

If the particle remains inside a thin tube of radius  $r_0$  centred about  $\Gamma$  where  $r_0$  is of the order of the characteristic distance of the field variation, then the particle is said to be confined and  $r_0$  is called the 'plasma radius'.

By taking into account the expansion parameters  $a_0$  (Larmor radius/plasma radius) and  $\varepsilon$  (plasma radius/mean radius of curvature) together, an adiabatic type theory can be developed, which takes finite Larmor radius effects into account in the basic periodic solution of the problem. The expansion parameter is in fact:

$$\varepsilon a_0 = \frac{\text{Larmor radius}}{\text{mean radius of curvature}}$$

In the first-order approximation  $\varepsilon a_0$ , the principal terms which appear only in the second-order approximation of the classical adiabatic theory (parameter  $a_0$ ) are found.

The equations are no longer valid if the particle departs too far from the tube, but this case is not of interest since these particles are no longer confined.

The 'guiding centre' equations obtained are thus valid in the thin tube approximation, even if the Larmor radius is of the same order as the plasma radius.

The effect of singularities in the magnetic field at integral values of the rotational transform on the magnetic axis, can be studied, as well as the reflection of particles with low  $v_{\parallel}$  due to a mirror effect of the toroidal field. Certain cases are completely integrable to this order, and provide a simple illustration of the phenomena.

### Non-linear resonance effects at high power in a cylindrical plasma\*

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A PLASMA column exhibits a main resonance and series of secondary resonances when a low power electromagnetic wave (with  $\vec{E}_{\perp}$  axis and  $\vec{B}_{\parallel}$  axis) is directed on it and when the density is varied by altering the discharge current. The corresponding phenomena are now investigated at high power ( $E$ -fields up to  $10^4$  Vm $^{-1}$ ). A luminosity curve gives an independent measurement of plasma density as a function of discharge current. When  $E > 10^3$  Vm $^{-1}$ , strong non-linear effects appear: deformation of the different resonance peaks and hysteresis, the corresponding luminosity curve exhibits kinks indicating greater perturbation of the average plasma density over a wider resonance domain. As the power increases, the kinks become plateaux indicating preferential absorption of energy at resonance and a tendency of the plasma to remain in a resonant state through h.f. ionization.

With high incident energy, the high amount of energy absorbed at resonance is sufficient to sustain the plasma in a resonant state in the absence of a d.c. discharge current. The absorbed h.f. energy remains practically constant as a function of increasing incoming energy when the plasma is self-sustained in one of its eigenmodes; the corresponding density also remains practically constant as indicated by the luminosity. When the incoming h.f. energy is decreased below a certain minimum, there is a jump to the next resonance characterized by a lower density.

\* Presented by P. E. VANDENPLAS. Paper published in *Phys. Lett.* 25A, 339 (1967).

### Kinetic equations for microscopic turbulence\*

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Two types of kinetic equations for the homogeneous microscopic turbulence of a one-dimensional plasma are compared with each other. These are the usual quasilinear approach and a theory based on the Hamilton-Jacobi formalism. Both theories should yield the same results as long as particle trapping can be neglected. We estimate that this becomes important for times of the order of

$$T = \frac{1}{\omega_p} \cdot \frac{m^2 v_p (\Delta v_p)^3}{e^2 \phi^2}$$

\* Presented by D. PFIRSCH.

here  $v_p$  is a typical phase velocity in a reference system with vanishing group velocity and  $\Delta v_p$  is the range of phase velocities for the unstable modes. For times small compared to  $T$  the two theories agree, whereas for times of the order of  $T$  the Hamilton-Jacobi approach deviates from the quasi-linear theory in a way which can be interpreted as a relaxation effect.

### Instability and anomalous diffusion of a weakly ionized magnetized radio-frequency plasma\*

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THE influence of a magnetic field on the transverse diffusion of a weakly ionized radio-frequency hydrogen plasma has been investigated by measuring the density decay of the plasma along the magnetic field, directed parallel to the discharge tube (inner diameter 5 cm, length 200 cm). The plasma is created at one end of this tube. The axial decay length is inversely proportional to the square root of the transverse diffusion coefficient. The range of pressures used was 0.02–0.07 torr. At low values of the magnetic field ( $\leq 50$ –70 G) the transverse diffusion coefficient behaves classically, while at higher fields the transverse diffusion becomes anomalous, its coefficient being independent of gas pressure and magnetic field. At the critical field  $B_{cr}$ , where anomalous diffusion begins, the plasma becomes unstable, showing oscillations with frequencies of about 100 kc/s and higher harmonics of this frequency. At higher fields the oscillations disappear in strong random fluctuations.  $B_{cr}$  is proportional to the square root of the gas pressure. An explanation of the above phenomena, based on the theory of the ion acoustic instability of a weakly ionized inhomogeneous plasma by Timofeev, can be given, from which the values of the critical field and the oscillation spectrum can be derived. The results of the diffusion coefficient measurements can be explained by introducing the turbulent mixing length concept in the theory of a plasma, subjected to ion acoustic instability.

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### Fusion efficiency for revolving fields\*

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IN A d.c. magnetic bottle a hydrogen plasma has been created by the magnetic and electric fields induced by two sets of four coils. One phase of a two-phase 4.0-MHz oscillator energizes two coils of each set, and the other phase energizes the other two coils. The strength of the magnetic bottle field is that required for ion-cyclotron resonance. Thomson scattering from a ruby laser beam indicates a particle density of  $2 \times 10^{13} \text{ cm}^{-3}$  to date. The salient ability of the coil ensemble to (a) walk ions to the centre of the plasma and (b) to give them an energy corresponding to that in an orbit of a radius which is 5 per cent less than the radius of the vacuum chamber but no more energy, is corroborated.

An extension of this coil ensemble is proposed. Two such ensembles as above would be placed in the same magnetic bottle field and made to operate upon two parallel and continuously connected vacuum chambers. One coil ensemble would be energized at a cyclotron frequency of  $A$  and the other ensemble, properly phased, at a frequency of  $B$ , where  $A/B = 3/2$ . A mixture of deuterium and tritium would be admitted. The two cyclotronic beams would react in the connected regions of the chambers. Calculations indicate that for the two ion beams, each at the terminal energy of 44.5 keV and at regional densities of about  $10^{16} \text{ cm}^{-3}$ , an energy yield could be obtained which is equal to the energy inherent in the beams.

\* Presented by M. L. POOL.

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