV, Liu C, Deserno M, Potestio R. Optimal coarse-grained site selection in elastic network models of biomolecules. *J Chem Theory Comput.* 2018;15(1):648-664.

31. Gowers RJ, Carbone P. A multiscale approach to model hydrogen bonding: the case of polyamide. *J Chem Phys*. 2015;142(22):224907.
32. Abrams CF, Delle Site L, Kremer K. Dual-resolution coarse-grained simulation of the bisphenol-*a*-polycarbonate/nickel interface. *Phys Rev E*. 2003;67:021807.
33. Han W, Wan C-K, Jiang F, Wu Y-D. Pace force field for protein simulations. 1. Full parameterization of version 1 and verification. *J Chem Theory Comput*. 2010;6(11):3373-3389.
34. Han W, Wan C-K, Wu Y-D. Pace force field for protein simulations. 2. Folding simulations of peptides. *J Chem Theory Comput*. 2010;6(11):3390-3402.
35. Shelley MY, Selvan ME, Zhao J, et al. A new mixed all-atom/coarse-grained model: application to melittin aggregation in aqueous solution. *J Chem Theory Comput*. 2017;13(8):3881-3897.
36. Kar P, Gopal SM, Cheng Y-M, Predeus A, Feig M. Primo: a transferable coarse-grained force field for proteins. *J Chem Theory Comput*. 2013;9(8):3769-3788.
37. Darré L, Machado MR, Brandner AF, González HC, Ferreira S, Pantano S. Sirah: a structurally unbiased coarse-grained force field for proteins with aqueous solvation and long-range electrostatics. *J Chem Theory Comput*. 2015;11(2):723-739.
38. Orsi M, Essex JW. The elba force field for coarse-grain modeling of lipid membranes. *PLoS One*. 2011;6(12):1-22.
39. Praprotnik M, Delle Site L, Kremer K. Adaptive resolution molecular-dynamics simulation: changing the degrees of freedom on the fly. *J Chem Phys*. 2005;123(22):224106-14.
40. Praprotnik M, Delle Site L, Kremer K. Adaptive resolution scheme for efficient hybrid atomistic-mesoscale molecular dynamics simulations of dense liquids. *Phys Rev E*. 2006;73:066701.
41. Praprotnik M, Delle Site L, Kremer K. A macromolecule in a solvent: adaptive resolution molecular dynamics simulation. *J Chem Phys*. 2007;126:134902.
42. Potestio R, Fritsch S, Español P, et al. Hamiltonian adaptive resolution simulation for molecular liquids. *Phys Rev Lett*. 2013;110:108301.
43. Potestio R, Español P, Delgado-Buscalioni R, Everaers R, Kremer K, Donadio D. Monte carlo adaptive resolution simulation of multicomponent molecular liquids. *Phys Rev Lett*. 2013;111:060601.
44. Fogarty AC, Potestio R, Kremer K. Adaptive resolution simulation of a biomolecule and its hydration shell: structural and dynamical properties. *J Chem Phys*. 2015;142(19):195101.
45. Netz PA, Potestio R, Kremer K. Adaptive resolution simulation of oligonucleotides. *J Chem Phys*. 2016;145(23):234101.
46. Fiorentini R, Kremer K, Potestio R, Fogarty A. Using force-based adaptive resolution simulations to calculate solvation free energies of amino acid sidechain analogues. *J Chem Phys*. 2017;146(06):244113.
47. Heidari M, Kremer K, Cortes-Huerto R, Potestio R. Spatially resolved thermodynamic integration: an efficient method to compute chemical potentials of dense fluids. *J Chem Theory Comput*. 2018;14(7):3409-3417.
48. Kirkwood JG. Statistical mechanics of fluid mixtures. *J Chem Phys*. 1935;3(5):300-313.
49. Baysal C, Atilgan AR. Relaxation kinetics and the glassiness of proteins: the case of bovine pancreatic trypsin inhibitor. *Biophys J*. 2002;83(2):699-705.
50. Frauenfelder H, McMahon B. Dynamics and function of proteins: the search for general concepts. *Proc Natl Acad Sci*. 1998;95(9):4795-4797.
51. Tarenzi T, Calandrini V, Potestio R, Carloni P. Open-boundary molecular mechanics/coarse-grained framework for simulations of low-resolution g-protein-coupled receptor-ligand complexes. *J Chem Theory Comput*. 2019;15(3):2101, 30763087-2109.
52. Weeks JD, Chandler D, Andersen HC. Role of repulsive forces in determining the equilibrium structure of simple liquids. *J Chem Phys*. 1971;54(12):5237-5247.
53. Kreis K, Fogarty AC, Kremer K, Potestio R. Advantages and challenges in coupling an ideal gas to atomistic models in adaptive resolution simulations. *Eur Phys J Spec Top*. 2015;224(12):2289-2304.
54. Parrinello M, Rahman A. Polymorphic transitions in single crystals: a new molecular dynamics method. *J Appl Phys*. 1981;52(12):7182-7190.
55. Hornak V, Abel R, Okur A, Strockbine B, Roitberg A, Simmerling C. Comparison of multiple amber force fields and development of improved protein backbone parameters. *Proteins Struct Funct Bioinf*. 2006;65(3):712-725.
56. Jorgensen WL, Chandrasekhar J, Madura JD, Impey RW, Klein ML. Comparison of simple potential functions for simulating liquid water. *J Chem Phys*. 1983;79(2):926-935.
57. Kirschner KN, Yongye AB, Tschampel SM, et al. Glycam06: a generalizable biomolecular force field. *Carbohydrates. J Comput Chem*. 2008;29(4):622-655.
58. Abraham MJ, Van Der Spoel D, Lindahl E, Hess B. The gromacs development team gromacs user manual version 5.0.4. 2014. <https://ftp.gromacs.org/pub/manual/manual-5.0.4.pdf>.
59. Shirts MR, Pitera JW, Swope WC, Pande VS. Extremely precise free energy calculations of amino acid side chain analogs: comparison of common molecular mechanics force fields for proteins. *J Chem Phys*. 2003;119(11):5740-5761.
60. Miyamoto S, Kollman PA. Settle: an analytical version of the shake and rattle algorithm for rigid water models. *J Comput Chem*. 1992;13(8):952-962.
61. Andersen HC. Rattle: a "velocity" version of the shake algorithm for molecular dynamics calculations. *J Comput Phys*. 1983;52(1):24-34.
62. Halverson JD, Brandes T, Lenz O, et al. Espresso++: a modern multiscale simulation package for soft matter systems. *Comput Phys Commun*. 2013;184:1129-1149.
63. Guzman HV, Tretyakov N, Kobayashi H, et al. Espresso++ 2.0: advanced methods for multiscale molecular simulation. *Comput Phys Commun*. 2019;238:66-76.
64. Hess B, Kutzner C, van der Spoel D, Lindahl E. Gromacs 4: algorithms for highly efficient, load-balanced, and scalable molecular simulation. *J Chem Theory Comput*. 2008;4(3):435-447.
65. Student. The probable error of a mean. *Biometrika*. 1908;6(1):1-25.

SUPPORTING INFORMATION

Additional supporting information may be found online in the Supporting Information section at the end of this article.

How to cite this article: Fiorentini R, Kremer K, Potestio R. Ligand-protein interactions in lysozyme investigated through a dual-resolution model. *Proteins*. 2020;88:1351-1360. <https://doi.org/10.1002/prot.25954>