Sensitivity of chemical tracers to meteorological parameters in the MOZART-3 chemical transport model

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1 The Model for Ozone and Related Chemical Tracers, version 3 (MOZART-3), which represents the chemical and physical processes from the troposphere through the lower mesosphere, was used to evaluate the representation of long-lived tracers and ozone using three different meteorological fields. The meteorological fields are based on (1) the Whole Atmosphere Community Climate Model, version 1b (WACCM1b), (2) the European Centre for Medium-Range Weather Forecasts (ECMWF) operational analysis, and (3) a new reanalysis for year 2000 from ECMWF called EXP471. Model-derived tracers (methane, water vapor, and total inorganic nitrogen) and ozone are compared to data climatologies from satellites. Model mean age of air was also derived and compared to in situ CO 2 and SF 6 data. A detailed analysis of the chemical fields simulated by MOZART-3 shows that even though the general features characterizing the three dynamical sets are rather similar, slight differences in winds and temperature can produce substantial differences in the calculated distributions of chemical tracers. The MOZART-3 simulations that use meteorological fields from WACCM1b and ECMWF EXP471 represented best the distribution of long-lived tracers and mean age of air in the stratosphere. There was a significant improvement using the ECMWF EXP471 reanalysis data product over the ECMWF operational data product. The effect of the quasi-biennial oscillation circulation on long-lived tracers and ozone is examined.


1. Introduction

During the last decade, three-dimensional (3-D) chemical transport models (CTMs) have been developed and used to quantify the global budget of tropospheric and stratospheric chemical compounds, and to simulate their atmospheric transport, chemical transformations and removal from the atmosphere. These numerical tools have also been used to assess the impact of human activities on global air quality. Many of the earlier chemical transport models that were focused on tropospheric physical and chemical processes, extended in altitude only to the tropopause (typically 100-50 hPa) [e.g., Levy et al., 1985; Crutzen and Zimmerman, 1991; Penner et al., 1991; Roelofs and Lelieveld, 1995; Müller and Brasseur, 1995; Wang et al., 1998] or into the midstratosphere (typically 10-5 hPa) [e.g., Brasseur et al., 1998; Lawrence et al., 1999; Horowitz et al., 2003], and the influence of stratospheric compounds on tropospheric chemistry was expressed by specified upper boundary conditions. For example, in the Model for Ozone and Related Chemical Tracers, version 2 (MOZART-2) [Horowitz et al., 2003], the mixing ratio of species like ozone (O 3 ), inorganic nitrogen oxides (NO x ), and nitric acid (HNO 3 ) were specified in the stratosphere according to climatological values, while in the Harvard GEOSCHEM model [Bey et al., 2001], a net stratospheric production was specified for O 3 , which determined the downward flux of this compound through the tropopause [McLinden et al., 2000]. Several studies [Horowitz et al., 2003; Emmons et al., 2003; Rotman et al., 2004; Gettelman et al., 2004] have shown that the formulation of transport conditions in the vicinity of the tropopause has a large influence on the calculated O 3 concentrations in the troposphere, particularly during winter. Different models provide different quantitative values of the cross-tropopause O 3 flux, and hence of the global budget of tropospheric O 3.
Since the climate-related interactions between chemical composition and radiative processes are particularly strong near the tropopause, and the vertical concentration gradients of radiatively active species are large in this region of the atmosphere, much care must be put in the representation of dynamical and chemical processes in the upper troposphere/lower stratosphere (UTLS) region. For example, the vertical resolution adopted in this transition region should lead to a sufficiently accurate representation of cross tropopause exchanges as well as detrainment of air from convective regions. A major challenge for modern models is to accurately describe the exchange of chemical compounds across the tropopause, which is strongly affected by synoptic-scale events and other small-scale mixing processes [Pan et al., 2004, 2007].

Several recent studies have highlighted the difficulty of accurately representing dynamical and transport processes in the upper troposphere and lower stratosphere. Douglass et al. [1999], for example, have shown that different meteorological fields used to drive chemical transport models can lead to very different distributions of chemical species in the UTLS region. Specifically, they provide different estimates of the impact on the stratosphere of the exhaust released by aircraft engines in the vicinity of the tropopause. Douglass et al. [2003] showed that, although chemical transport models driven by assimilated fields provide realistic short-term simulations, the use of dynamical fields produced by atmospheric general circulation models might be more appropriate for long-term simulations of the stratospheric composition. Schoeberl et al. [2003] showed that the assimilation process introduces additional noise in the meridional motions, leading to excessive vertical dispersion and hence to an overestimation of the Brewer-Dobson circulation in the stratosphere. As a result, the mean age of air in the stratosphere [Hall and Plumb, 1994; Hall et al., 1999] is shifted toward values lower than suggested by the observations. Gettelman and Sobel [2000] also showed that assimilation systems could bias estimates of stratospheric-tropospheric exchange. Recently, van Noije et al. [2004] showed that the enhanced Brewer-Dobson circulation produced by the ERA-40 reanalysis distributed by the European Centre for Middle-Range Weather Forecast (ECMWF) produces an excessive downward flux of O$_3$ across the tropopause. They concluded that simulations could be improved if tropospheric chemical transport models driven by the ECMWF reanalysis, ERA-40, constrain the O$_3$ concentrations in the lower stratosphere by relaxation to climatology (rather than by specifying the net stratospheric ozone source), and use the “forecast” fields rather than the “analysis” (first-guess) fields to calculate the transport of chemical species.

The mean meridional circulation of the lower stratosphere, which affects vertical mass fluxes in the UTLS region, is driven primarily by the “drag” provided by planetary and gravity wave momentum deposition in the stratosphere and mesosphere. The “downward control principle” [Haynes et al., 1991] states that the mean vertical velocity in the vicinity of the tropopause is influenced by momentum sources in the stratosphere, and even up to the lower thermosphere. Therefore global three-dimensional chemical transport models, even if they address issues related to tropospheric processes, should extend to the upper levels of the middle atmosphere. This is a condition for these chemical transport models to adequately represent the meridional transport of atmospheric tracers, the mean age of air in the lower stratosphere, and the flux of chemical constituents through the tropopause [Holton et al., 1995]. Such models need also to treat accurately chemical transformations and, specifically, heterogeneous processes on polar stratospheric cloud (PSC) and aerosol particles.

Three-dimensional chemical transport models with a focus on the middle atmosphere were developed over the last decade [e.g., Rose and Brasseur, 1989; Chipperfield et al., 1993; Rasch et al., 1994; Lefèvre et al., 1994; Brasseur et al., 1997; Douglass et al., 1997; Rotman et al., 2001; Mahowald et al., 2002; Marsh et al., 2003; Rotman et al., 2004; Considine, 2004; Douglass et al., 2004]. These chemical transport models have included gas-phase and heterogeneous reactions that represent realistically the distribution of radicals contained in the O$_3$, HO$_3$, NO$_3$, ClO$_3$, and BrO$_3$ family of constituents that are important for the O$_3$ balance in this region. Only one of these models represented the chemical and physical processes of both the troposphere and stratosphere [Rotman et al., 2004], while another represented in detail the chemical and physical processes of the mesosphere and lower thermosphere [Rose and Brasseur, 1989; Marsh et al., 2003]. Chemical models covering the “entire atmosphere” have recently been coupled “online” to general circulation models. The CTM used in the present study, the Model for Ozone and Related Chemical Tracers, version 3 (MOZART-3) chemical formulation has, for example, been incorporated into the Whole Atmosphere Community Climate Model (WACCM) [Sassi et al., 2005; Lamarque et al., 2006; Eyring et al., 2006; Garcia et al., 2007; Marsh et al., 2007; Tilmes et al., 2007] at the National Center for Atmospheric Research, in Boulder, Colorado, and into the Hamburg Model for Neutral and Ionized Atmosphere (HAMMONIA) [Schmidt et al., 2006] at the Max Planck Institute for Meteorology, in Hamburg, Germany.

Several application studies using MOZART-3 have recently been published [Forkman et al., 2003; Park et al., 2004; Sassi et al., 2004; Gettelman et al., 2004; Kulawik et al., 2006]. The purpose of the present paper is (1) to provide the first detailed description of MOZART-3 (see auxiliary material), (2) to evaluate the distributions of chemical compounds calculated using different meteorological fields, with an emphasis on the middle atmosphere region, (3) to highlight the role including a quasi-biennial oscillation has in modifying stratospheric ozone and its precursors, and (4) to assess whether the new reanalysis meteorological products provided by ECMWF are suitable for future studies that include satellite observations to better understand stratospheric physical, dynamical, and chemical processes. 1 Specifically, this paper investigates the sensitivity of the calculated tracer distributions and mean age of air to the adopted temperature and wind fields, and assesses the impact of different physical and dynamical processes on the distribution of water vapor ($H_2O$) and O$_3$ in the middle atmosphere.

1Auxiliary material data sets are available at ftp://ftp.agu.org/apend/jd/ 2006jd007879. Other auxiliary material files are in the HTML.
2. Model Description

[8] The MOZART model, version 3 (MOZART-3), was designed to represent chemical and physical processes from the troposphere through the lower thermosphere. This model is an extension to the middle atmosphere of the global tropospheric models MOZART-1 [Brasseur et al., 1998; Hauglustaine et al., 1998] and MOZART-2 [Horowitz et al., 2003; Emmons et al., 2003]. MOZART-3 is designed in a framework similar to the earlier versions of MOZART, but accounts for additional processes such as (1) vertical mixing associated with gravity wave breaking in the upper stratosphere/mesosphere, (2) molecular diffusion of constituents important above 80 km, (3) source gases, radicals, chemical and photochemical reactions to represent the ClO, and BrO families of constituents, (4) stratospheric sulfate aerosols and polar stratospheric clouds (PSCs) to represent stratospheric heterogeneous processes and subsequent O3 loss, (5) photolysis to wavelengths as short as 120 nm, important for upper stratosphere and mesosphere photochemistry, and (6) auroral NOx production for better representation of high-latitude NOx and O3 photochemistry.

[9] MOZART-3 is similar to MOZART-2 [Horowitz et al., 2003] in that a detailed formulation of physical processes in the troposphere is included (e.g., boundary layer exchanges, convective transport, and hydrological cycle) (based on the Model for Atmospheric Transport and Chemistry (MATCH) [Rasch et al., 1997]). Meteorological fields (e.g., temperature and horizontal wind components) are needed to drive the transport of chemical compounds as well as the hydrological cycle (cloudiness, precipitation, etc.; see auxiliary material). These dynamical fields are provided either from a general circulation model simulation or from meteorological analyses produced through operational weather prediction or reanalysis.

[10] The vertical velocities are derived in MOZART-3 from the divergence of the horizontal velocity fields, using an advection routine based on the work of Lin and Rood [1996]. This routine is a numerically fast, flux form semi-Lagrangian transport scheme. Therefore the vertical velocities of the parent General Circulation Model (GCM) are not used. One consequence of using meteorological fields from GCMs is that the time evolution of the divergence of the meteorological fields may not be consistent with the time evolution of the surface pressure field. This issue is discussed in detail by Horowitz et al. [2003] and Rotman et al. [2004]. For the results shown in this paper, the error associated with adding in a surface pressure adjustment on the mass conservation of stratospheric consistsents (e.g., ozone) is negligible.

[11] A detailed description of the approaches adopted in MOZART-3 to account for meteorological and transport processes, surface emission and deposition of chemical species, gas-phase and heterogeneous processes, as well as the numerical algorithms used to solve the corresponding continuity transport equations is contained in the auxiliary material.

3. Model Simulations

[12] Three sets of meteorological fields are used to drive MOZART-3 in this study. Below is a brief description of each set and of how the fields were used in the simulations reported here. See also Table 1.

3.1. Case 1: Dynamical Fields From the WACCM1b Model

[13] The Whole Atmosphere Community Climate Model, version 1b (WACCM1b) is a comprehensive model, spanning the range of altitudes from the Earth’s surface to the lower thermosphere [Sassi et al., 2004]. This model is built upon the numerical framework of the NCAR Community Climate System Model (CCSM). The horizontal resolution is $2.8^\circ \times 2.8^\circ$, with 128 x 64 points in a quasi-linear grid [Williamson, 1997]. There are 66 vertical levels from the surface to approximately 150 km. The vertical resolution of WACCM1b is less than 1 km in the troposphere, and approximately 1.2 km in the UTLS region; it increases to 2–3 km in the stratosphere and to ~3 km in the mesosphere/lower thermosphere (M/LT). The dynamical governing equations are solved with a semi-Lagrangian technique [Williamson and Olson, 1994].

3.2. Case 2: Dynamical Fields From ECMWF Operational Analysis

[14] Operational meteorological fields for the year 2000 produced by the European Centre for Medium-Range Weather Forecasts (ECMWF; http://www.ecmwf.int/) are used in this study. Like WACCM1b, the assimilation model used to derive these meteorological fields is based on a semi-Lagrangian treatment of atmospheric dynamics and a comprehensive representation of physical processes. The operational analysis system uses four-dimensional variational (4D-Var) assimilation. During 2000, this 4D-Var system was changed from 6- to 12-hour cycling in September and spectral horizontal resolution was changed from
T319 to T511 truncation in November. For use in MOZART-3, the archived meteorological fields are reduced to a uniform lower resolution by truncation to T63, consistent with an N48 Gaussian grid of horizontal resolution 1.9° × 1.9°. They are defined on 60 levels from the surface to approximately 65 km, with vertical resolution about 0.5 km in the midtroposphere, around 1 km in the UTLS region, and close to 1.5 km for most of the stratosphere.

3.3. Case 3: Dynamical Fields From the ECMWF EXP471 Reanalysis

[15] The ECMWF reanalysis called ERA-40 was produced by a 3D-Var data assimilation technique using 6-hourly cycling. The assimilating atmospheric model had T159 spectral truncation in the horizontal (using a linear reduced Gaussian grid with about 125 km spacing), and the same 60-level vertical resolution as used in the operational analysis system. The model and the assimilation system are described in detail by Uppala et al. [2005].

[16] Because of known deficiencies in ERA-40 fields for simulation of long-term chemical transport [e.g., van Noije et al., 2004; Uppala et al., 2005], we here use an updated set of ECMWF reanalysis fields for the year 2000. The updated reanalysis, which is designated EXP471, is one of several carried out for particular time segments as part of the development of an improved assimilation system to be used for a new reanalysis of the period from 1989 onward. The assimilation system for EXP471 used the same resolution as ERA-40, but used 4D-Var rather than 3D-Var analyses, with 12-hourly rather than 6-hourly cycling. It benefited also from many changes developed for the operational ECMWF system over the period since 2001, when the ERA-40 system was fixed. These changes include revisions in the amount, type, bias correction and quality control of ingested data, in the humidity analysis, in the dynamical background error constraint, and in the assimilating model’s parameterizations of convection and radiation. It was not feasible to evaluate the changes individually over test periods long enough to quantify the separate impact of each on key features such as the implied age of stratospheric air. Monge-Sanz et al. [2007] discuss the changes most likely to have had significant effect. As done for the operational analyses, the reanalysis data have been truncated to T63 and evaluated on the corresponding N48 grid.

3.4. Simulations

[17] The WACCM1b model was integrated from 1979 to 2001 using observed sea surface temperatures and greenhouse gas abundances [Sassi et al., 2004]. High-frequency (3-hour average) meteorological output from this simulation was created for MOZART-3. The WACCM1b year 1999 was chosen for this study because it had meteorological conditions similar to the ECMWF year 2000 for the stratosphere in Northern Hemisphere winter/spring. MOZART-3 was integrated for multiple years (using repeatedly this meteorological year) until the model was in steady state with respect to stratospheric O3. The last year of the simulation was used in this study. In what follows, this MOZART-3 WACCM1b simulation will be designated as MZ3-WA1b (case 1).

[18] The MOZART-3 ECMWF operational and ECMWF EXP471 simulations were integrated in a similar manner. Input meteorology was on a 6-hour average interval. The MOZART-3 model was integrated for multiple years for each set of meteorological fields, until the derived O3 reached steady state in the stratosphere. The last year of the simulation was used in this study. In the following discussion, these simulations will be designated as MZ3-Op (case 2) and MZ3-EXP471 (case 3). MZ3-ECMWF will refer to both of the simulations.

[19] For all three cases examined in the paper, the sulfate surface area density was derived from year 1999 observations [Thomason et al., 1997] (as updated by World Meteorological Organization (WMO) [2003]). Surface boundary conditions for CO2, CH4, N2O, H2, CH3Cl, CCl4, CH2CCL3, CFC12, CF2Cl2, CFC-113, HCFC-22, CH3Cl, CH4Br, CF3Br, and CF2ClBr are constrained by observations (see auxiliary material for references). These surface boundary conditions were set to 1999 conditions for the MZ3-WA1b case and 2000 conditions for the MZ3-ECMWF cases.

[20] For the mean age-of-air simulations, a tracer version of MOZART-3 was created. Here a flux is specified that increases the atmospheric burden linearly in time. This tracer is conserved (i.e., there are no loss processes in the atmosphere). The behavior of such a tracer is similar to the observed temporal evolution of SF6. The mean age of air in the middle atmosphere was derived relative to the tropical tropopause. This procedure is discussed in more detail by García et al. [2007]. The age of air was integrated with the meteorological fields corresponding to the 3 model cases until steady state conditions were reached (approximately 15 years).

3.5. Observations

[21] Three observational data sets are used in this study for comparison to the model simulations. For CH4, H2O, NOx, and O3 we use a climatology derived from instruments on board the Upper Atmosphere Research Satellite (UARS). For comparisons to this satellite climatology, we are only interested in how representative the MOZART-3 simulation results are with respect to climatological means (i.e., we do not attempt to compare specific features in the simulation to satellite products on a given date). We also examine the evolution of O3 and H2O from an occultation instrument that samples high-latitude regions in both hemispheres. Here we give both the climatological mean (for comparison to MZ3-WA1b) and the results from year 2000 (to compare with the MZ3-ECMWF cases). We also examine the evolution of total column ozone. In this case we are primarily interested in how the MZ3-ECMWF simulations agree with total column ozone derived for year 2000. Below is a brief description of each data set.

3.5.1. POAM III: Ozone and Water Vapor

[22] The Polar Ozone and Aerosol Measurement (POAM) III instrument [Lucke et al., 1999] was launched into a polar, Sun synchronous orbit in March of 1998, and operated until December of 2005. POAM III measured vertical profiles of O3, NO2, H2O, and aerosol extinction using the solar occultation technique with nine wavelength channels ranging from 353 nm to 1018 nm. Because of its high-inclination orbit, measurements could be obtained in the polar regions throughout the year. Each day, 14 occultations separated by about 25 degrees in longitude were made around a circle of approximately constant latitude. Latitudes
varied slowly throughout the year, between about 54° and 71° in the Northern Hemisphere, and 63° and 68° in the Southern Hemisphere. (A map of latitude points where observations are made is included in Figure 14.) There are 43,404 total occultations incorporated in the 1998–2005 climatology used in this study. The POAM III inversion algorithm, which utilizes the optimal estimation technique to simultaneously retrieve vertical profiles of all of the measured constituents, has been described in detail by Lumpe et al. [2002].

[25] Lumpe et al. [2006] validated the POAM III v4.0 H2O retrievals. The H2O profiles extend from 5 to 50 km, with a vertical resolution of 1 km in the lower stratosphere, 3 km in the upper stratosphere, and 5–7% precision throughout the stratosphere. Comparisons with correlative measurements performed by Lumpe et al. [2006] suggest a high bias in the Northern Hemisphere POAM measurements (satellite sunrise occultations) of about 5–10% in the altitude range of 12–35 km. The Southern Hemisphere measurements (satellite sunset occultations) exhibit a high bias of about 15–25% in the same range of altitude. Differences between POAM measurements and correlative data decrease at higher altitudes. In the Northern Hemisphere, POAM agrees with other measurements at 40 km, and is 10% low by 50 km. The high bias in the Southern Hemisphere decreases to about 10% by 50 km. In spite of these systematic differences in the correlative comparisons, the POAM H2O measurements were found to be self-consistent, with no long-term trends in accuracy or precision [Lumpe et al., 2006].

[24] Randall et al. [2003] validated the POAM III v3.0 O3 retrievals, showing that the POAM data agree with correlative measurements to within ±5% from 13–60 km. The O3 retrievals exhibit a high bias compared to correlative measurements of up to ~0.1 ppmv from 10–12 km. POAM III O3 retrievals have a vertical resolution of 1 km from 15 to 50 km, increasing to 1.5 km at an altitude of 10 km and 2.5 km at an altitude of 60 km. Random errors are approximately ±5% or less above 12 km. As described by Lumpe et al. [2006], the primary improvements in the version 4.0 retrievals used here over version 3.0 involve changes in how the instrument pointing was calculated, and in the parameterization of the spectral dependence of aerosol extinction. These improvements had only very small effects on the POAM III O3 retrievals, however, and the conclusions of Randall et al. [2003] apply to version 4.0.

3.5.2. EP/TOMS: Total Column Ozone

[25] The Earth Probe Total Ozone Mapping Spectrometer (EP/TOMS) is currently operating in a 740 km, Sun synchronous orbit. The inclination is 98.4° with a period of 99.65 minutes. The equatorial crossing time is 11:16 AM. The EP/TOMS field of view at nadir is 39 km latitude by 39 km longitude. The current orbital characteristics of EP/TOMS give 90% daily global coverage. Data used in this study are taken from the NASA Goddard Space Flight Centers web page http://toms.gsfc.nasa.gov/epftoms/ep.html). Daily zonal mean total column ozone data (in Dobson Units) were averaged into 5° latitude bins for year 2000 by the EP/TOMS science team. The total ozone retrieval gives an absolute error of ±3 percent; the random error is approximately ±2 percent. More detailed information on accuracy, precision, and coverage for EP/TOMS is given by McPeters et al. [1998].

3.5.3. UARS Climatology: CH4, H2O, NOy, and O3

[26] Details regarding the UARS climatology used in this study are given by Randel et al. [1998]. These data were mapped onto an equivalent latitude coordinate. The monthly mean methane (CH4) and H2O climatologies are based primarily on version 18 Halogen Occultation Experiment (HALOE) data between 1991 and 1997. To fill in the polar regions not sampled by HALOE, the Cryogenic Limb Array Etalon Spectrometer (CLAES) version 7 data and the Microwave Limb Sounder (MLS) version 4 data were used for CH4 and H2O, respectively. Data from these two instruments were obtained during 1992–1993. Ozone data were primarily taken from HALOE, with MLS data for 1991–1997 used in the polar regions. Total inorganic odd nitrogen (NOy) from observations are defined by summing the NO + NO2 sunset data observed by HALOE and adding it to the HNO3 data observed by CLAES. The definition of NOy in the model is slightly different, as explained in section 4.2.

4. Model Results and Discussion

[27] The evaluation of MOZART-3 driven with the above mentioned meteorological fields is divided into four sections. The first section examines similarities and differences between the input zonal winds and temperature fields. The second section compares modeled and observed long-lived tracers as surrogates for the global circulation and transport processes. The third section examines model representation of total column ozone and polar processes. The goal of these sections is to show that accurate representation of temperature, dynamics, and transport is critical for achieving realistic chemical distributions, and specifically for obtaining correct distributions of long-lived tracers and of ozone.

4.1. Zonal Winds and Temperature Comparisons

[28] In Figure 1, the latitude-altitude monthly mean zonal wind distributions are shown for January, March, July, and October. Solid lines represent zonal wind propagation from the west (westerlies); dotted lines represent zonal wind propagation from the east (easterlies). In January the WACCM1b zonal winds are similar to the ECMWF zonal winds, except for the presence in the ECMWF data of a westerly jet in the tropics, centered at 50 hPa. It is well known that the zonal winds in the lower stratosphere shift direction from easterlies to westerlies with a period of approximately 27 months, a feature known as the quasi-biennial oscillation (QBO). The year 2000 ECMWF zonal winds show a clear QBO westerly structure in the tropical stratosphere (as well as an enhancement of the easterly winds near 10 mb). The WACCM1b zonal winds are lacking these features because the QBO is not simulated in the parent WACCM1b model. Instead, the tropical stratosphere in the WACCM1b simulation is in a permanent easterly phase. The importance of the QBO on tracer distributions will be examined in section 4.2. In March, there are differences between the Op and EXP471 zonal winds near 1 hPa. Here the EXP471 zonal winds are westerly, whereas the Op zonal winds are easterly. This region of the tropical stratosphere is influenced by a semiannual oscillation (SAO) in the zonal winds, and
Figure 1. Latitude-altitude cross sections of zonal mean zonal winds (m s$^{-1}$) for January, March, July, and October for the ECMWF Operational, ECMWF EXP471, and WACCM1b meteorological fields.
WACCM1b has a SAO signature similar to EXP471. WACCM1b zonal winds in the winter hemisphere are generally stronger than in the ECMWF meteorological fields; this is especially true in the Southern Hemisphere winter (July) and in the transition from winter to spring (October), which occurs too late in the model. Thus WACCM1b retains westerlies through 0.1 hPa, while in the ECMWF meteorological fields there are easterlies above 1 hPa. The importance of this bias in WACCM1b meteorology can be seen in the Southern Hemisphere temperature distribution.

[29] In Figure 2 the zonal monthly mean temperatures are shown for the three MOZART-3 simulations during January, March, July, and October. Year 2000 was a cold year in the Northern Hemisphere [Sabutis and Manney, 2000], with minimum ECMWF temperature approaching 190 K in January. Results from both Op and EXP471 are similar in this figure. The WACCM1b temperatures for January, in the polar Northern Hemisphere are also similar to ECMWF results. At these temperatures in the polar lower stratosphere, one would expect significant activation of inorganic chlorine reservoir species. In March, the WACCM1b zonal monthly mean Northern Hemisphere polar temperatures are 5 K colder than the ECMWF meteorological products in the polar lower stratosphere. This promotes more springtime O\textsubscript{3} loss in the MZ3-WA1b simulation relative to the two MZ3-ECMWF simulations. In July, the ECMWF and WACCM1b temperatures approach 180 K in the Southern Hemisphere polar lower stratosphere. These cold temperatures extend lower in pressure (higher in altitude) in the WACCM1b fields relative to the ECMWF fields; this will have consequences for dehydration between the simulations. In October, the Southern Hemisphere temperatures in WACCM1b in the polar lower stratosphere are colder and this cold region has a vertical extent that is broader in altitude than the ECMWF meteorological products. This will tend to promote more O\textsubscript{3} loss over a larger vertical area in the MZ3-WA1b simulation relative to MZ3-Op and MZ3-EXP471.

[30] The evolution of the daily zonal mean temperatures in the polar lower stratosphere and equatorial UTLS region are shown in Figure 3. At 84°S there are noticeable differences between EXP471 and Op temperatures, with EXP471 being 3–4 K colder in winter/spring. Temperatures from EXP471 are more similar to WACCM1b in this region up to mid-October, when there is a final warming in EXP471, but not in WACCM1b. In the latter case, the Southern Hemisphere vortex stays isolated late into the spring, which causes the O\textsubscript{3} distribution in MZ3-WA1b simulation to be different from those in the MZ3-ECMWF simulations. The cold bias in temperature and isolation of the Southern Hemisphere vortex in WACCM1b is consistent with the zonal wind distributions shown in Figure 1. At 84°N, the Op and EXP471 temperatures are very similar throughout the year. The ECMWF temperatures are also fairly consistent with WACCM1b temperatures until late winter/early spring when the ECMWF temperatures become warmer than the WACCM1b temperatures. In the tropics, at the cold point (85 hPa and 96 hPa for WACCM1b and ECMWF respectively), January and February differ among the Op, EXP471, and WACCM1b cases, with WACCM1b being the coldest and EXP471 being the warmest. WACCM1b also has a larger seasonal amplitude in temperature than the ECMWF temperature products. The most significant difference between WACCM1b and the ECMWF temperatures are seen in November and December, when both ECMWF temperature products have temperature distributions that are colder by up to 4 K relative to WACCM1b. This difference in temperatures will be seen in the evolution of the H\textsubscript{2}O in section 4.2.

[31] It should be noted that in the following sections the comparisons of calculated chemical fields for specific meteorological years are compared to various observational chemical climatologies. Agreement or disagreement between model results and observational climatologies in a broad sense can be instructive, although specific features in either the climatology or model results should not be overinterpreted. This approach can help discriminate between the effects of the different dynamical fields on chemical distributions. An example of how this analysis approach has been successful is given by Douglass et al. [1999].

4.2. Global Circulation and Transport

[32] One measure of the representation of the global circulation and transport is the mean age of air derived in a model [Hall and Plumb, 1994; Hall et al., 1999; Waugh and Hall, 2002]. In Figure 4 (top), the annual average mean age of air is shown for the Op, EXP471, and WACCM1b meteorological fields. Figure 4 (bottom) shows the difference between the Op and WACCM1b cases relative to the EXP471 case. Corresponding comparisons for the Op and EXP471 analysis, and for ERA-40 and UK Met Office analysis, using the TOMCAT transport model have been reported by Monge-Sanz et al. [2007]. The mean age of air from the MZ3-Op simulation is significantly younger than both the MZ3-EXP471 and MZ3-WA1b simulations. The maximum mean age of air for the MZ3-Op simulation is less than 4 years. The maximum mean age of air for the MZ3-EXP471 and MZ3-WA1b simulations is 4.0 and 4.5 years, respectively. The shapes of the mean age-of-air isopleths are also very different for each meteorological product. The derived MZ3-Op mean age of air shows weak meridional gradients throughout the lower stratosphere. In contrast, the derived MZ3-EXP471 mean age of air has strong gradients in the lower stratosphere at 20°N and 20°S. Here the mean age-of-air isopleths are essentially constant in the MZ3-EXP471 simulation at altitudes between 18–28 km, suggesting a strong lower stratosphere midlatitude barrier to horizontal transport between the tropics and subtropics. The derived mean age of air for the MZ3-WA1b simulation does not show such a strong apparent barrier to transport in this region. However, the annual average gradients are larger in the ±45–60° and 10–15 km region in MZ3-WA1b relative to both MZ3-ECMWF simulations. In the tropics, between 30 and 40 km, the mean age-of-air isopleths in the MZ3-WA1b simulation have strong meridional gradients, suggestive of a more isolated tropical region in this simulation. This is consistent with long-lived tracer distributions discussed below. In Figure 5a, a mean age-of-air latitudinal cross section is shown at 20 km for both models and observations [Boering et al., 1996; Andrews et al., 2001; Elkins et al., 1996; Ray et al., 1999; see also Monge-Sanz et al., 2007]. The observational estimates of the mean age of air are based on measurements...
Figure 2. Latitude-altitude cross sections of zonal mean temperature (K) for January, March, July, and October for the ECMWF Operational, ECMWF EXP471, and WACCM1b meteorological fields. Contour intervals are 180–200 K by units of 5 K and 200–300 K by units of 10 K.
of CO$_2$ and SF$_6$. Overall, the MZ3-WA1b simulation better represents the observations at 20 km. In Figure 5b, the mean age-of-air $7^\circ$S profiles are shown. At $7^\circ$S, the MZ3-EXP471 simulation better represents the observations between 17 and 22 km. At $7^\circ$S, all three model simulations greatly underestimate the mean age of air above 22 km. Even at $12^\circ$S (not shown), the three simulations underestimate the observed mean age of air at $7^\circ$S (above 22 km). Therefore,

![Graphs](image_url)

**Figure 3.** Annual cycle of zonal mean temperatures (K) at 43 hPa for ECMWF operational (blue), ECMWF EXP471 (red), and WACCM1b (black) meteorological fields (a) for $84^\circ$S, (b) at the tropical cold point temperature (85 hPa for WACCM1b and 96 hPa for ECMWF meteorological fields), and (c) for $84^\circ$N.

![Graphs](image_url)

**Figure 4.** Latitude-altitude cross sections of the annual average mean age of air (in years) for (a) MZ3-Op, (b) MZ3-EXP471, and (c) MZ3-WA1b. The mean age-of-air isopleths are separated by 0.5 years. (bottom) The difference between the MZ3-Op and MZ3-WA1b simulations relative to the MZ3-EXP471 simulation. For this comparison, the mean age-of-air isopleths are separated by $\pm0.2$ years.
above 22 km, the tropics are younger in a broader region than observations would suggest. At 24°S (not shown), the three simulations have air that is significantly older; however, only the MZ3-WA1b simulation mean age of air approaches the mean age of air derived from observations at 7°S. The model dynamics and transport that leads to having too young mean age of air will be discussed in more detail below.

[33] One diagnostic of the effect of input dynamics and temperature on model derived chemical constituents is the photochemical lifetime. In Table 2, the model photochemical lifetime with respect to loss processes are shown for CH4, N2O, CFC-11, CFC-12, CFC-113, HCFC-22, CCl3, CH3CCl3, CH3Cl, and CH3Br. The best estimates of the lifetimes inferred from observations are also listed for each species. In this study, for molecules that have a strong hydroxyl radical loss (e.g., CH4), the MZ3-EXP471 simulation gives the best result, followed by the MZ3-WA1b and MZ3-Op simulations. This would suggest the tropospheric OH abundance is best represented in MOZART-3 when the EXP471 meteorological fields are used. For species whose loss is dependent on stratospheric photolysis and/or loss with O(1D), the MZ3-WA1b simulation give the best comparison to lifetimes inferred from observations, followed by MZ3-EXP471 and MZ3-Op. This result is consistent with the mean age of air being older in MZ3-WA1b simulation relative to the other two simulations. Using lifetimes analysis as a diagnostic for meteorological field performance should be considered as only a rough metric, since there is considerable uncertainty in the results inferred from observations.

[34] In Figure 6, the tropical time-height propagation of H2O in the equatorial region (averaged between 12°S and 12°N) is shown. The MZ3-WA1b simulation has strong propagation of dry and wet regimes up to the midstratosphere. This propagation pattern is known as the H2O “tape recorder signal”. The MZ3-WA1b wet phase of the tape recorder signal is too moist compared to observations. This has been discussed by Park et al. [2004] and Gettelman et al. [2004]. The MZ3-EXP471 and MZ3-Op simulations are too dry relative to observations based on the HALOE climatology and they do not have the correct seasonal evolution at the equatorial tropopause. The difference in the seasonal evolution of H2O can be related to the seasonal evolution of cold point temperatures (Figure 3b). In the MZ3-ECMWF simulations, the minimum in temperature at the cold point occurs in November and December. It is well known that the year 2000 cold point temperatures during November and December are colder than the climatological mean [Randel et al., 2004]. Therefore the dry phase of the tape recorder signal starts earlier than the UARS climatology in both MZ3-ECMWF simulations. The wet phase (near 100 hPa) of the tape recorder is similar to the HALOE

### Table 2. Global Annual Average Photochemical Lifetime for Selected Species

<table>
<thead>
<tr>
<th>Species</th>
<th>Inferred From Observations</th>
<th>MZ3-Op</th>
<th>MZ3-EXP471</th>
<th>MZ3-WA1b</th>
</tr>
</thead>
<tbody>
<tr>
<td>CH4</td>
<td>8.9°</td>
<td>10.0</td>
<td>8.7</td>
<td>9.4</td>
</tr>
<tr>
<td>N2O</td>
<td>120°</td>
<td>110</td>
<td>113</td>
<td>114</td>
</tr>
<tr>
<td>CFC-11</td>
<td>45°</td>
<td>42</td>
<td>50</td>
<td>56</td>
</tr>
<tr>
<td>CFC-12</td>
<td>100°</td>
<td>81</td>
<td>86</td>
<td>93</td>
</tr>
<tr>
<td>CFC-113</td>
<td>85°</td>
<td>73</td>
<td>79</td>
<td>85</td>
</tr>
<tr>
<td>HCFC-22</td>
<td>12.9°</td>
<td>12.6</td>
<td>10.9</td>
<td>11.7</td>
</tr>
<tr>
<td>CH3CCl3</td>
<td>5.0°</td>
<td>5.3</td>
<td>4.7</td>
<td>5.1</td>
</tr>
<tr>
<td>CH3Cl</td>
<td>1.5°</td>
<td>1.5</td>
<td>1.3</td>
<td>1.4</td>
</tr>
<tr>
<td>CH3Br</td>
<td>1.7°</td>
<td>1.8</td>
<td>1.5</td>
<td>1.7</td>
</tr>
</tbody>
</table>

* Lifetimes are in years.

*From Intergovernmental Panel on Climate Change [2001, Table 4.1]. For CH4, this includes reactions with OH and O(1D). It does not include soil sink term. For N2O, this lifetime is the global mean atmospheric lifetime, not the perturbation lifetime.

*From WMO [2003, Table 1.6]. For CH3Cl, this lifetime only includes reaction with OH. Additional sinks (e.g., chlorine radicals in the marine boundary layer, soil uptake, or degradation in seawater) reduce this lifetime to 1.3 years.

*From Happel and Roche [2003]. For CCl3, this includes only the atmospheric lifetime (i.e., does not includes ocean or soil uptake). Including these other sinks reduces the lifetime 20 years. The CCl3 atmospheric lifetime is based on Volk et al. [1997].

*From Yvon-Lewis and Butler [1997]. For CH3Br, this lifetime is due to in situ reactions only. It does not include additional sinks (e.g., ocean uptake), which would reduce the lifetime to 0.7 years.
c patriotism in the MZ3-ECMWF simulations. The dry and wet phases of the tape recorder signal also propagate too rapidly in the MZ3-ECMWF simulations. In Figure 7a, the phase lag for the dry phase of the tape recorder signal is shown for all three simulations and the UARS climatology (see Waugh and Hall [2002] for discussion of phase lag and model results). It takes about 4 months and 8 months to propagate the dry signal up to 20 hPa in the MZ3-Op and MZ3-EXP471 simulations, respectively. Observations suggest a period of 14 months. In the MZ3-WA1b simulation, the propagation timescale of the tape recorder signal to 20 hPa is a little over one year, slightly less than observations would suggest. In Figure 7b, the attenuation of the H₂O vapor abundance is shown. Here the MZ3-EXP471 simulation is attenuated most rapidly between 100 hPa and approximately 20 hPa relative to the other simulations. The MZ3-WA1b simulation is attenuated slightly more than

Figure 6. Height-time sections of water vapor volume mixing ratio derived in the equatorial region (averaged between 12°S and 12°N). Observations are based on HALOE data. Two consecutive seasonal cycles are shown in each diagram.

Figure 7. Vertical profiles of transport diagnostics for water vapor averaged between 12°S and 12°N latitude. (a) The phase lag of the dry phase of the vertically propagating H₂O vapor abundance. (b) The attenuation of the H₂O vapor abundance. The attenuation is forced to be unity at 100 hPa for cases.
Figure 8. Zonal mean latitude-altitude cross sections of CH$_4$, NOy, and O$_3$ in January. Observations are based on UARS climatology. Both observations and model results have been mapped to an equivalent latitude pressure grid.
A simple method for assessing the representation of transport processes in the model is to analyze the distribution of relatively inert tracers. In Figure 8, we show as a function of latitude and height the zonally averaged mixing ratio of CH$_4$, NO$_y$, and O$_3$ between 100 hPa and 0.1 hPa in January. Figures 8a shows values measured by the Upper Atmosphere Research Satellite (UARS), Figures 8b-8d represent the distributions calculated when the model is driven by the Op, EXP471, and WACCM1b winds and temperatures, respectively. In Figure 9, the difference in CH$_4$, NO$_y$, and O$_3$ relative to the UARS climatology is shown for the three simulations. The differences between the distributions highlight the influence of uncertainties in the three-dimensional transport on the calculated tracer distributions. The observational climatology can help in
evaluating the model results and, specifically, the simulated transport.

[36] When comparing the fields of CH₄ obtained for the three sets of dynamical meteorological fields, it immediately appears that the transport generated by the Op analyzed winds leads to meridional tracer distributions that are somewhat “smoother” in latitude than when the winds produced by either EXP471 or WACCM1b are adopted. This is also seen in the difference plot for the Op case, which has larger CH₄ values in the extratropics. There is also less CH₄ in the deep tropics between 10 hPa and 5 hPa. This is consistent with less CH₄ reaching this altitude because of excessive horizontal transport below this altitude range. In the WACCM1b case, the dynamical barriers separating tropical from midlatitude air, and midlatitude from polar air are more pronounced, in particular in the winter Northern Hemisphere. In January (Figure 8), for example, a clear “surf zone” located between the dynamical barriers near 15°N and 50°N can be seen in the tracer distributions only when the WACCM1b winds are used. This feature is also clearly visible in the observed zonal monthly mean CH₄ distribution, although it is more pronounced in the MZ3-WA1b simulation. In the MZ3-WA1b case, the upward branch of the Brewer-Dobson circulation seems to be very narrow around the equator, and the associated vertical flow of air appears to be stronger, or the horizontal mixing weaker, than suggested by the observations of tracers like CH₄. This will be discussed in more detail below. In the tropical middle stratosphere, the absolute mixing ratio of CH₄ calculated with EXP471 fields is closer to the observed values than when calculated with WACCM1b meteorology. In the extratropics, the observed mixing ratio of CH₄ in the lower stratosphere (100-50 hPa) is best reproduced in January when the model is driven by WACCM1b meteorology. However, in the MZ3-WA1b simulation polar CH₄ volume mixing ratios in Northern Hemisphere winter are significantly lower than suggested by the UARS climatology. This feature in the MZ3-WA1b simulation is indicative of strong descent within the polar vortex and/or less mixing with midlatitude air. However, Northern Hemisphere polar descent is known to have strong interannual variability; therefore any agreement (or lack of) with respect to an observational climatology is not particularly revealing. Conversely, the high extratropical CH₄ concentration values produced by the MZ3-Op and MZ3-EXP471 cases are suggestive of too rapid meridional transport of air masses from the tropics. The circulation and transport of tracers appear therefore to be very different in the three model simulations, and the differences are broadly consistent with the calculated mean age of air. In the MZ3-WA1b case, the mean age of air (Figures 4 and 5) that reaches the polar lower stratosphere during winter is considerably older (4.5 years at 20 km) than in the Op case (2 years at 20 km), which suggests that the meridional exchanges derived from this latter data set are substantially faster than in the former case. This conclusion regarding the MZ3-Op meteorological fields is consistent with the findings of van Noije et al. [2004].

[37] In Figures 8 (center) and 9 (center), total inorganic nitrogen (NO₃) is shown. There is good agreement between the peak mixing ratio of NO₃ (NO + NO₂ + 2*N₂O₅ + HNO₃) for the MZ3-WA1b simulation (approximately 18 ppbv), and the peak NO₃ (NO₃ss + NO₂ss + HNO₃) derived from UARS observations. This approach of comparing model NO₃ (as defined above) with NO₃ derived from observations of HNO₃ summed with NO and NO₂ sunset data follows the recommendations of Park et al. [1999]. In contrast, the peak NO₃ abundance is approximately 14 ppbv and 16 ppbv for the MZ3-Op and MZ3-EXP471 simulations respectively. The MZ3-Op simulation does a poor job of representing the NO₃ vertical and horizontal structure relative to the UARS climatology. The NO₃ isopleths are too flat relative to observations, consistent with the mean age-of-air results shown in Figure 4. At pressures equal to 10 hPa and less (lower in altitude), all three simulations underestimate NO₃ (and NO₂, not shown). This affects the O₃ distributions discussed below.

[38] In the case of O₃, the calculated latitude-height distributions for the month of January shown in Figure 8 (and the difference distributions in Figure 9) show noticeable differences from the UARS climatology. The maximum O₃ abundance in all model simulations in the tropical stratosphere near 10 hPa is approximately 1–2 ppbv higher than the UARS climatology. The O₃ abundance is also high in the middle stratosphere extratropics. The high bias in model O₃ is consistent with the low bias in NO₃. One can clearly see this when comparing Figures 8(middle) and 8(right) with (Figures 9(middle) and 9(right).

[39] It is well known that the main stratospheric production of NO₃ is from the reaction of O³ (1D) with nitrous oxide (N₂O) forming two nitric oxide (NO) molecules. In Figure 10a–10c, the annual average production rate of NO₃ is shown for cases 1–3. There clearly is a progression of increasing NO₃ production rate from MZ3-Op to MZ3-EXP471 to MZ3-WA1b. The production rate of NO₃ is primarily influenced by the transport of N₂O (not shown, but similar to the CH₄ distribution shown in Figure 8). For MZ3-Op simulation, since less N₂O reaches the middle stratosphere, this has the net effect of producing too little NO₃ globally less than 1 hPa (lower in altitude). The altitude where the peak NO₃ production rate is a maximum also shifts upward when going from MZ3-Op to MZ3-EXP471 to MZ3-WA1b, where the peak NO₃ production rate is at approximately 15 hPa in MZ3-Op, at 12 hPa in MZ3-EXP471, and at 8 hPa for MZ3-WA1b. The tropical abundance of NO₃ relative to the N₂O abundance is shown in Figure 10d. The N₂O abundance can be considered an altitude coordinate, with altitude increasing with decreasing abundance of N₂O. As the N₂O abundance is reduced, more NO₃ is formed, at least up to the peak of the NO₃ versus N₂O curve. At some point, NO₃ loss becomes the dominant processes and NO₃ decreases with decreasing N₂O. What is clear from Figure 10d is that the N₂O altitude coordinate is not consistent across simulations, that is, the transport processes that control N₂O are very different and the rate of oxidation of N₂O to NO₃ varies considerably among the simulations considered in this paper.

[40] In Figure 11, the Transformed Eulerian Mean (TEM) vertical velocity is shown for four seasons, averaged between 20°S and 20°N latitudes (see Brassey and Solomon [2005, and the references therein] for discussions of the TEM velocities). In this figure, the derived TEM vertical velocities are fairly consistent amongst the three simulations in the 100 hPa to 50 hPa altitude region; however, the MZ3-
WA1b simulation departs from the other two simulations near 20 hPa. At lower pressures (higher altitude), the MZ3-WA1b vertical velocities can be up to a factor of two greater. This is especially true during the equinox seasons (MAM and SON). The larger tropical TEM vertical velocity in the 20 hPa to 5 hPa region in the MZ3-WA1b simulation is consistent with more rapid vertical transport of N$_2$O and, as a consequence, an NO$_y$ production region located at higher altitude than in the two ECMWF simulations.

We will now investigate in more detail the processes responsible for the differences in the latitudinal and temporal evolution of ozone near the 10 hPa ozone peak. In Figure 12a, the O$_3$ time-latitude distribution is shown near the stratospheric O$_3$ peak (near 10 hPa) for the three model simulations. The corresponding UARS climatology time-latitude distribution is shown in Figure 13. The UARS climatology shows the maximum in O$_3$ at the spring and fall equinoxes with values approaching 10.5 ppmv. The MZ3-WA1b simulation time-latitude peak O$_3$ distributions are also centered on the equator where NO$_x$ catalytic cycle is the main odd oxygen loss process in this region.

In Figure 12b, the time-latitude distribution of NO$_y$ is shown for the three model simulations (near 10 hPa). The corresponding NO$_y$ UARS climatology is shown in Figure 13. The MZ3-WA1b NO$_y$ time-latitude minima are also centered on the equator and significantly underestimate NO$_y$ by ~4–5 ppbv relative to the UARS climatology. The MZ3-ECMWF NO$_y$ time-latitude minima values are more similar to the UARS climatology, but also significantly underestimate the NO$_y$ abundance (~3–4 ppbv). For the MZ3-WA1b simulation, the tropical, 20 hPa to 10 hPa region is too isolated (e.g., see CH$_4$ distributions in Figures 8 and 9), allowing low NO$_y$ values to be transported from the lower stratosphere to the midstratosphere, without mixing in higher extratropical NO$_y$ air. This is clearly seen in the MZ3-WA1b latitude-height NO$_y$ distributions (Figure 8). The MZ3-EXP471 simulation in the tropical stratosphere near 10 hPa is less isolated (e.g., see CH$_4$ isopleths (Figures 8 and 9); and there is additional abundance of NO$_y$ in the 10 hPa region in the MZ3-EXP471 case relative to the MZ3-WA1b case. The MZ3-Op simulation also underestimates the NO$_y$ abundance relative to the UARS climatology, but not as much as the MZ3-EXP471 and MZ3-WA1b simulations. Nonetheless, in all three simulations, the NO$_y$ abundance at the equatorial O$_3$ maximum is significantly underestimated.

In Figure 12c, the time-dependent TEM vertical velocities are shown between 50°S and 50°N latitudes. One can clearly see that the vertical velocities for all three simulations evolve in a manner consistent with NO$_y$ distributions in this figure. For example, in the MZ3-Op simulation, NO$_y$ is high where the TEM vertical velocities are downward (negative in Figure 12), bringing higher
abundances of NOx from above to the 10 hPa level. When NOy abundances are higher, O3 abundances are lower, consistent with increased catalytic odd oxygen loss from the NOx family. The MZ3-EXP471 simulation also shows a similar relationship between the evolution of the TEM vertical velocities and NOy abundances. However, it should be pointed out that the MZ3-EXP471 TEM vertical velocities tend to be less variable (i.e., less “noisy”) than the MZ3-Op values. The MZ3-WA1b simulation also shows a strong correlation between the TEM vertical velocities and NOx abundances. However, the MZ3-WA1b TEM vertical velocities are larger (as shown in Figure 11) and the temporal distribution of the vertical TEM velocities is mostly upward in the tropics.

In Figure 12d, the 10 hPa latitudinal and temporal evolution of the TEM horizontal velocities are shown. Here a negative velocity designates flow toward the south; a positive velocity is designates flow toward the north. It is clear from this figure that the MZ3-WA1b simulation has the smallest TEM meridional magnitudes and temporal gradients. This is consistent with the high CH4 (Figure 8) and low NOx abundances in the MZ3-WA1b simulation being isolated in the deep tropics.

The interesting behavior of the NOy distribution in the MZ3-WA1b simulation relative to the MZ3-ECMWF simulations can be partially reconciled if one remembers the WACCM1b meteorological fields do not resolve a QBO structure in the zonal winds (Figure 1). The secondary meridional circulation that accompanies the QBO is downward in the region of westerly shear. This will tend to transport high NOx from above and consequently reduce ozone. This effect has been discussed recently by Tian et al. [2006] who used a chemistry climate model that internally generated a QBO. One of the main conclusions from this study was that variations in NOy (and NOx) induced by the secondary circulation of the QBO were responsible for modulating O3 in the 10 hPa region of the tropical stratosphere. We have obtained similar results with a fully interactive version of WACCM, run with and without an imposed QBO (not shown). The simulation without a QBO was very similar to the MZ3-WA1b simulation in this paper, whereas the simulation that included a QBO resembled the MZ3-ECMWF results, especially for those years that had the same phase of the QBO as the year 2000 ECMWF meteorological fields.

4.3. Column Ozone and Polar Processes

An examination of the behavior of tracers at high latitudes in both hemispheres provides additional information on the ability of the WACCM1b and ECMWF wind and temperature fields to properly simulate the transport of long-lived species. In the case of polar H2O (Figure 14), the model simulations are compared with POAM III observations from the year 2000 and a POAM III climatology (1998–2005 average). In Figure 14, model results have been sampled at the POAM III measurement locations. The POAM III H2O in 2000 is shown for comparison to the MZ3-ECMWF simulations, while the MZ3-WA1b simulation is not specific to any 1 year and is more appropriately compared to the POAM III climatology. In the Arctic, the

**Figure 11.** Vertical profiles of seasonal TEM vertical velocities (cm s\(^{-1}\)) for seasons (a) December–January–February (DJF), (b) March–April–May (MAM), (c) June–July–August (JJA), and (d) September–October–November (SON).
Figure 12. Latitude-time distributions at 10 hPa for O$_3$ (ppmv), NO$_y$ (ppbv), TEM vertical velocity (cm s$^{-1}$), and TEM meridional velocities (cm s$^{-1}$): (left) MZ3-Op, (middle) MZ3-EXP471, and (right) MZ3-WACCM1b.
Figure 13. Zonal mean latitude-time sections of ozone (ppmv) and NO$_x$ (ppbv) at 10 hPa. Observations are based on UARS HALOE and CLAES data. See text for details.

Figure 14. Height-time cross sections of water vapor volume mixing ratio at the POAM III measurement locations. POAM III water vapor is shown for year 2000 (Figure 14b) and climatology made of years 1998–2005 (Figures 14c-14f) from the model simulations (c) MZ3-Op, (d) MZ3-EXP471, and (e) MZ3-WA1b for high-latitude (left) Northern Hemisphere and (right) high-latitude Southern Hemisphere data and model results.
mixing ratio of H\textsubscript{2}O does not change substantially during the year along the measurement track of POAM III at any altitude. The weak maximum of near 7 ppmv in the POAM III H\textsubscript{2}O climatology between 1 hPa and 5 hPa from August to November is not captured when the Op winds are adopted, but is reproduced when the EXP471 and WACCM1b fields are used. The H\textsubscript{2}O minimum observed near 100 hPa from March to August in the Northern Hemisphere in the POAM III climatology is not low enough in the MZ3-WA1b or the MZ3-Op simulation, but is better simulated when adopting ECMWF EXP471 meteorological fields. Overall, the MZ3-EXP471 simulation of Arctic H\textsubscript{2}O shows the best agreement with the POAM III climatology. In the Antarctic, the descent of high H\textsubscript{2}O volume mixing ratios (6–8 ppmv) from the stratopause to 10 hPa between January and June is reproduced in both the MZ3-WA1b and MZ3-EXP471 simulations, but the model H\textsubscript{2}O volume mixing ratios are smaller than in POAM III. The MZ3-Op simulation does a poor job of representing this feature. All model simulations produce realistic dehydration in the Southern Hemisphere lower stratosphere in spring, however, the MZ3-WA1b dehydration is more severe and extends higher in altitude than in the MZ3-ECMWF cases and in the POAM III observations.

It is important to assess how the dynamical fields affect the seasonal evolution and latitudinal distribution of the total ozone column, and specifically the formation and dissipation of the springtime O\textsubscript{3} hole in Antarctica. Here, the model simulations will be compared with data provided by the EP/TOMS instrument for the year 2000. The general patterns derived by the model (Figure 15) for the three meteorological simulations are relatively similar, and characterized by low tropical values, maximum values during spring in the Arctic and near 50°S, and minimum values during springtime in the Antarctic. At most locations, the magnitude of the column calculated with the WACCM1b meteorological fields is, however, in better agreement with the observations than the values calculated with the Op and EXP471 meteorological fields. The MZ3-WA1b does an adequate job representing tropical total column ozone. In the MZ3-Op and MZ3-EXP471 simulations, the tropical total column ozone is approximately 25 DU higher (10%) than observed.

The Southern Hemisphere O\textsubscript{3} in the polar lower stratosphere is shown in Figure 16. Here, the model simulations are sampled at the POAM III measurement locations and daily results are shown for three different potential temperature surfaces: 400 K (~15 km); 450 K (~18 km); and 500 K (~20 km). Since air motions are nearly isentropic in the lower stratosphere, it is useful to examine the evolution of ozone using this vertical coordinate. The solid black line is for POAM III year 2000 (for comparison to the ECMWF runs). The solid gray line is the average for the entire POAM III data record (for comparison to the
The differences in the O$_3$ hole simulations noted above when using the three sets of meteorological fields are likely related to differences in temperature and/or in transport. To examine this in more detail, the annual evolution of Southern Hemisphere total column ozone and local constituents at 84°S are shown in Figure 17. This figure displays the temporal evolution of total column ozone, plus local change at 43 hPa (~20 km) in O$_3$, inorganic reservoir species (HCl, ClONO$_2$), total inorganic chlorine (ClO$_x$ = ClO + Cl + ClO + ClO$_2$ + OCl + HOCl), HClO, and HNO$_3$. The figure shows that the available inorganic chlorine contained in HCl and ClONO$_2$ is significantly activated to ClO$_x$ species in all three simulations. However, ClO$_y$ is not the same in the three simulations. There is a larger abundance of ClO$_y$ in the MZ3-EXP471 simulation, followed by the MZ3-EXP471, and MZ3-Op simulations. The amount of ClO$_x$ is strongly dependent on the dynamics and transport, and is proportional to the mean age of air (Figures 4 and 5) [Waugh and Hall, 2002]. The MZ3-EXP471 had the oldest mean age of air in the polar latitudes and, consistent with, the greatest amount of ClO$_y$. The MZ3-EXP471 simulation underestimated the ClO$_y$ relative to MZ3-EXP471 case by 0.2–0.5 ppbv. However, at this altitude prior to 1 October (0.75 year fraction) the temporal evolution of the temperatures is very similar in MZ3-EXP471 and MZ3-EXP471 (Figure 3). This condition, along with similar ClO$_x$, is why the ClO$_y$ temporal distributions are consistent and the O$_3$ loss is comparable (but not identical) between the two simulations. A similar conclusion can be made for BrO, and BrO$_x$ species (not shown). It is also interesting to compare the temporal evolution of H$_2$O and HNO$_3$. In Figure 3, the final warming in the Southern Hemisphere polar stratosphere occurs in mid-October when using ECMWF meteorological fields; the WACCM1b case never experienced an abrupt final warming. This allowed the MZ3-EXP471 Southern Hemisphere polar vortex to stay isolated longer and slowed the O$_3$ recovery. The persistent MZ3-EXP471 Antarctic vortex also affected the late spring/summer distributions of ClONO$_2$, HNO$_3$, and H$_2$O. There was significant denitrification during the winter-spring period and, without replenishment of inorganic nitrogen species from outside the vortex, ClONO$_2$ and HNO$_3$ could not rapidly recover to prewintertime values (contrary to what occurred in the MZ3-ECMWF simulations). For H$_2$O, dehydration processes stopped in midwinter and the amount of H$_2$O in the vortex remained relatively constant until mixing of H$_2$O rich midlatitude air occurred later in the year (see Figure 14). Overall, the more realistic representation of the final warming in the MZ3-ECMWF simulations (Figure 3) explains the better simulation of the dehydration and recovery of H$_2$O relative to observations (see also Figure 14).

5. Summary and Conclusions

[50] In order to assess the influence of meteorological input parameters on the calculated atmospheric composition of the atmosphere, we have used an atmospheric chemical transport model driven by three different dynamical fields. The first dynamical data set comes from simulations performed with the Whole Atmosphere Community Climate Model, version 1b. The second and third sets were provided by the European Centre for Medium-Range Weather Forecasts, and included operational products for year 2000 and a
recent reanalysis, called EXP471, which is an update for year 2000 of the ERA-40 reanalysis. The fields retained in the case of WACCM1b correspond to a specific simulation year that bears resemblance to observed conditions in year 2000; however, it should be emphasized that it is not a detailed simulation of this or any other specific year.

A detailed analysis of the chemical fields simulated by MOZART-3 shows that, even though the general features characterizing the three dynamical sets are rather similar, slight differences in winds and temperature can produce substantial differences in the calculated distributions of chemical tracers. This study shows that, in most cases, the ECMWF operational wind fields produce too strong latitudinal mixing and therefore too weak dynamical barriers against meridional transport. This behavior is reflected in the calculated mean age of air, which at 20 km in the extratropics is 1–3 years lower than the values derived from the observations of long-lived tracers. The EXP471 reanalysis benefits from many of the improvements made to ECMWF’s operational analysis system since the year 2000, and produces much more realistic distributions of stratospheric tracers consistent with much improved values of the mean age of air. The tracer distributions calculated when the WACCM1b meteorological fields are specified adequately represents observed tracer distributions and the mean age of air. Meteorological quantities provided by general circulation models with no observational constraints seem in many cases to provide more realistic results than quantities resulting from data assimilation procedures. This has been shown in previous studies to be due to undesirable noise or imbalance generated by the assimilation procedure.

Figure 17. Annual cycles of total column ozone (DU), Ozone (ppmv), HCl (ppbv), ClONO₂ (ppbv), total inorganic chlorine (ClOy), active inorganic chlorine (ClOX), water vapor (ppmv), and HNO₃ (ppbv) at 84°S on the 43 hPa pressure surface. Black line is for MZ3-WA1b; blue line is for MZ3-Op; and red line is for MZ3-EXP471.
which has considerable effects on the model-simulated transport of chemical species.

The photochemical lifetimes for selective short-lived (<10 years) and long-lived species are also shown for the three model simulations. The MZ3-WA1b simulation produces lifetimes of long-lived species (e.g., CFC-11, CFC-12, and N$_2$O) that are close to the values derived from observations, followed by the MZ3-EXP471 and MZ3-Op simulations respectively. In contrast, the CH$_4$ photochemical lifetime, which is less than 10 years and controlled mainly by the reaction with tropospheric OH, is best represented by MZ3-EXP471 simulation.

Even though the tropical total ozone columns calculated in the three model cases are in fair agreement with observations, the calculations of vertical and horizontal exchanges in the tropics are not satisfactory. In the tropical lower-to-middle stratosphere, there are differences between the three model cases when the H$_2$O tape recorder signal is analyzed. The phase lag of the dry component of the tape recorder signal is best represented by the MZ3-WA1b simulation. Both the MZ3-ECMWF simulations propagate the tape recorder signal too rapidly compared to observations. In the tropical midstratosphere region (near 10 hPa), the temporal evolution and magnitude of NO$_3$ is very different between the MZ3-ECMWF simulations and the MZ3-WA1b simulation. In the latter case, minimum NO$_3$ values are centered on the equator. These values tend to underestimate substantially the observed UARS climatology. Since O$_3$ is primarily controlled by NO$_3$ chemistry in this region, the O$_3$ distribution for the MZ3-WA1b simulation is centered on the equator and is overestimated relative to the UARS climatology. The MZ3-ECMWF simulations minimum NO$_3$ values are not centered on the equator and their distributions are influenced by the QBO secondary circulation for year 2000. The O$_3$ abundance in these simulations also follows the latitudinal and temporal structure of NO$_3$. This was confirmed by examining the TEM vertical and horizontal velocities. If the WACCM1b parent GCM had included a QBO, the tracer distributions and O$_3$ would have been represented much better. We have recently confirmed the importance of the QBO circulation in simulations with a fully coupled version of WACCM, which will be the focus of a future study.

Large differences between model results are also apparent at high latitudes. In all cases, the total column ozone appears to be overestimated during all seasons north of 40°N and in the latitude band of 60°S where the southern hemisphere maximum in the total column ozone is observed. Again, this suggests that the Brewer-Dobson circulation in the three model cases is too strong. Only the MOZART-3 simulation driven with WACCM1b dynamical data produces total column ozone amounts in reasonable agreement with observations in the Southern Hemisphere.

The differences in the three meteorological data sets also lead to very different representations of the springtime Antarctic O$_3$ hole. Although the Southern Hemisphere polar winter and early spring lower stratosphere temperatures are similar in the three cases, the intensity of the meridional circulation and the degree of isolation of the polar vortex are different. In the two ECMWF cases, the final warming occurs at the time constrained by observations, while in WACCM1b, the transition between winter and summer conditions occurs in a considerably less abrupt fashion and takes place about a month later than observations for year 2000 would indicate. The depth of the O$_3$ hole is considerably underestimated in the two cases using assimilated dynamics, especially in the case of the operational data. This is attributed to the excessively vigorous transport of O$_3$, which partly replenishes the O$_3$ that has been catalytically destroyed by chlorine compounds in the polar region. The isolation of the polar vortex is more pronounced when WACCM1b data are used, so that more ozone is destroyed. In addition, the level of inorganic chlorine and bromine is quite different in the three cases under consideration. In simulations with relatively low mean age of air, the level of inorganic chlorine and bromine at high latitudes is lower than in cases with higher mean age of air. As a result, the calculated chemical O$_3$ destruction rate by halogen radicals is four times higher when the MOZART-3 model is driven by WACCM1b data than, for example, when the model uses the ECMWF operational data. In the MOZART-3 simulation using WACCM1b meteorological fields, the vortex remains stable for several weeks in spring and large O$_3$ depletion is derived. However, since the vortex remains isolated later into the spring, the O$_3$ hole persists until the end of November.

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